

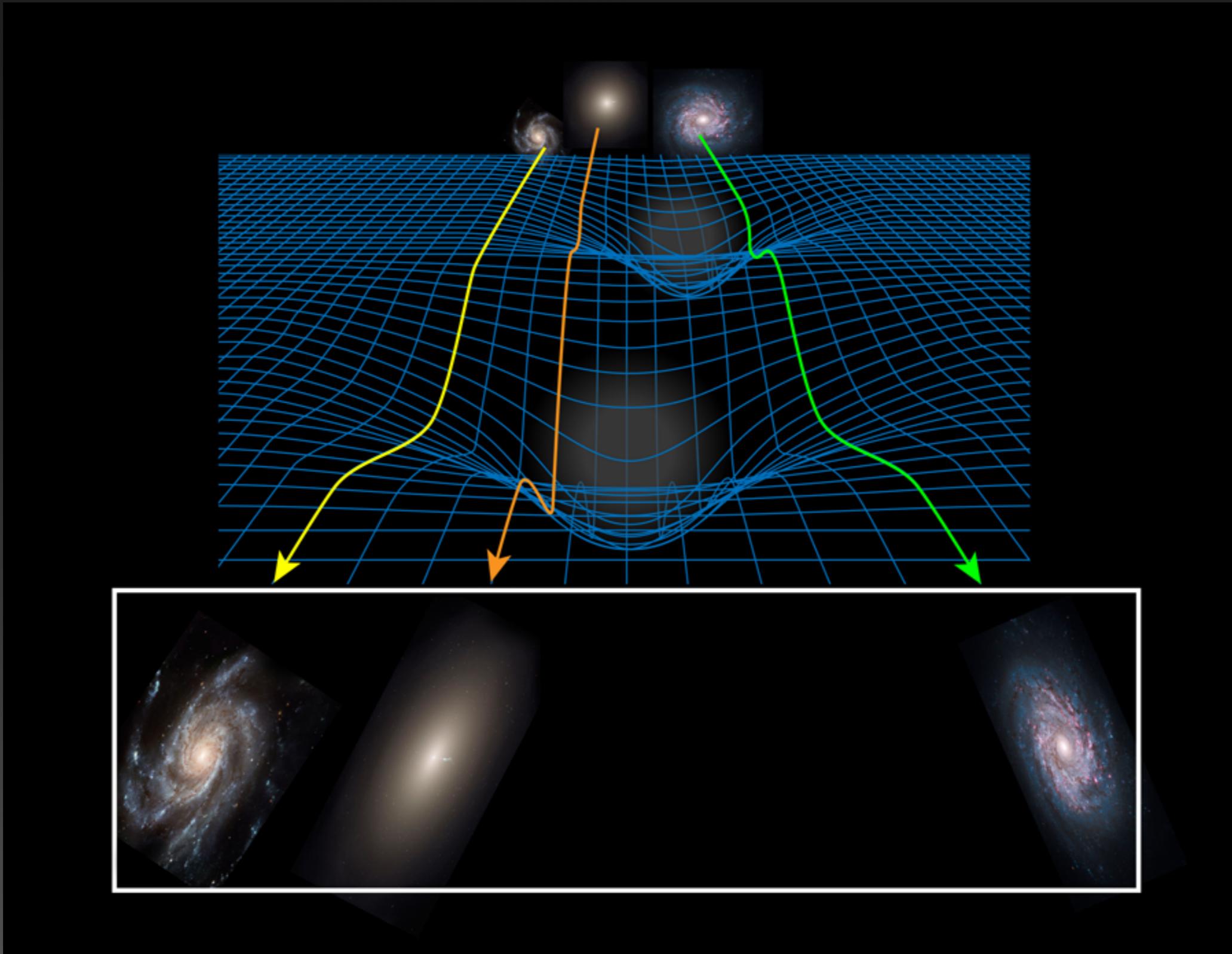
Michele Grasso
CFT Pan, Warsaw

Ray tracing, Parallax, Position and Redshift drift in Numerical Relativity

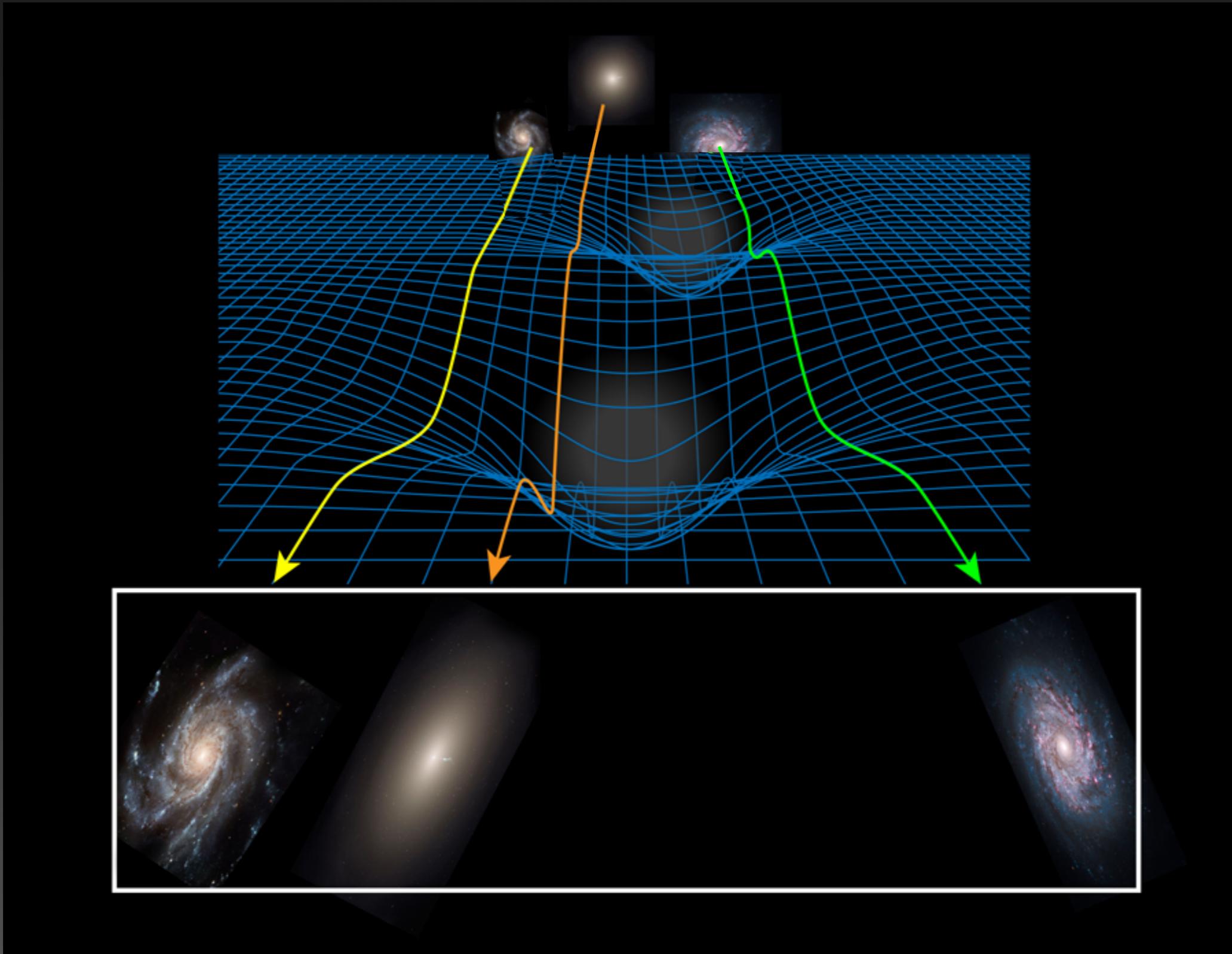
Inhomogeneous Cosmologies IV

July 18th 2019, Torun

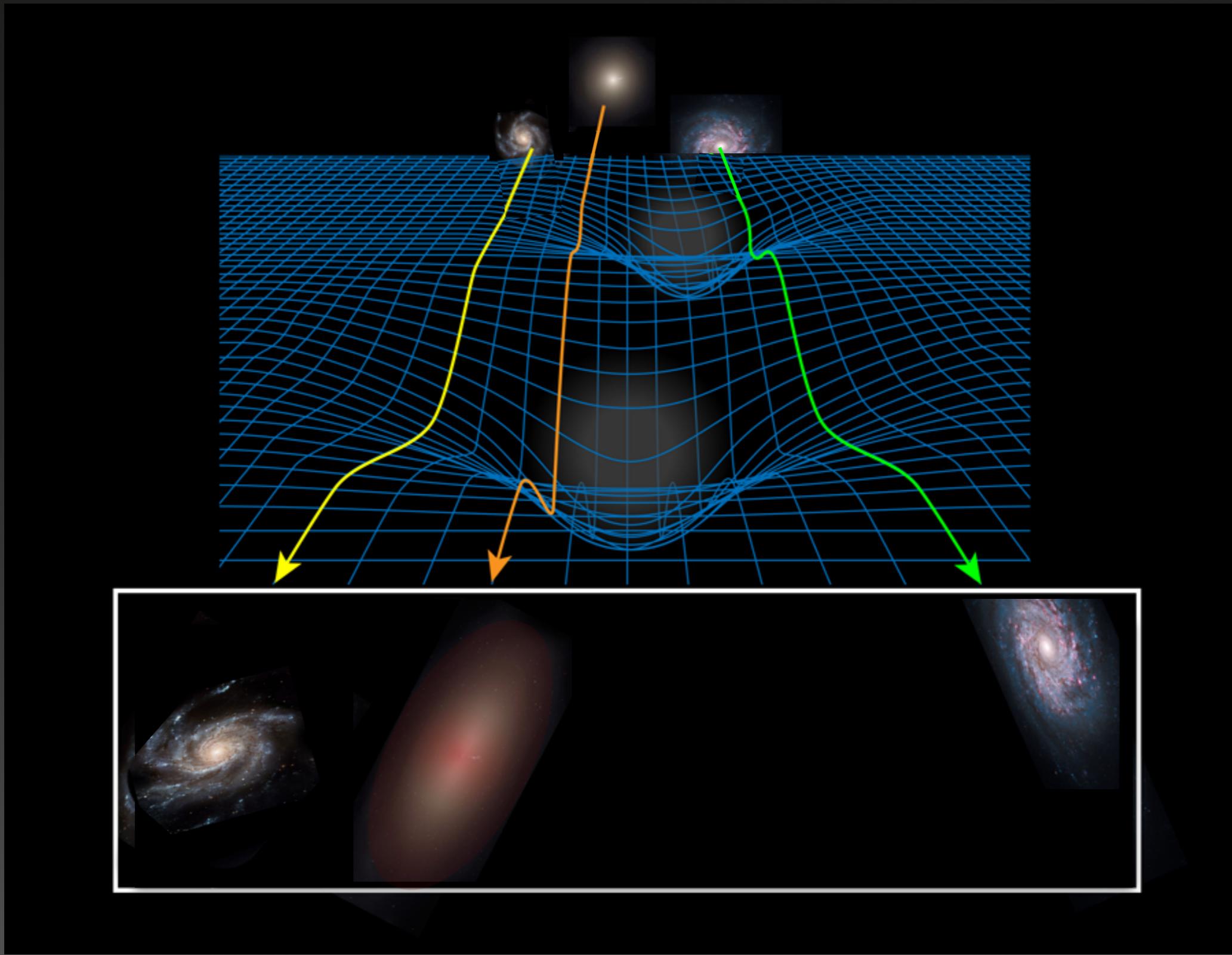
A long time ago in a galaxy far far away...



