

GRAMSES: a new route to general relativistic N -body simulations in cosmology

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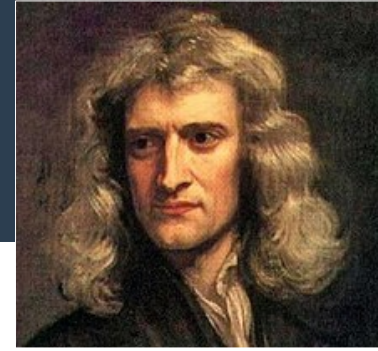
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Talk Layout

- Introduction and motivation
- A fully constrained formulation of GR
- GRAMSES overview
- Results from Λ CDM simulations
- On the generation of Initial Conditions
- Conclusions and outlooks

Cosmic Web à la Newton



In standard cosmological N -body simulations we solve a Newtonian system of equations:

Gravity sector

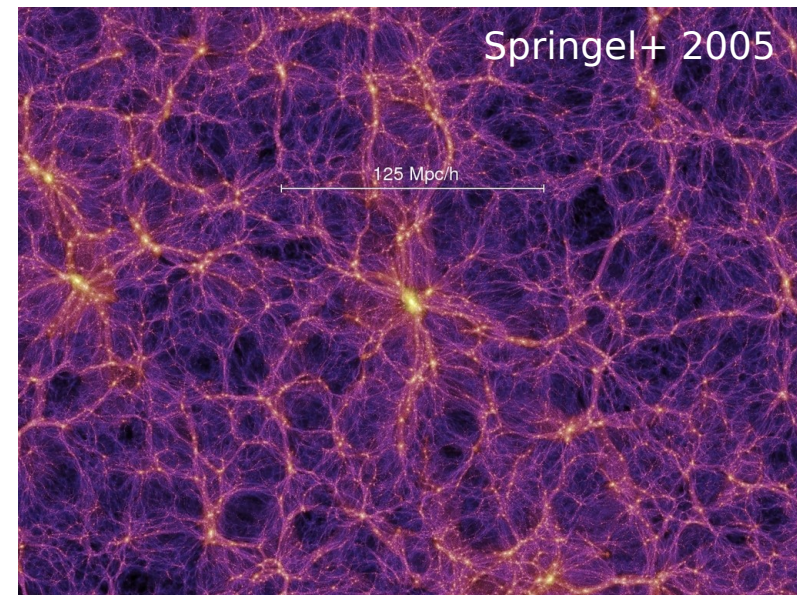
$$\nabla^2 \Phi = 4\pi G a^2 \delta\rho$$

Matter sector

$$\dot{v} + H v = -\nabla\Phi$$

where $a = a(\Omega_m, \Omega_\Lambda)$

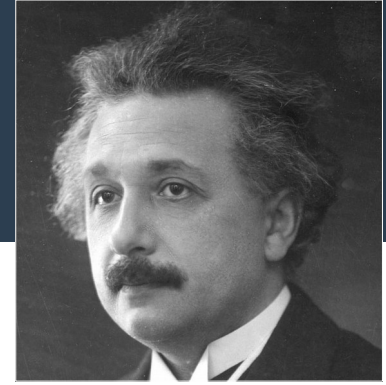
- ✓ Great success to explain our observable Universe.
- ✓ (Deceivingly?) Intuitive picture.
- ✗ Fixed FLRW background expansion.
- ✗ Structures do not back-react on spacetime.



Motivation for GR simulations

- In the era of precision cosmology, is Newtonian approximation going to remain sufficient for the upcoming years?
- Test whether **GR** makes any measurable difference on observables.
- How good is the **FLRW** background? Can we shed some light on the current 5.3σ tension in H_0 ? (Wong+ 2019).
- Cosmology with ‘exotic objects’ (beyond Λ CDM) such as relativistic particles, scalar fields, cosmic strings, gravitational waves, etc.

The Ultimate Cosmic Web



In principle, we would like to solve the full theory of gravity without any approximation:

Gravity sector

$$G_{\mu\nu} = T_{\mu\nu}$$

Matter sector

$$\nabla_{\mu} T^{\mu\nu} = 0$$

- ✓ Fully nonlinear and background-independent model for the Universe.
- ✓ Structures do interact with spacetime.
- ✗ Computationally expensive.
- ✗ Not very intuitive.

Known routes for GR cosmological simulations

- Perturbative Expansions route:
 - evolution (Adamek+ 2016)

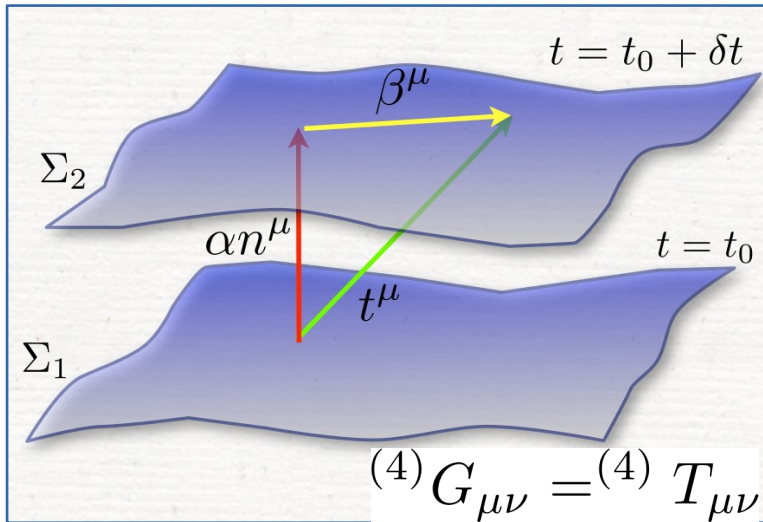
- ✓ Shell-crossing for DM.
- Linearisation procedure.
- No AMR (fixed resolution).