A long time ago in a galaxy far far away...

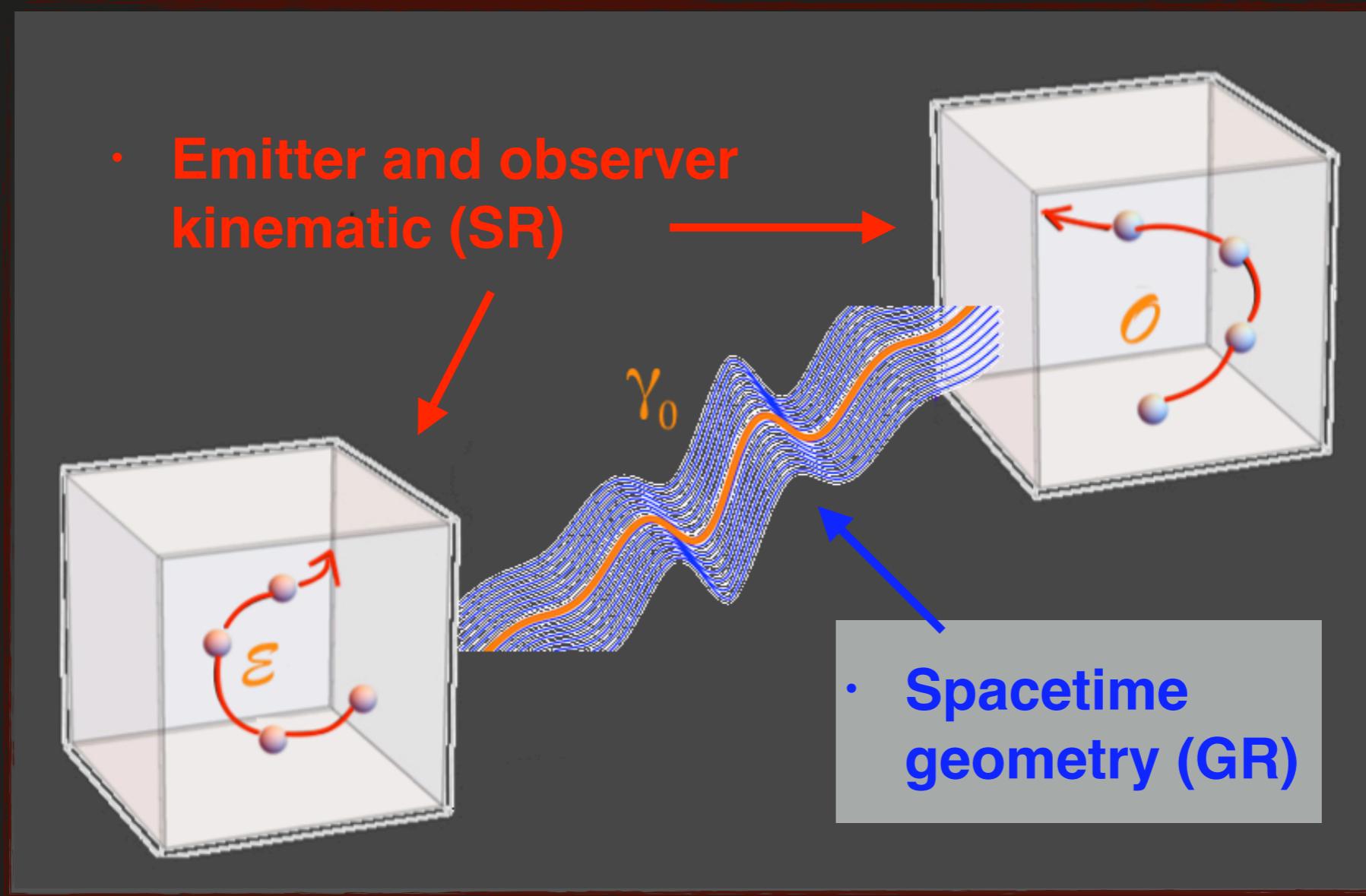


A long time ago in a galaxy far far away...



Advantage of our approach

Observables depending on:



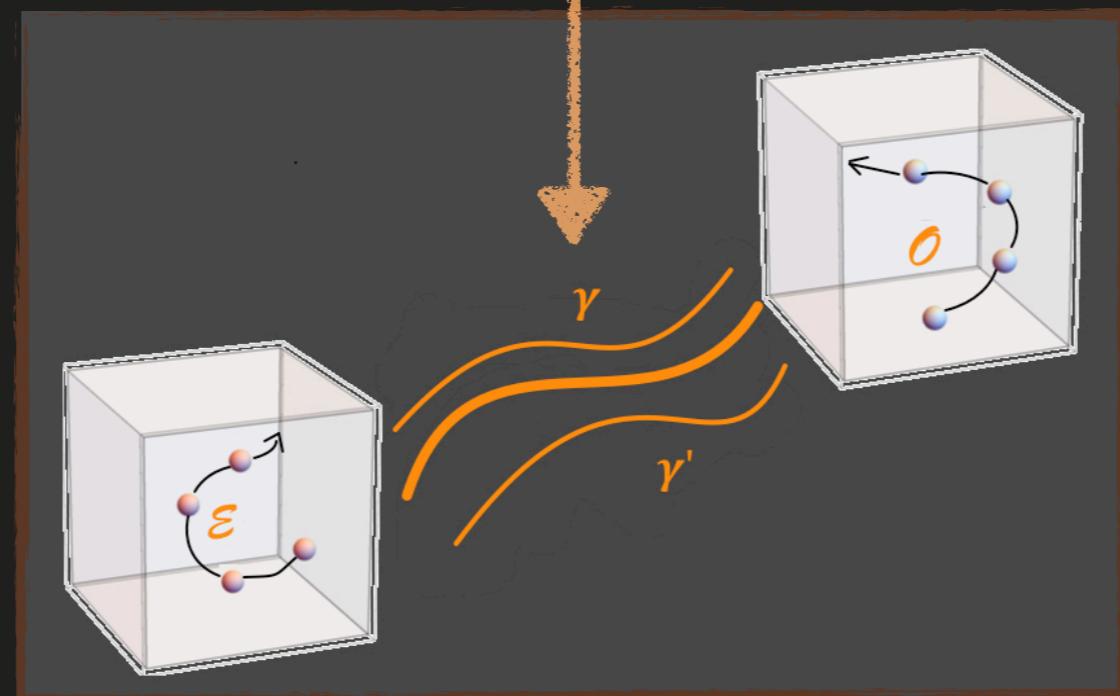
We separate the **effects of motion**
from the **effects of geometry**

MG,
Korzyński, M.
Sarbenta, J.
Phys.Rev. D, 99(6), 064038

The BGO's formalism

$$\nabla_l \nabla_l \xi^\mu - R^\mu{}_{\alpha\beta\nu} l^\alpha l^\beta \xi^\nu = 0 \quad , \quad \left\{ \begin{array}{l} \xi^\nu(\lambda_0) = \delta x_\mathcal{O}^\nu \\ \nabla_l \xi^\nu(\lambda_0) = \Delta l_\mathcal{O}^\nu \end{array} \right.$$

$$\begin{pmatrix} \delta x_\varepsilon \\ \Delta l_\varepsilon \end{pmatrix} = \mathcal{W} \begin{pmatrix} \delta x_\mathcal{O} \\ \Delta l_\mathcal{O} \end{pmatrix}$$

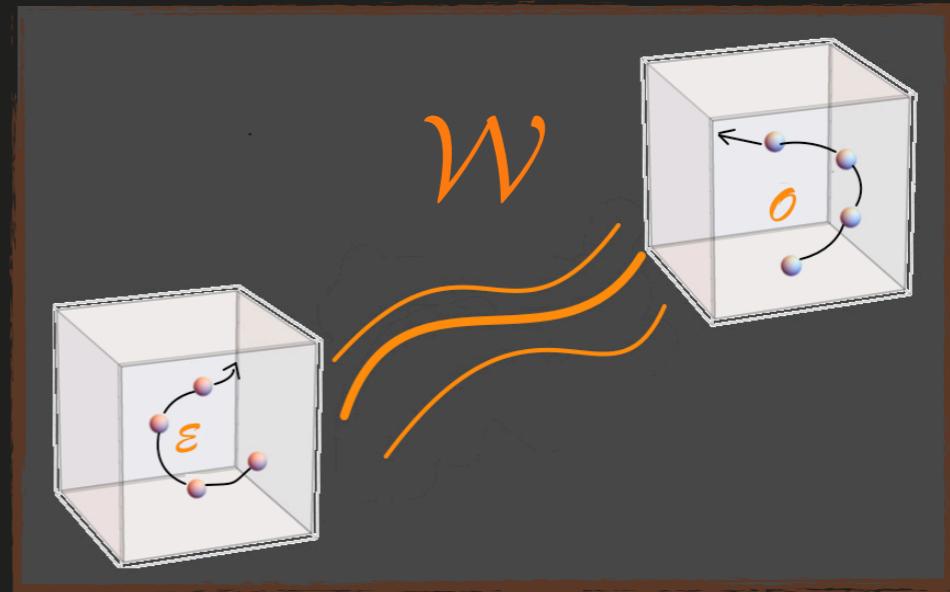


MG,
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Srebanta, J.
Phys. Rev. D, 99(6),
064038

Clear separation between curvature and E-O motion contributions in the observable expressions.

The BGO's formalism

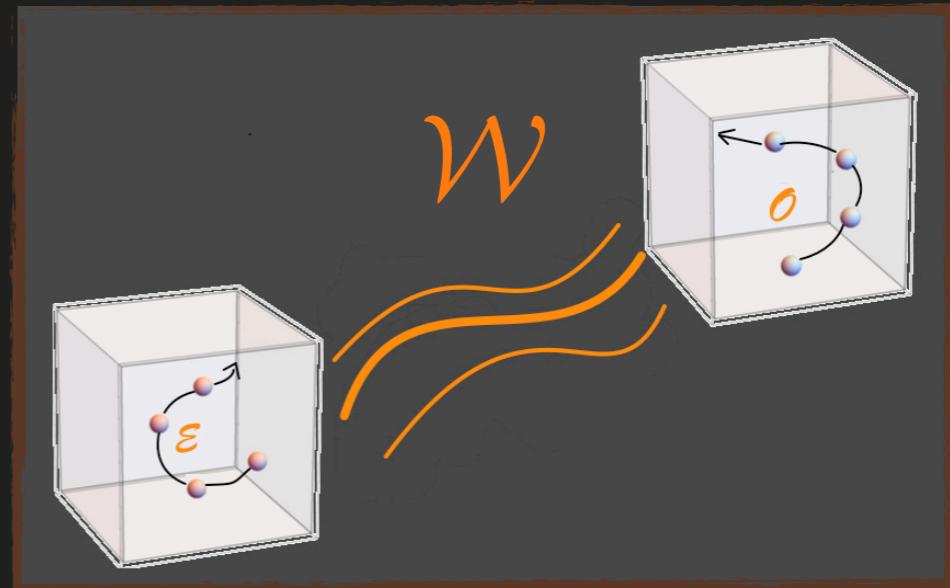
$$\begin{pmatrix} \delta x_{\mathcal{E}} \\ \Delta l_{\mathcal{E}} \end{pmatrix} = \mathcal{W} \begin{pmatrix} \delta x_{\mathcal{O}} \\ \Delta l_{\mathcal{O}} \end{pmatrix}$$



$$\mathcal{W} = \begin{pmatrix} W_{XX} & W_{XL} \\ W_{LX} & W_{LL} \end{pmatrix}$$

The BGO's formalism

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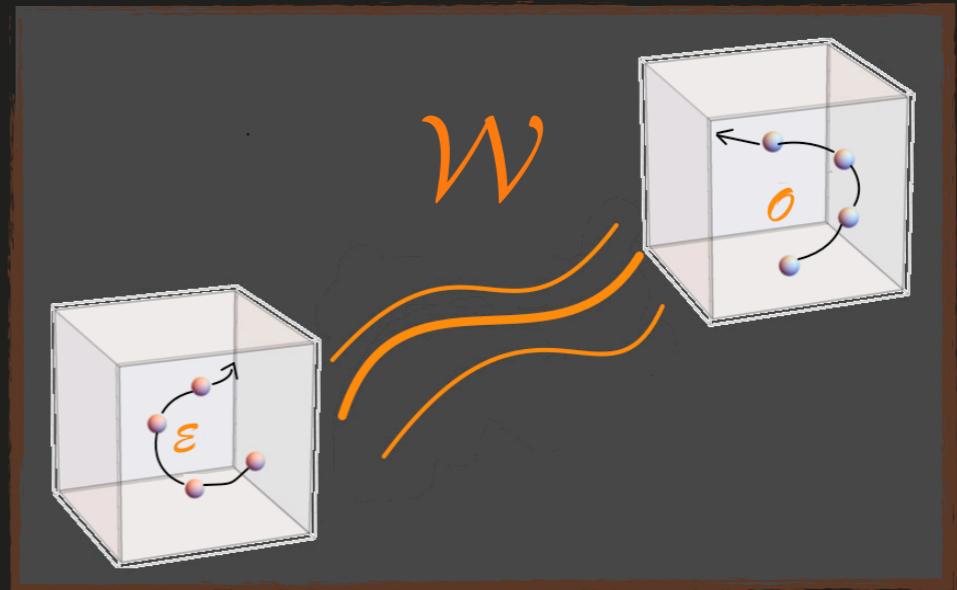


- Distortion / Magnification
- Angular Distance

$$\mathcal{W} = \begin{pmatrix} W_{XX} & \underline{W_{XL}} \\ W_{LX} & W_{LL} \end{pmatrix}$$

The BGO's formalism

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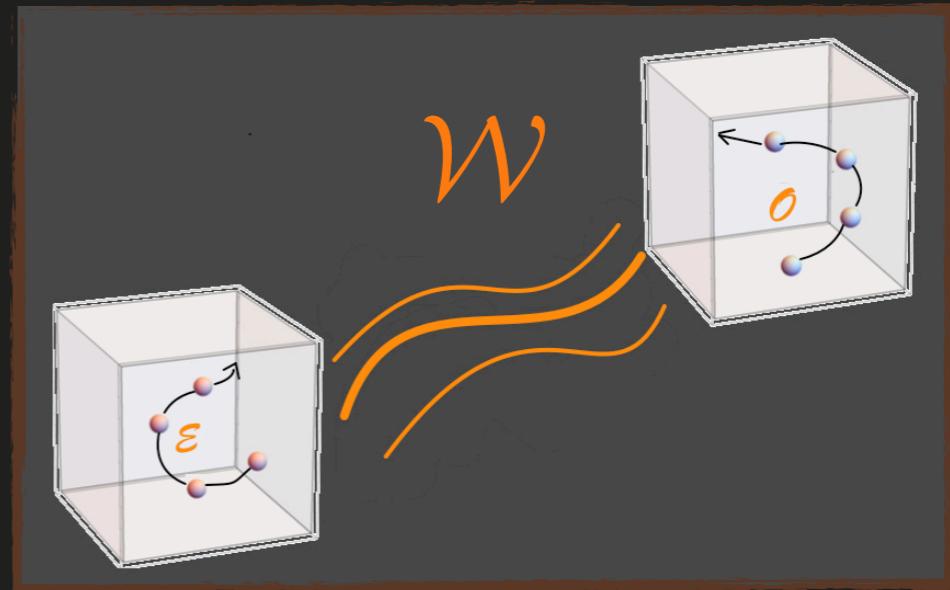


$$\mathcal{W} = \begin{pmatrix} \underline{W_{XX}} & \underline{W_{XL}} \\ W_{LX} & W_{LL} \end{pmatrix}$$