- Numerical Relativity route:
 - Bentivegna & Bruni 2016
 - Mertens+ 2017
 - Macpherson+ 2017

- ✓ Non-perturbative GR.
- No shell-crossing for DM.
- No Λ .

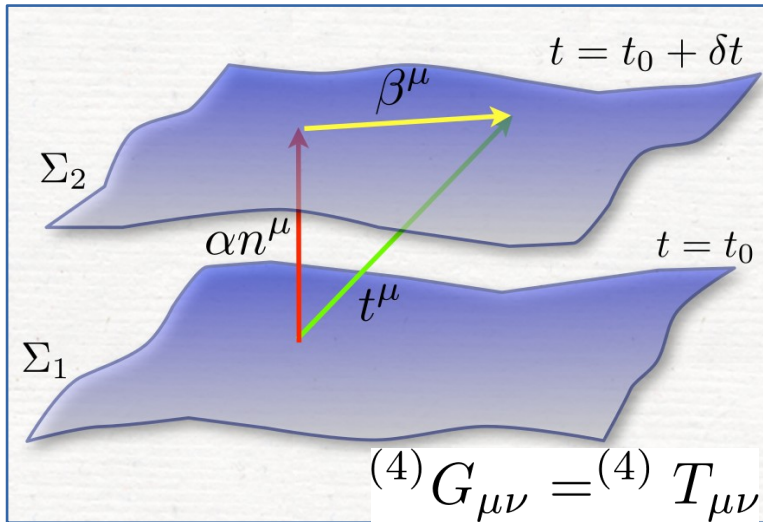
Let's try to have the best of both:
collisionless particles in non-perturbative GR.

Numerical Relativity: the 3+1 Formalism



$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij}(\beta^i dt + dx^i)(\beta^j dt + dx^j)$$

Numerical Relativity: the 3+1 Formalism



$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij}(\beta^i dt + dx^i)(\beta^j dt + dx^j)$$

CTT decomposition:

$$\gamma_{ij} = \psi^4 \bar{\gamma}_{ij},$$

$$\bar{A}^{ij} = \bar{A}_{TT}^{ij} + \bar{A}_L^{ij}$$

Evolution Equations:

$$(\partial_t - \mathcal{L}_\beta)\gamma_{ij} = -2\alpha K_{ij}$$

$$(\partial_t - \mathcal{L}_\beta)K_{ij} = -D_i D_j \alpha$$

$$+ \alpha(R_{ij} - 2K_{ik}K_j^k + K K_{ij})$$

$$- 8\pi\alpha \left[S_{ij} - \frac{1}{2}\gamma_{ij}(S - \rho) \right]$$

Constraint Equations:

$$R + K^2 - K_{ij}K^{ij} = 16\pi\rho$$

$$D_j(K^{ij} - \gamma^{ij}K) = 8\pi S^i$$

A Fully Constrained Formulation of GR

- However, the Einstein Equations in 3+1 are **weakly hyperbolic**, a problem that is fixed by introducing auxiliary variables as in the BSSN formalism (Shibata and Nakamura 1995, Baumgarte and Shapiro 1999).

A Fully Constrained Formulation of GR

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- For structure formation we might argue that tensor modes are subdominant.
- In order to kill tensor modes, we might take the approximation:

$$\bar{\gamma}_{ij} = \delta_{ij}, \quad \bar{A}_{TT}^{ij} = 0 \quad \forall t$$

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- In principle the conformal flatness approximation is very strong. However, we can use the gauge freedom so that we don't lose information about 1 scalar + 2 vector modes in the metric.

A Fully Constrained Formulation of GR

- Following Cordero-Carrion+ 2009, we can have a consistent fully constrained formulation by picking the Dirac gauge condition:

$$\partial_i \bar{\gamma}^{ij} = 0$$

- This gauge condition translates into a condition for the shift vector and makes the conformal flatness approximation ‘safe’.
- This is in the same spirit as the ‘waveless theories of gravity’ first explored by Isenberg in 1978 and later by Wilson and Mathews 1989.
- In this approach time evolution is due to particles (solving the geodesic equation).
- Even if no GW are originally present, these still can be added under certain approximations (Cordero-Carrion+ 2012).

System of equations for the gravity sector

In this approach the elliptic system of equations is decoupled into:

1. Momentum Constraint:

$$(\bar{\Delta}_L W)_i = 8\pi s_i \qquad \bar{A}_L^{ij} = (\bar{L}W)^{ij} \equiv \bar{A}^{ij}$$

2. Hamiltonian Constraint:

$$\bar{\nabla}^2 \psi = -2\pi\psi^{-1} s_0 - \frac{1}{8}\psi^{-7} \bar{A}_{ij} \bar{A}^{ij} + 2\pi\psi^5 \hat{\rho}_m,$$

3. Constant Mean Curvature (CMC) slicing condition:

$$\bar{\nabla}^2(\alpha\psi) = \alpha \left[2\pi\psi^{-1}(s_0 + 2s) + \frac{7}{8}\psi^{-7} \bar{A}_{ij} \bar{A}^{ij} + \psi^5 \left(\frac{5K^2}{12} + 2\pi\rho_\Lambda \right) \right] - \psi^5 \dot{K}$$