- Parallax
- Position Drift

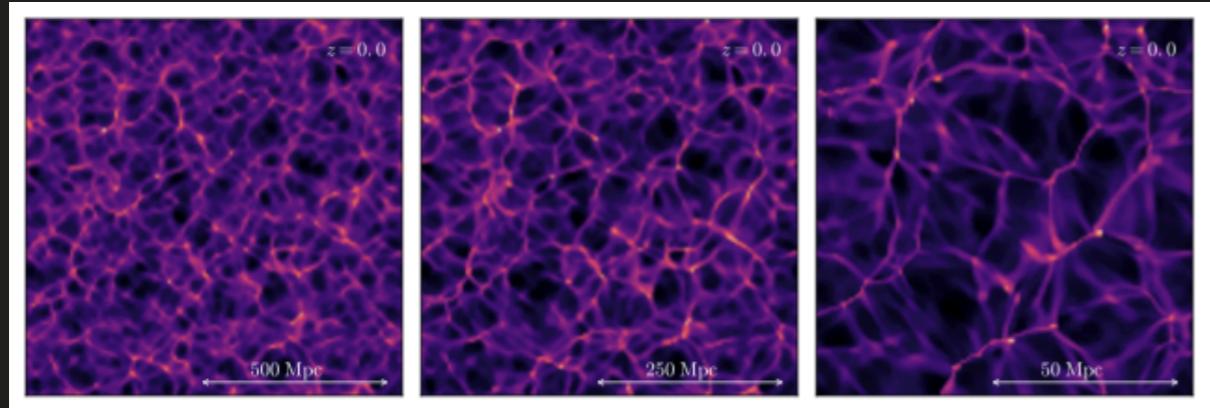
The BGO's formalism

$$\begin{pmatrix} \delta x_{\mathcal{E}} \\ \Delta l_{\mathcal{E}} \end{pmatrix} = \mathcal{W} \begin{pmatrix} \delta x_{\mathcal{O}} \\ \Delta l_{\mathcal{O}} \end{pmatrix}$$



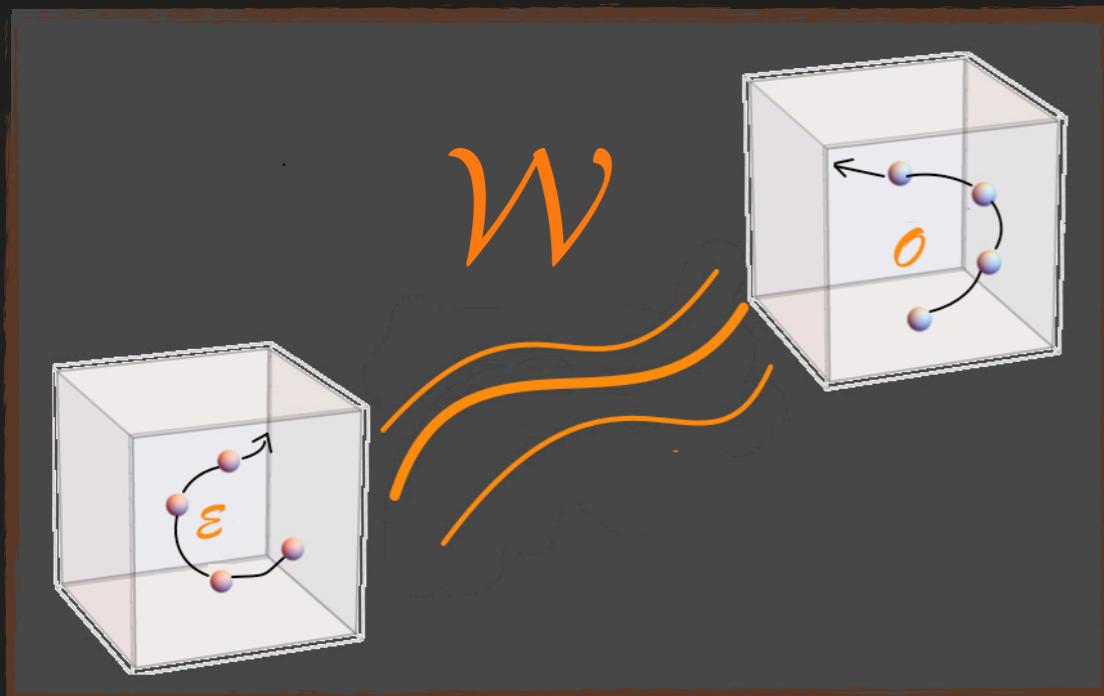
$$\mathcal{W} = \begin{pmatrix} W_{XX} & W_{XL} \\ \underline{W_{LX}} & \underline{W_{LL}} \end{pmatrix} \quad \begin{array}{l} \bullet \quad \text{Redshift Drifts} \\ \bullet \quad \text{Visual Prospective} \\ \dots \end{array}$$

Numerical Relativity and Light propagation

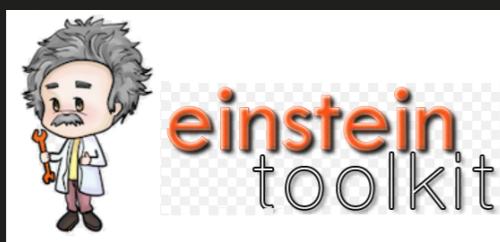


Cosmic web
from Numerical Relativity

H. Macpherson et. al., arXiv:1807.01711

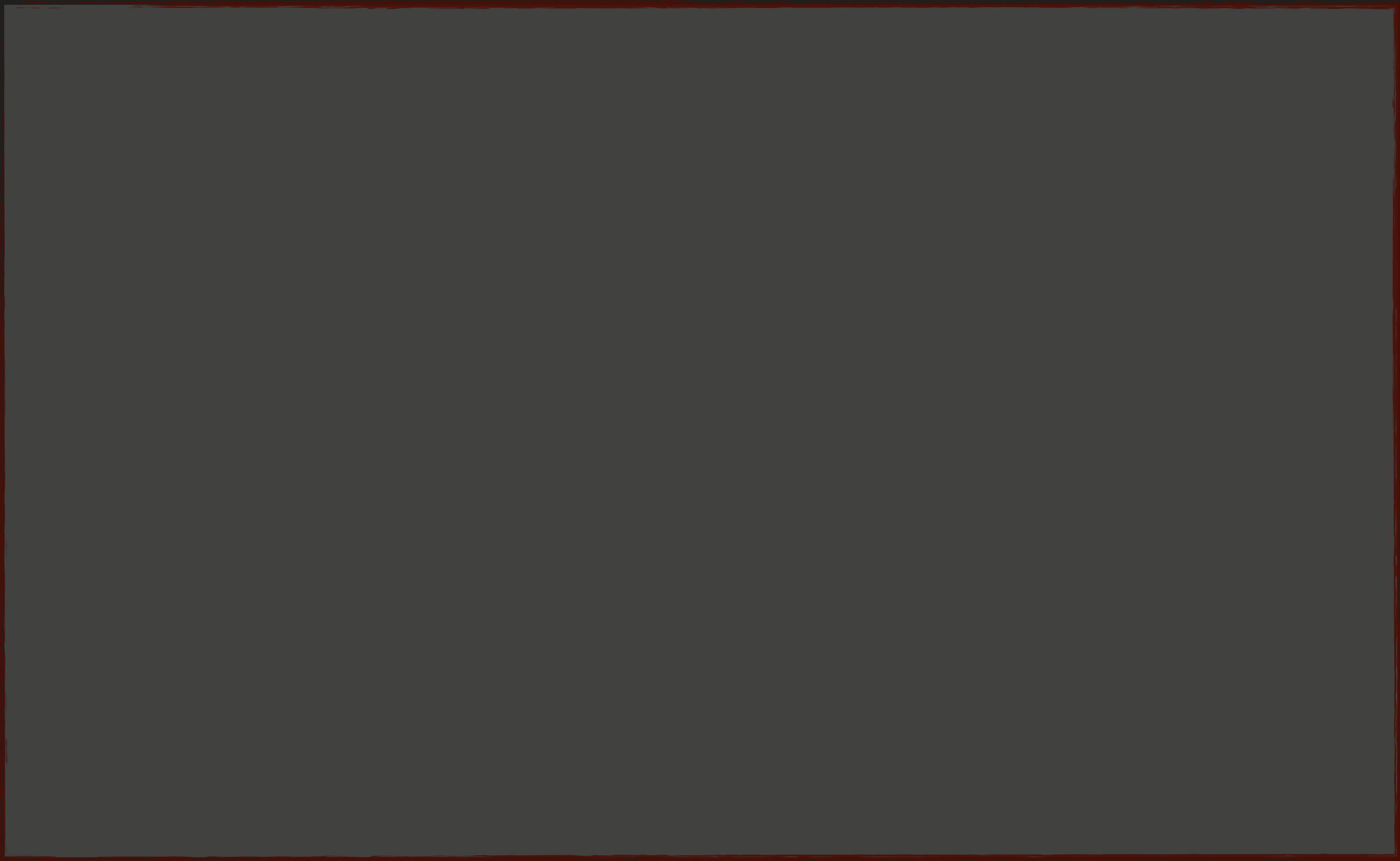


To Do



Light Propagation
in Numerical Relativity

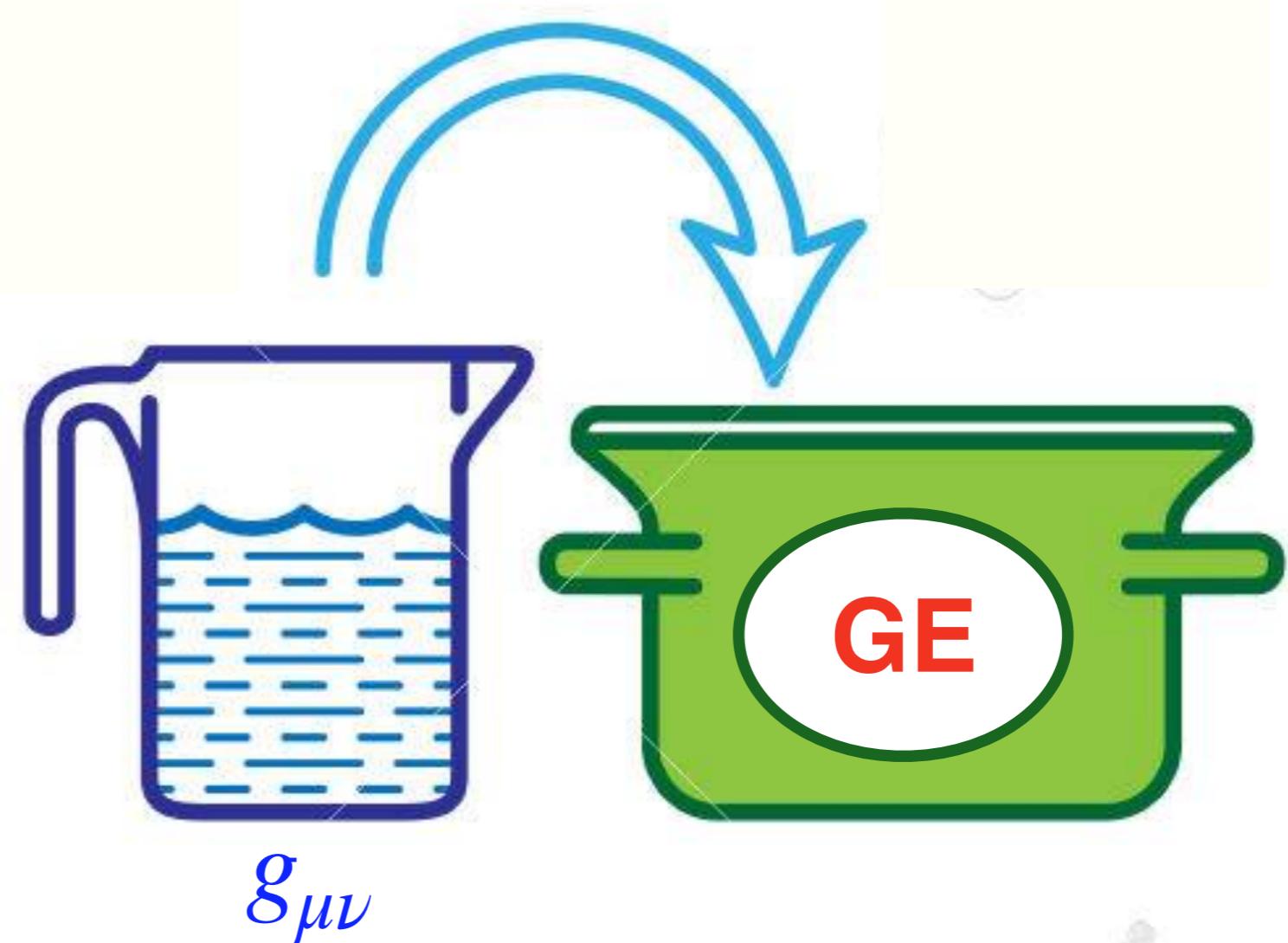
Recipe for cooking observables using BGO's



Recipe for cooking observables using BGO's

Necessary ingredient:

- Spacetime geometry



Recipe for cooking observables using BGO's

after some
time...

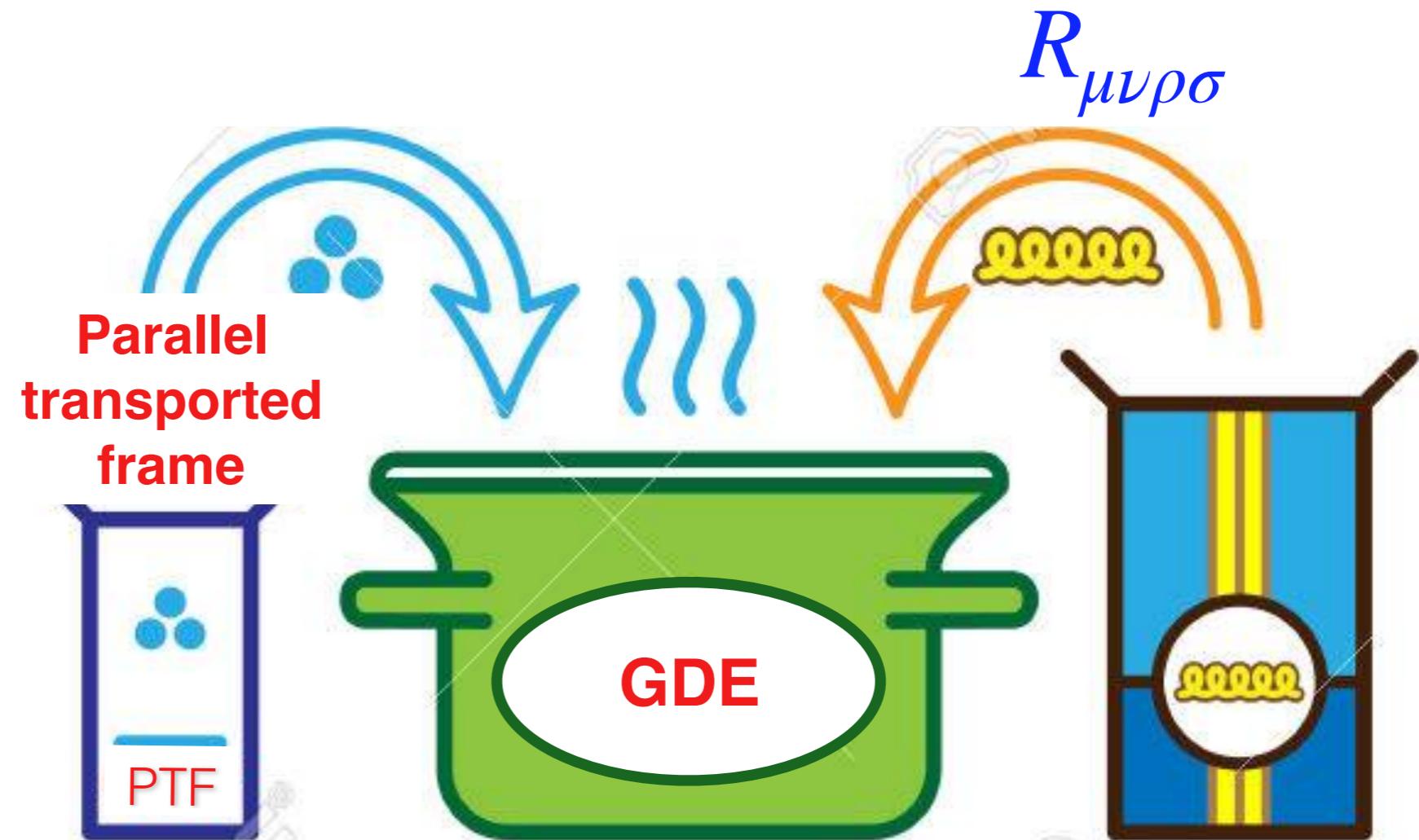
Geodesic
(Line of Sight)



Recipe for cooking observables using BGO's

Next ingredients:

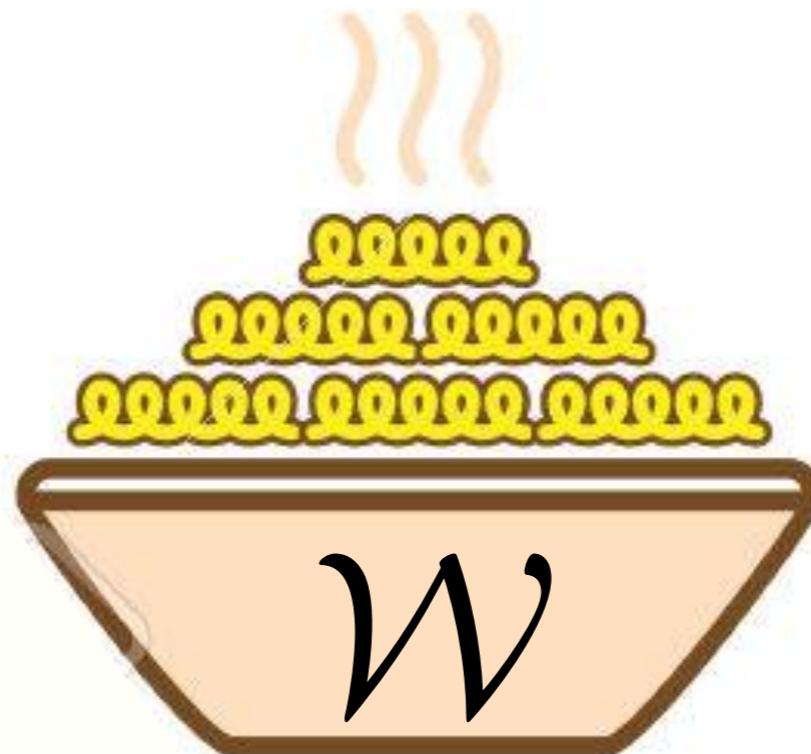
- Curvature
- a way to connect ϑ and ε



Recipe for cooking observables using BGO's

after some
time...