4. Determine shift vector associated to Minimal Distortion gauge (in CFA):

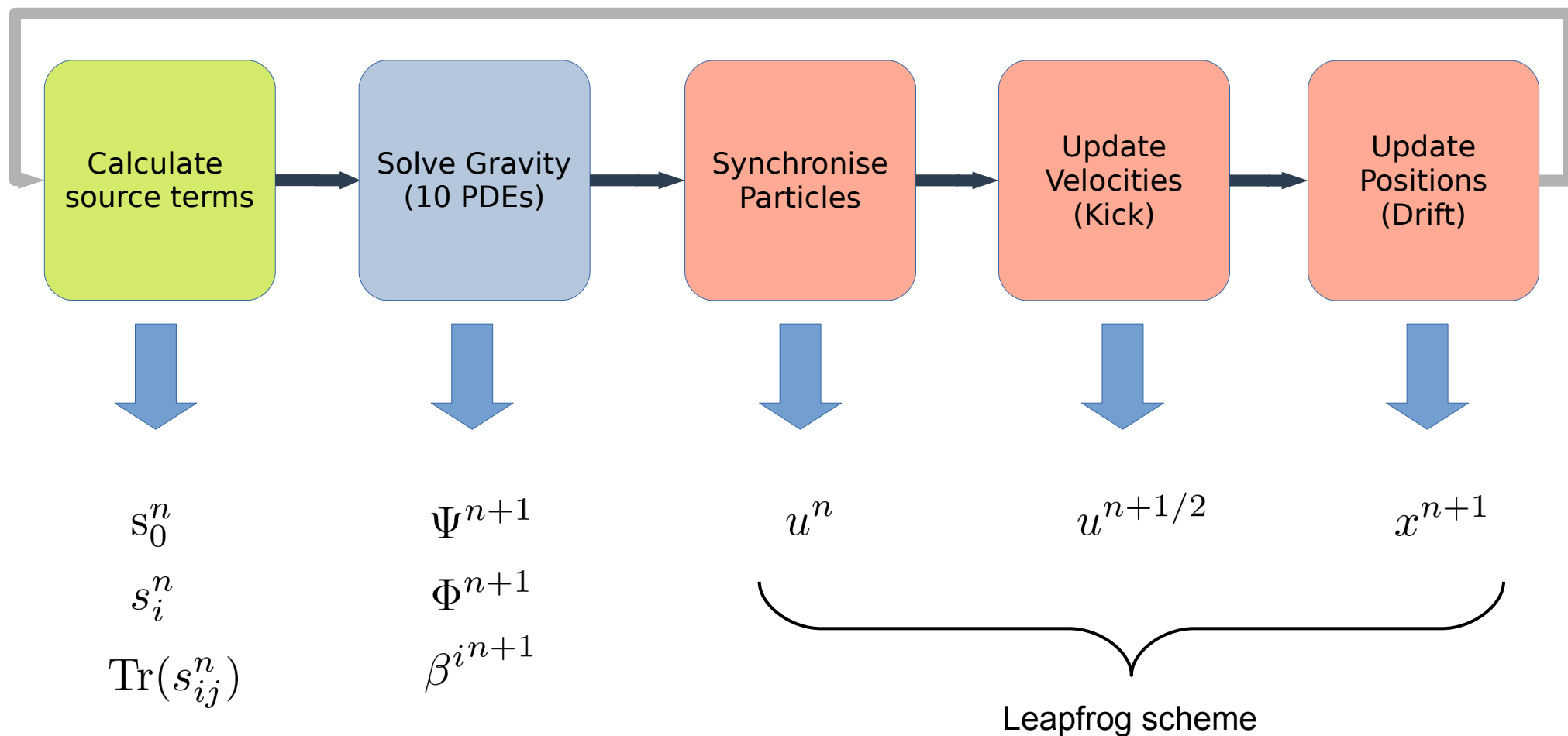
$$(\bar{\Delta}_L \beta)^i = 2\partial_j \left(\alpha\psi^{-6} \bar{A}_L^{ij} \right)$$

The **GRAMSES** code

- Based on (Newtonian) **RAMSES** code (Teyssier '02)
 - Grid-based code.
 - Parallelisation (MPI).
 - Multigrid + Adaptive Mesh Refinement (AMR).
 - Cloud In Cell (CIC) scheme for particles.
- The implementation/modification of several subroutines are needed for these GR simulations, including:
 - Calculation of relativistic source terms:
 - ♦ Matter source terms.
 - ♦ Geometric terms.
 - Multigrid Gauss-Seidel Solver for 8 Poisson-like equations.
 - Multigrid Gauss-Seidel Solver for 2 non-Poisson-like equations.