BGO

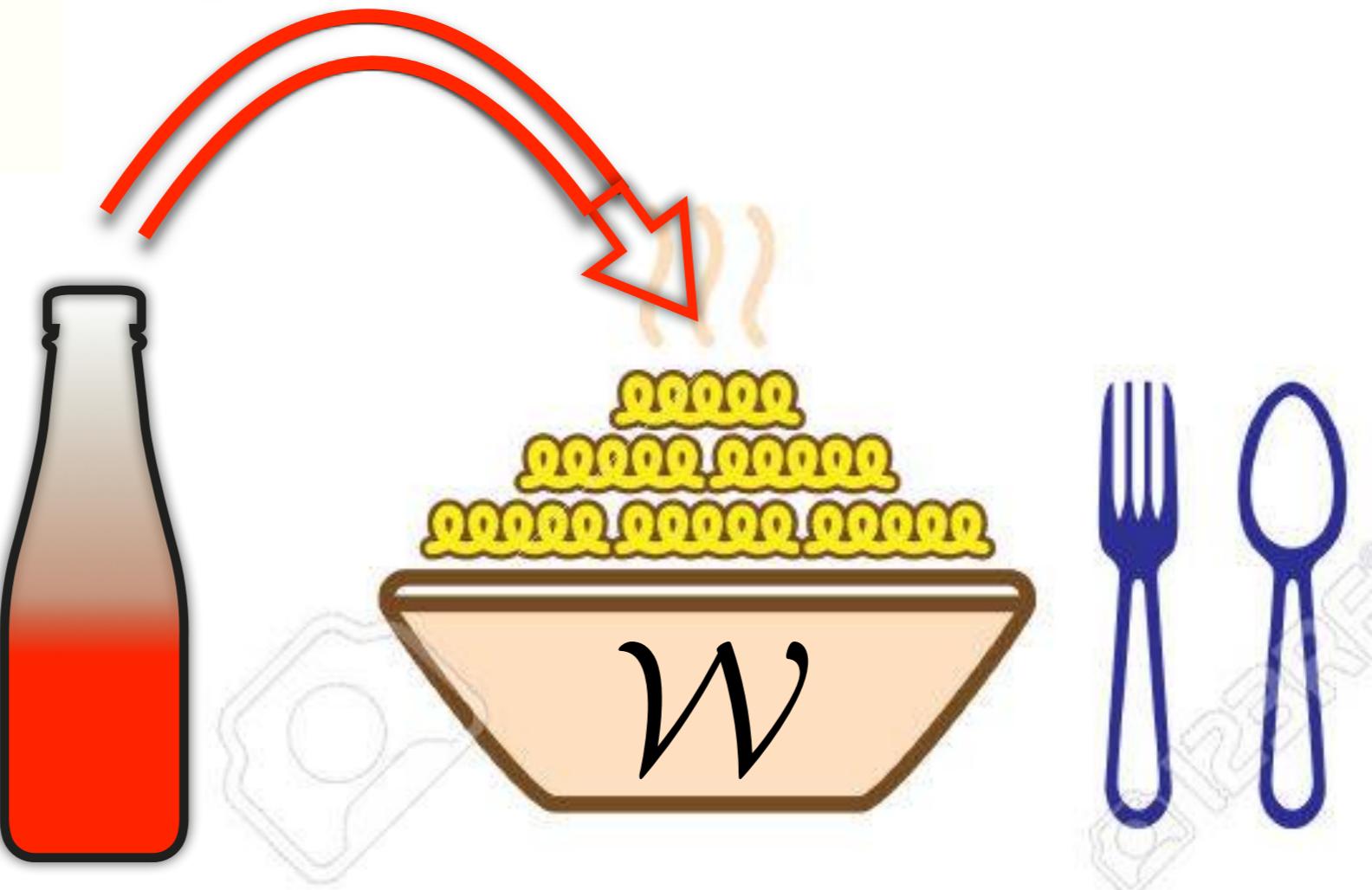


Recipe for cooking observables using BGO's

Last ingredient:

- Observer and source motion

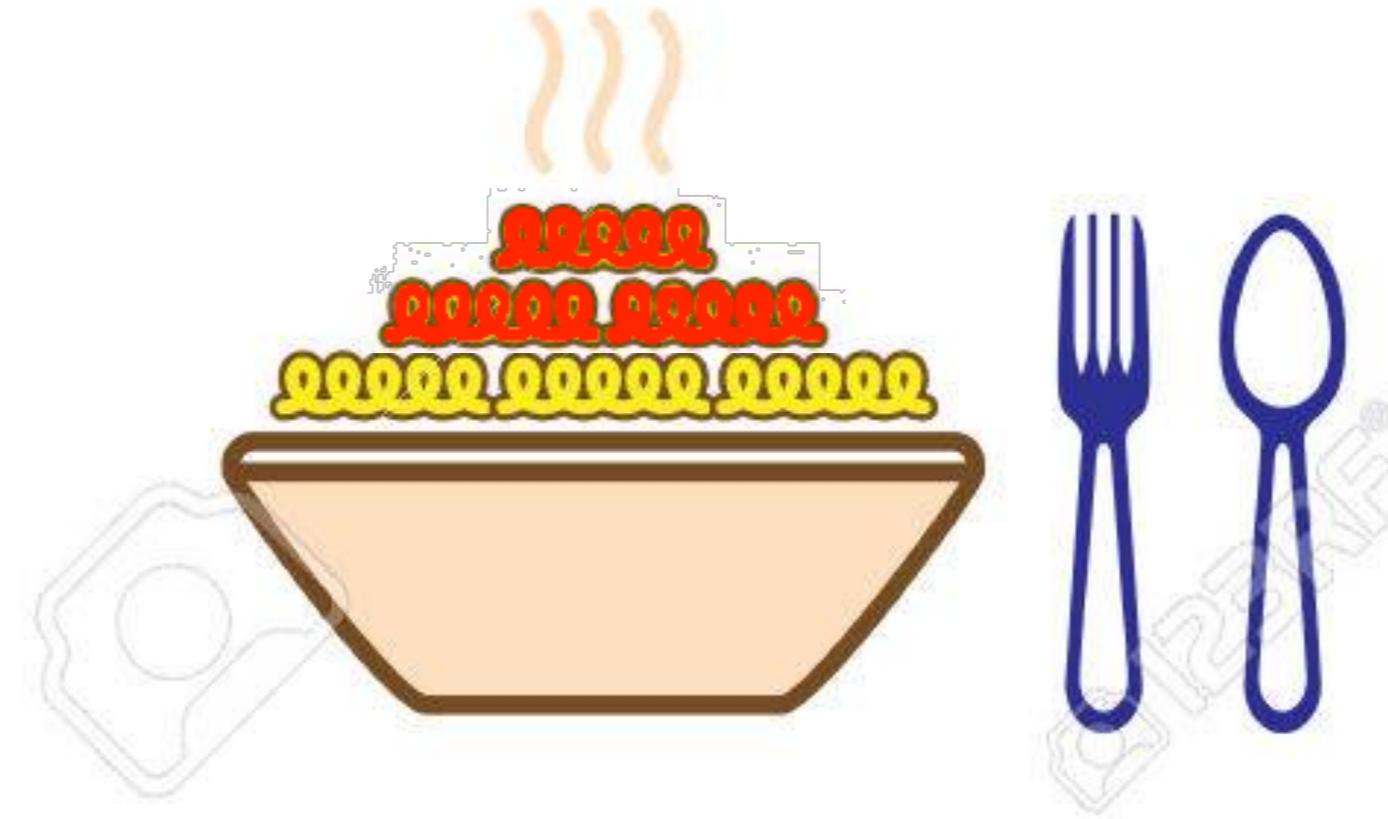
u^μ and a^μ



Recipe for cooking observables using BGO's

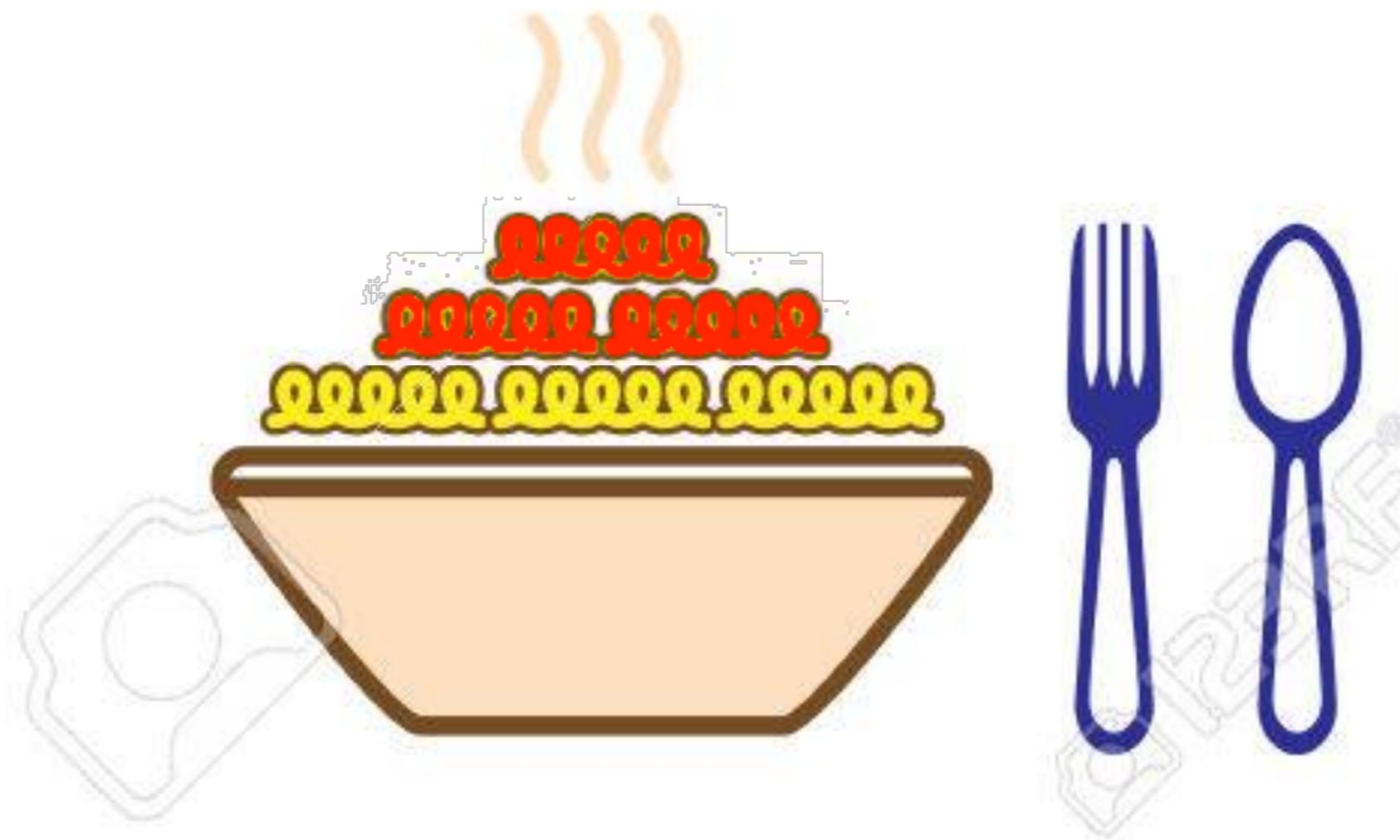
Distances
Parallax
Position drift
Redshift drift
...