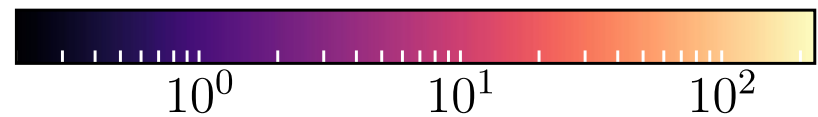
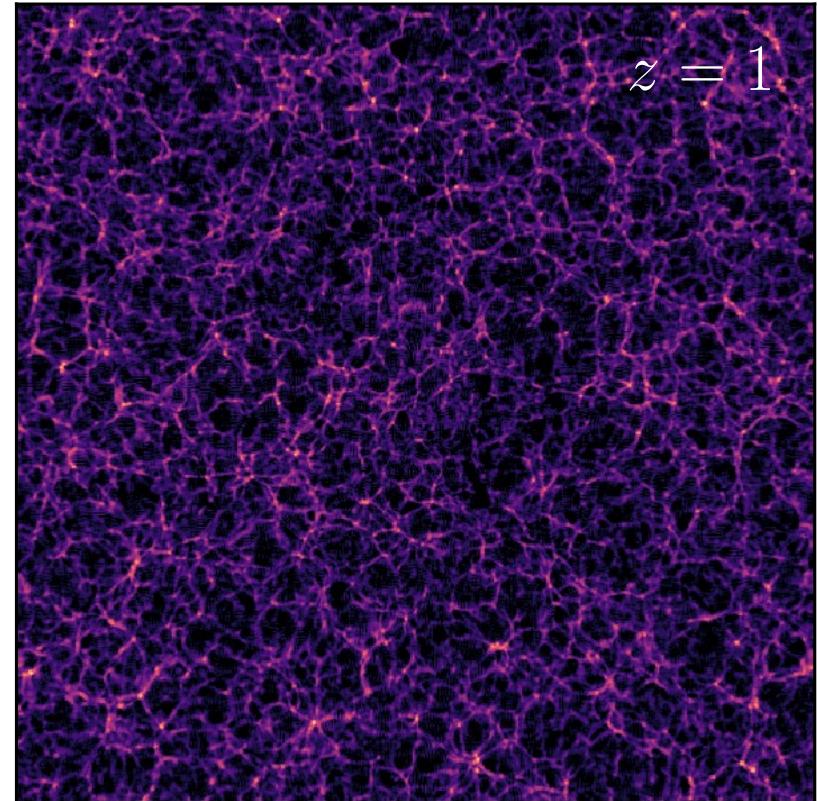
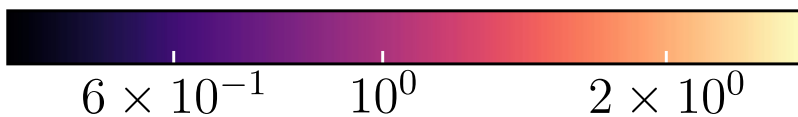
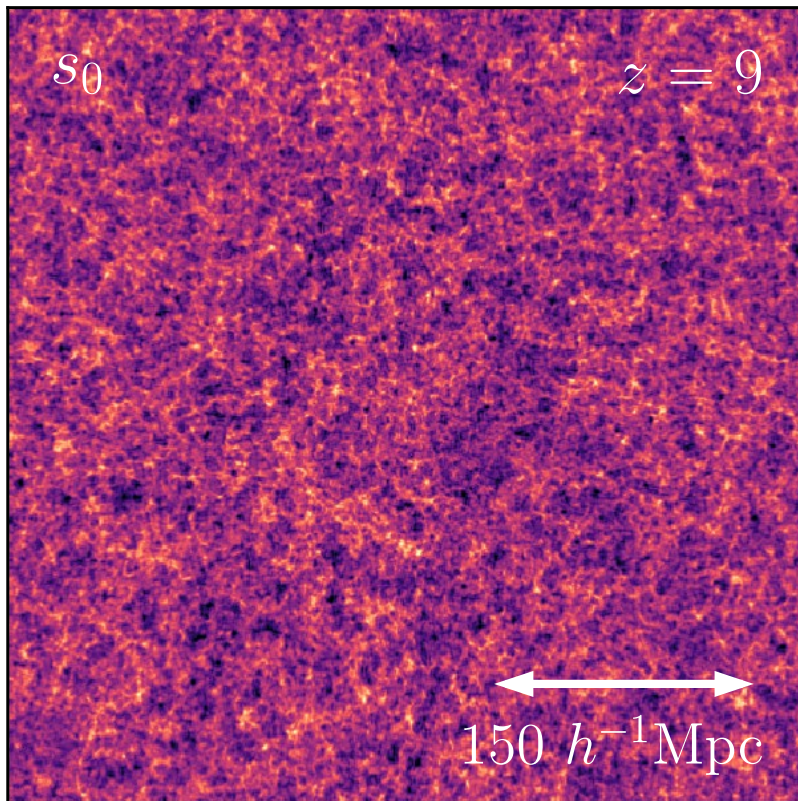
Flowchart of GRAMSES Solver

Given a snapshot of the system at a given time-step n , the solver follows the next logical sequence:



Cosmic web in a Λ CDM Universe

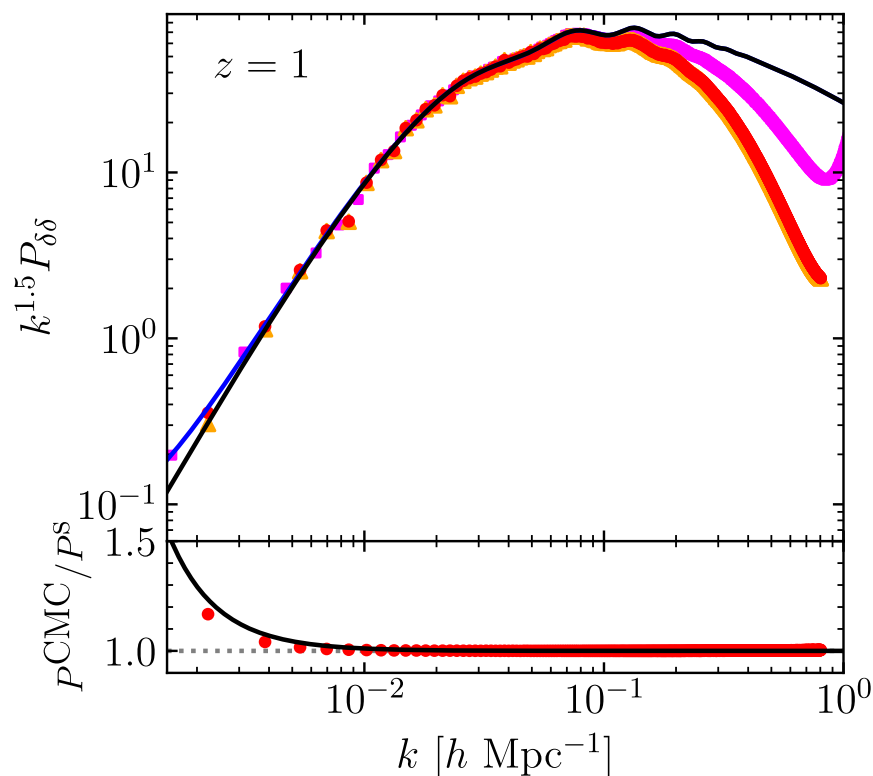
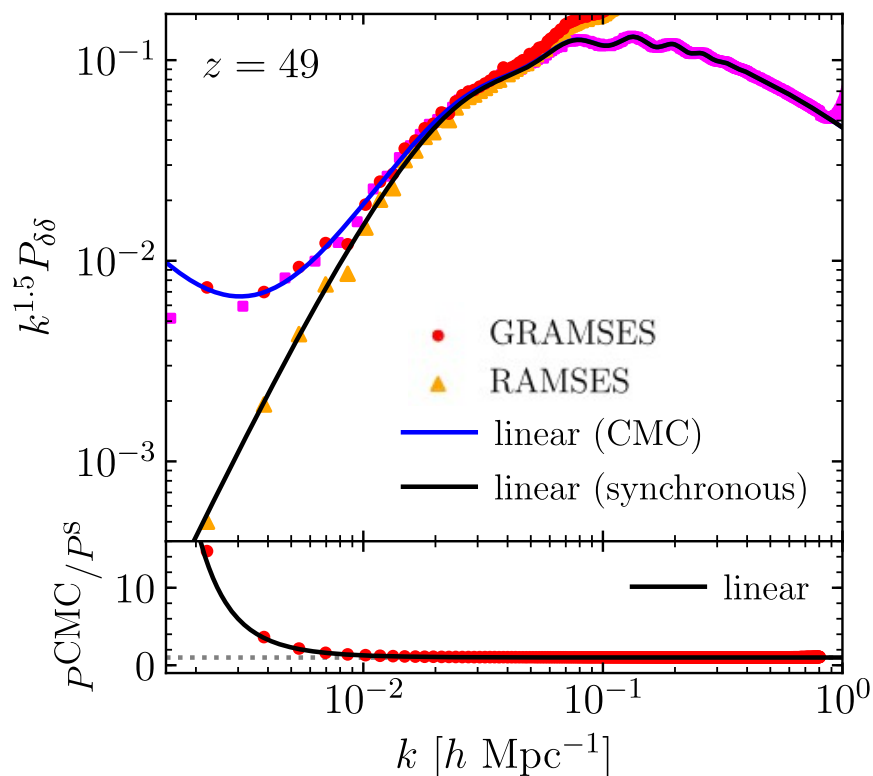
Setup: 512 Mpc/h box with 512^3 DM particles



Matter Power Spectrum

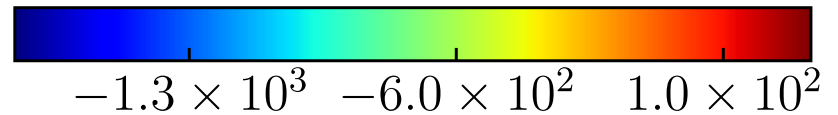
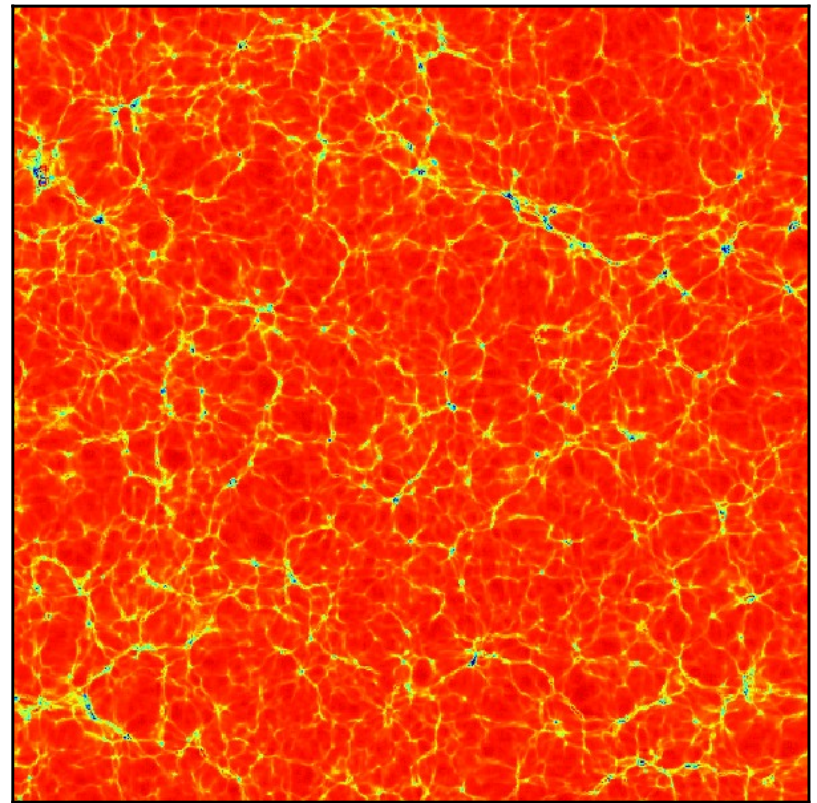
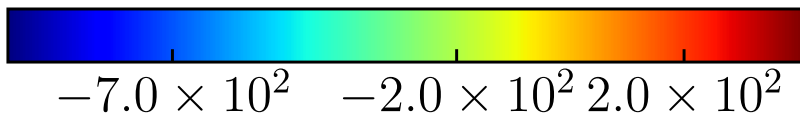
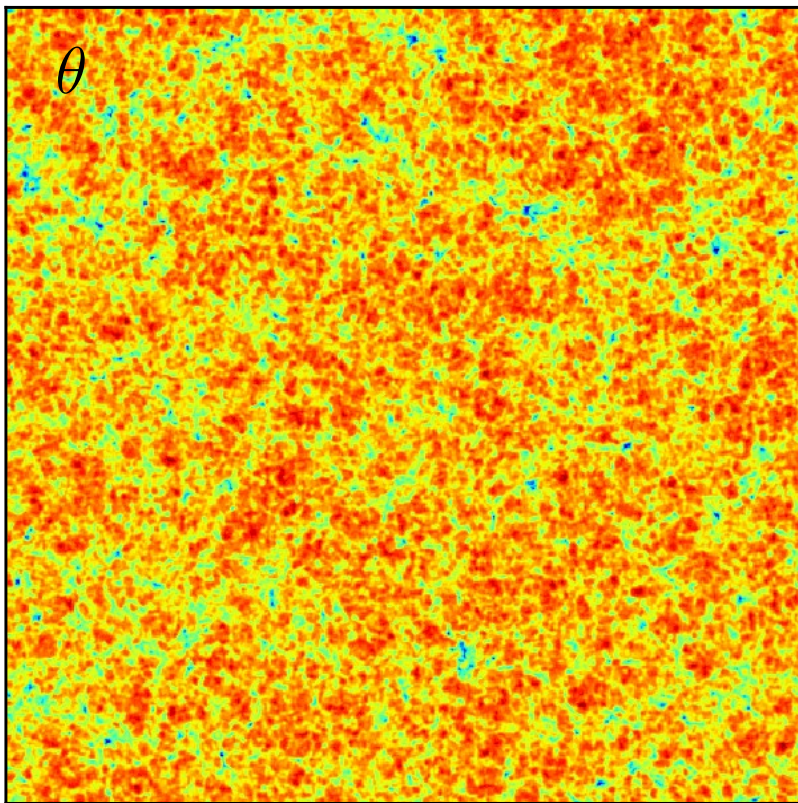
Setup: 4 Gpc/h box with 1024^3 DM particles.

$$\langle \delta(k)\delta(k') \rangle = \frac{2\pi^2}{k^3} P_{\delta\delta} \delta^{(3)}(k - k') \quad \leftarrow \text{gauge-dependent}$$



Velocity divergence

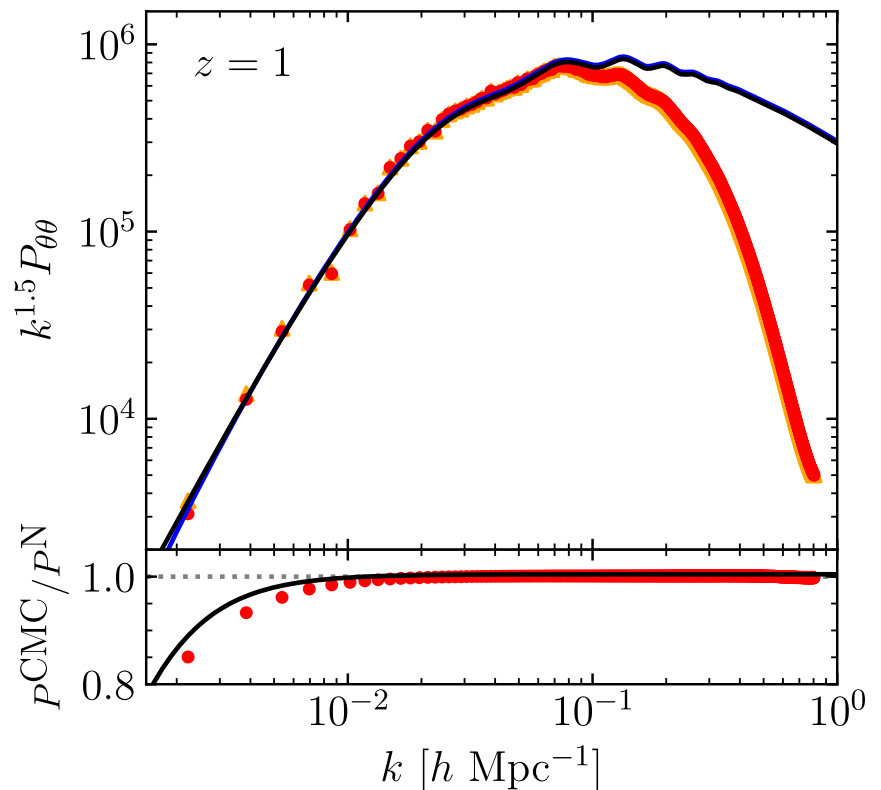
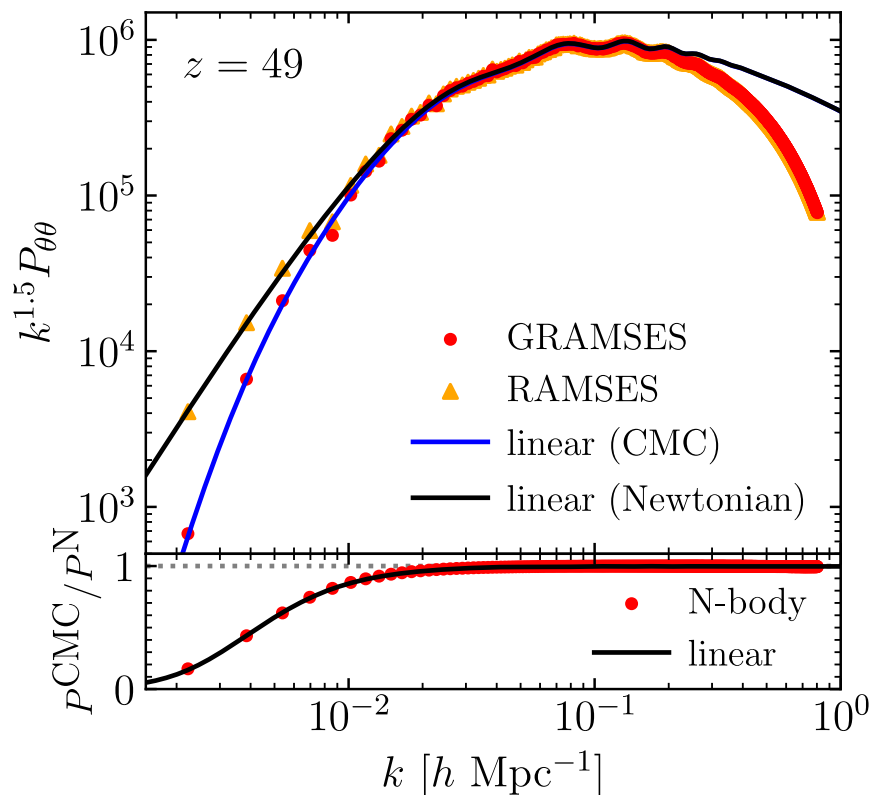
$$\theta = \nabla \cdot u$$



Velocity divergence Power Spectrum

Setup: 4 Gpc/h box with 1024^3 DM particles.

$$\langle \theta(k)\theta(k') \rangle = \frac{2\pi^2}{k^3} P_{\theta\theta} \delta^{(3)}(k - k') \quad \leftarrow \text{gauge-dependent}$$



Vector Potential

- The vector modes in the metric are responsible for a “frame dragging” effect which is absent in Newtonian gravity.
- In perturbation theory, β_i^V is a 2nd order quantity that is sourced by the rotational part of the momentum field.

$$\nabla^2 \beta_i^V = \frac{9\Omega_m^2 H_0^4}{2a^2} (\delta u_i)^V$$

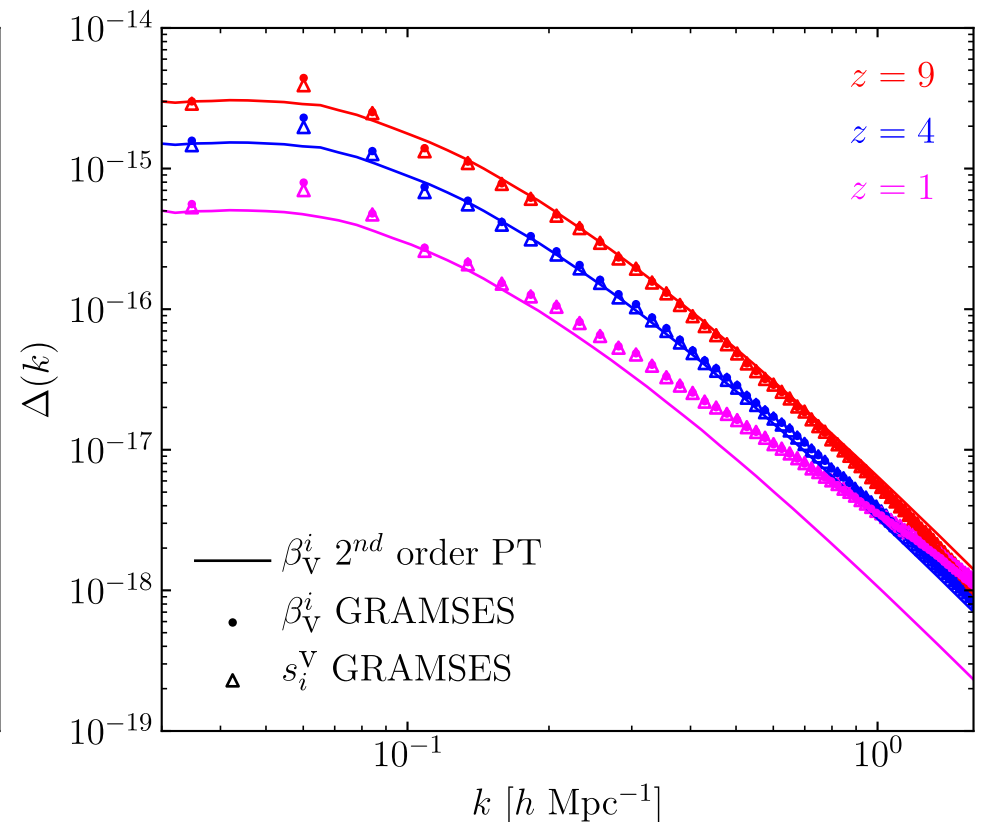
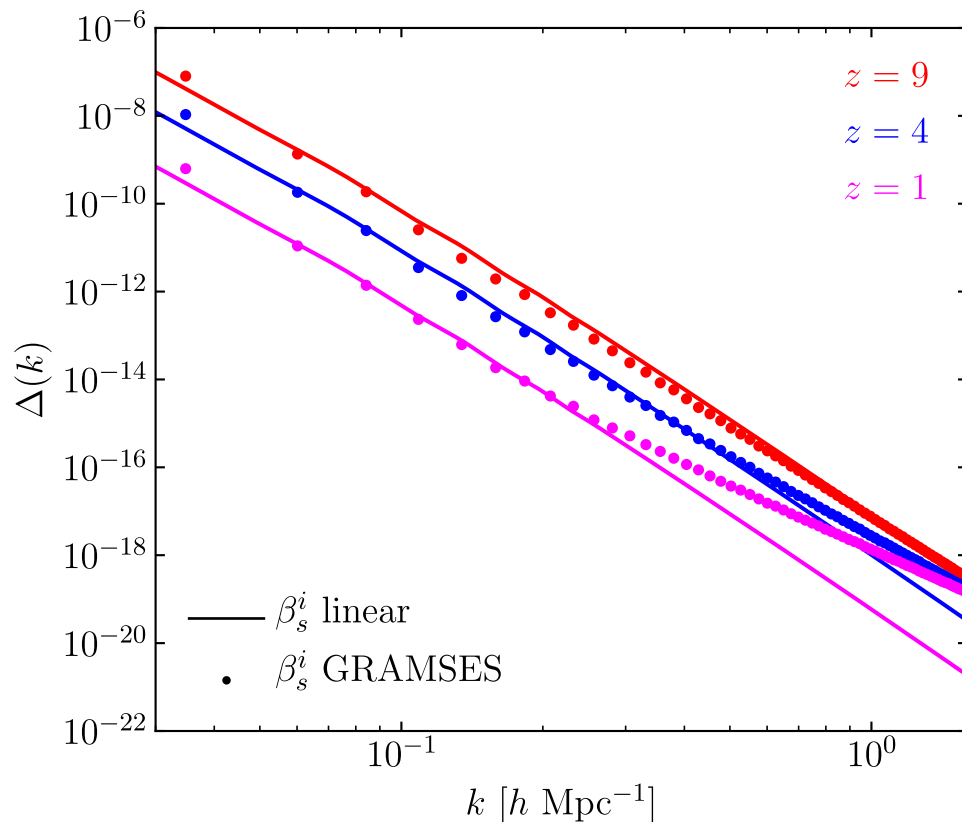
Using Wick’s theorem, the power spectrum of the vector mode is given by (Lu+ 2009):

$$\Delta_{\beta^V}(k) = \frac{9\Omega_m^2 H_0^4}{2a^2 c^2 k^2} \int_0^\infty dv \int_{|1-v|}^{1+v} du \Pi \left[\Delta_{\delta\delta}(ku) \Delta_{vv}(kv) - \frac{v}{u} \Delta_{\delta v}(ku) \Delta_{\delta v}(kv) \right]$$

Vector Potential Power Spectrum

In GRAMSES the shift vector has both scalar and vector modes

$$\beta^i = \beta_s^i + \beta_V^i$$



On the generation of Initial Conditions

- Given a density field from a linear code (CAMB, CLASS), we need to translate this into particle information.
- This can be approached as a density/displacement duality:


Positions: $x^i(q) = q^i + \Psi^i(q) : \delta(x) = -\partial_i \Psi^i$

Velocities: $v^i := \frac{dx^i}{dt} = \dot{\Psi}^i \quad \partial_i v^i = -\dot{\delta}$

- In order to generate velocities we need some information about the growth rate of density perturbations (both theory and gauge-dependent).

On the generation of Initial Conditions

- In order to avoid explicit parametrizations, we can calculate the initial velocities via finite differences of 2 neighbouring density fields:

$$v^i \approx \frac{\Delta x^i}{\Delta t} = -\frac{H}{a} \frac{\Psi^i(q_+) - \Psi^i(q_-)}{2\Delta z}$$


- We implemented this into `2LPTic` code (Crocce+ 2006).
- Since we want u_i rather than v^i we still need to correct inside GRAMSES by solving:

$$(\bar{\Delta}_L \beta_{\text{ini}})^i - 6 \frac{\Omega_m}{c^2} \beta_{\text{ini}}^i = 6 \frac{\Omega_m}{ac^2} v_{\text{ini}}^i$$

Conclusions

- We have implemented a fully non-linear **GR** N-body system on **RAMSES**.
- We cast **GR** into 10 elliptic PDEs.
- Results so far look promising.
- Due to gauge issues comparisons are difficult and the generation of IC is non-trivial.
- Upcoming/future work:
 - High resolution simulations.
 - Study back-reaction and GR effects.
 - Implement Scalar fields (dark matter models, early universe)
 - Add GW content following Cordero-Carrion 2012+.
 - Observables? (ray tracing).