Cosmological Observables

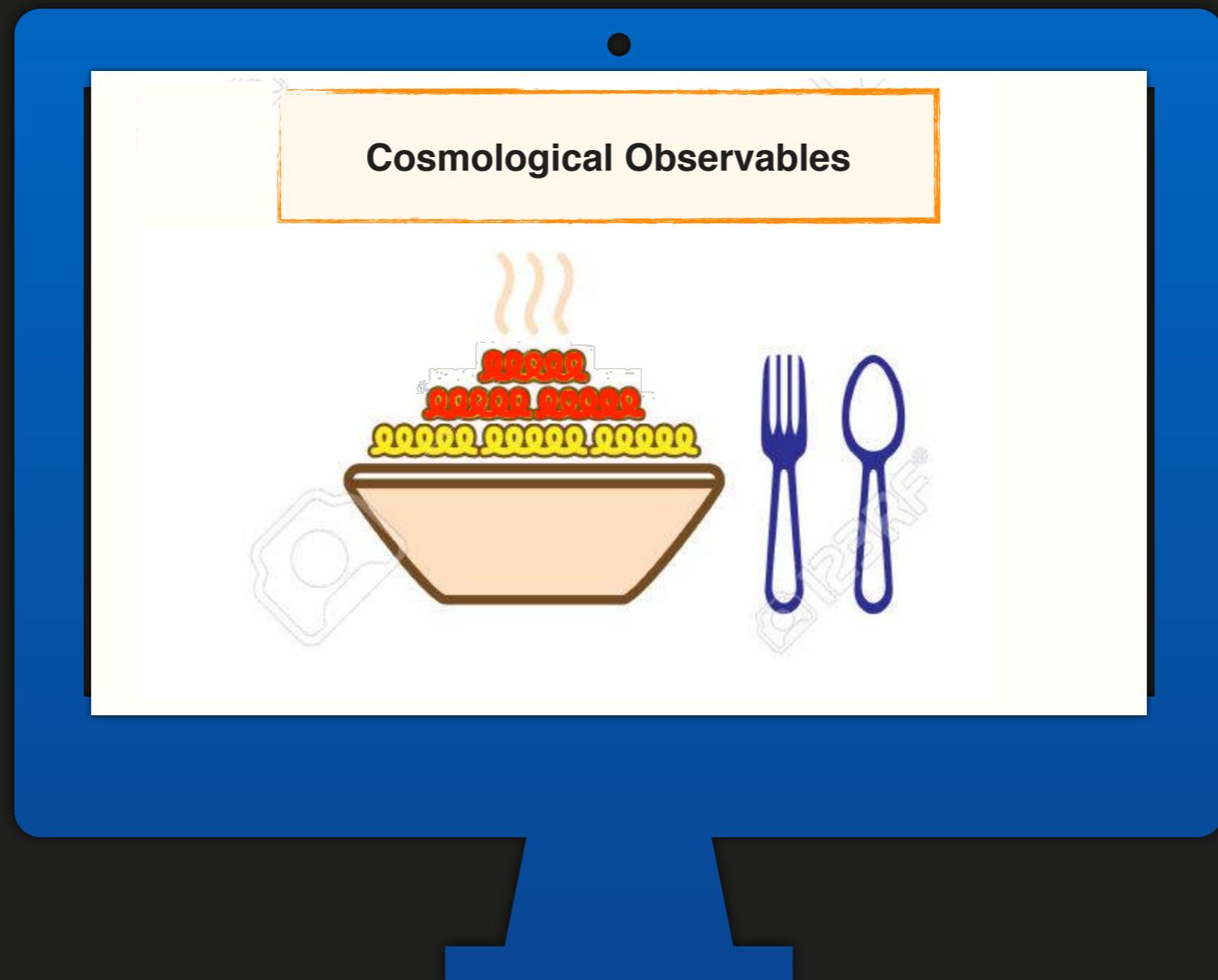


The actual computation

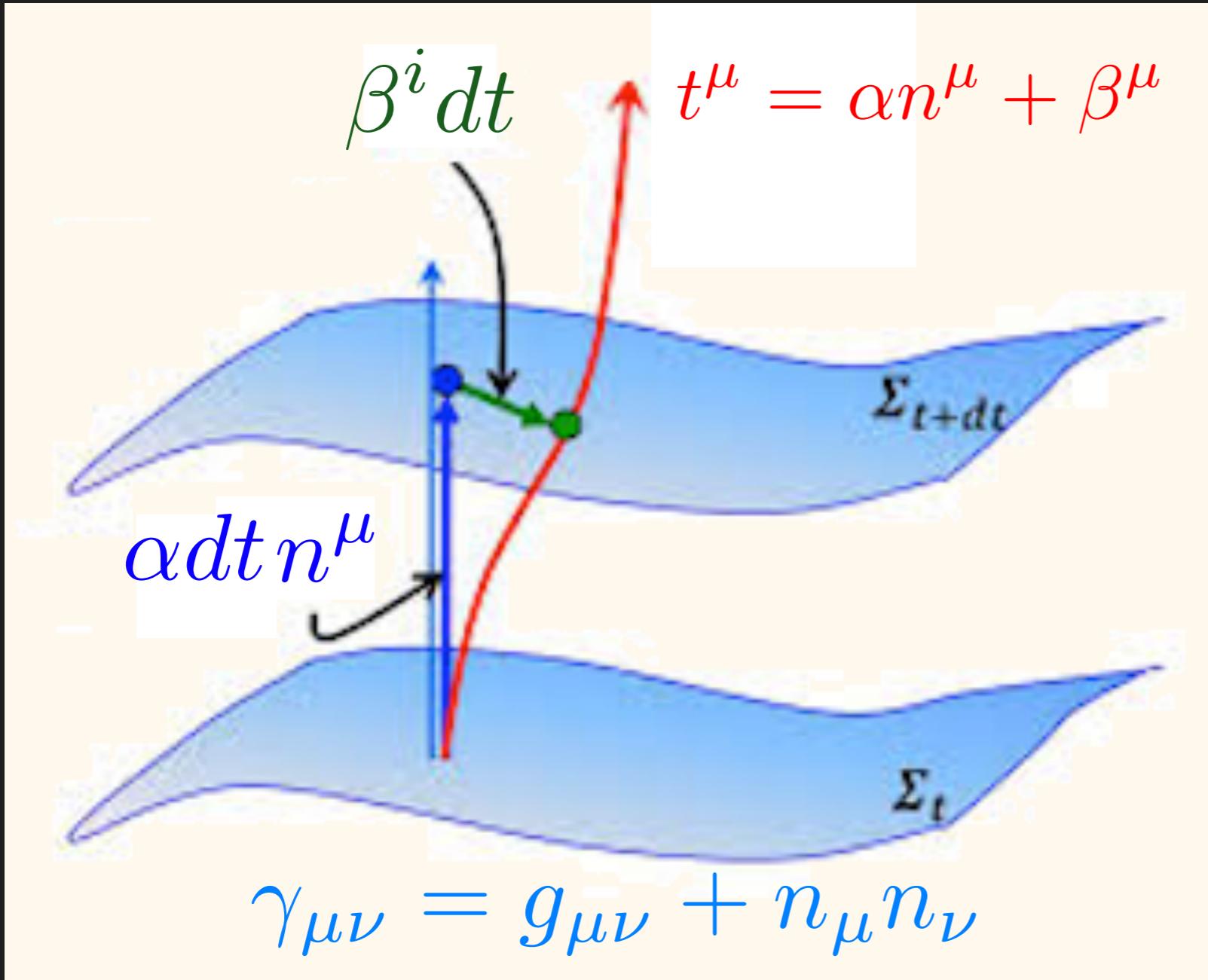
Cosmological Observables



The actual computation



3+1 Foliation of Spacetime

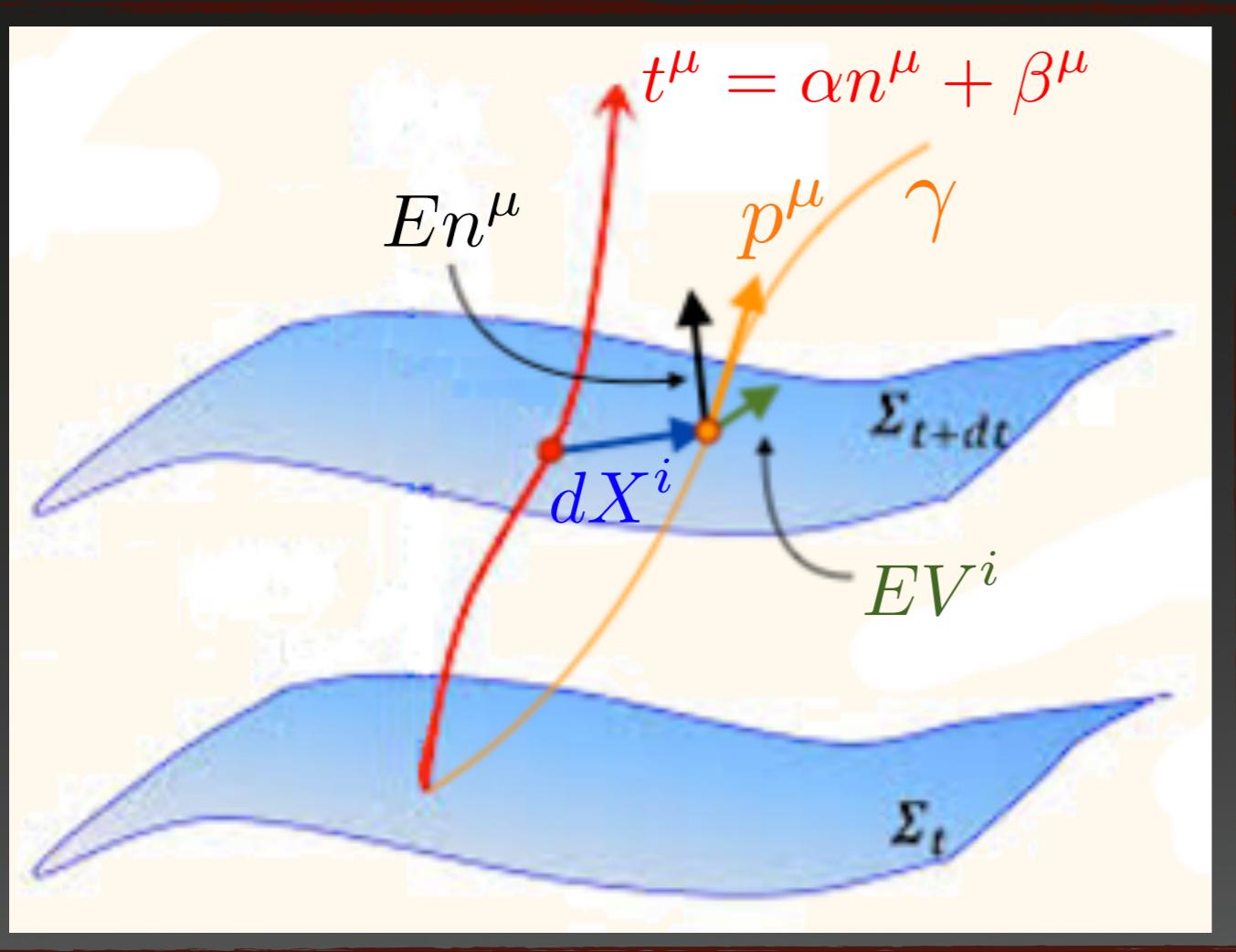


$$(\mathcal{M}, g_{\mu\nu})$$



$$(\alpha, \beta^i, \gamma_{ij}, K_{ij})$$

The boiling water: the line of sight



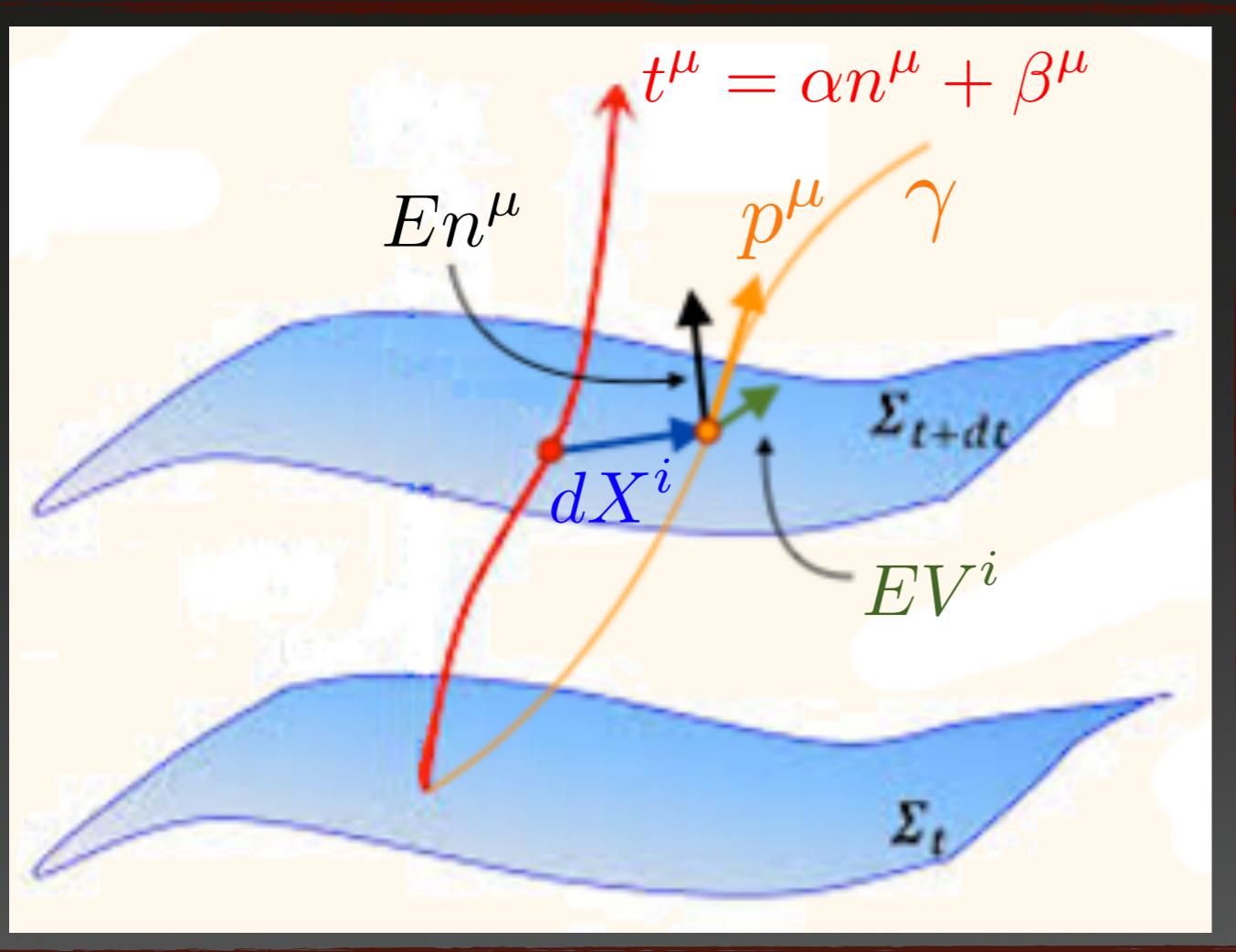
$$p^\mu = E(n^\mu + V^\mu)$$

$$p^\mu n_\mu = -E ; V^\mu n_\mu = 0$$

$$p^\mu \nabla_\mu p^\nu = 0$$

Vincent, F. H., Gourgoulhon, E., & Novak, J.
(2012). CQG 29(24), 245005.

The boiling water: the line of sight



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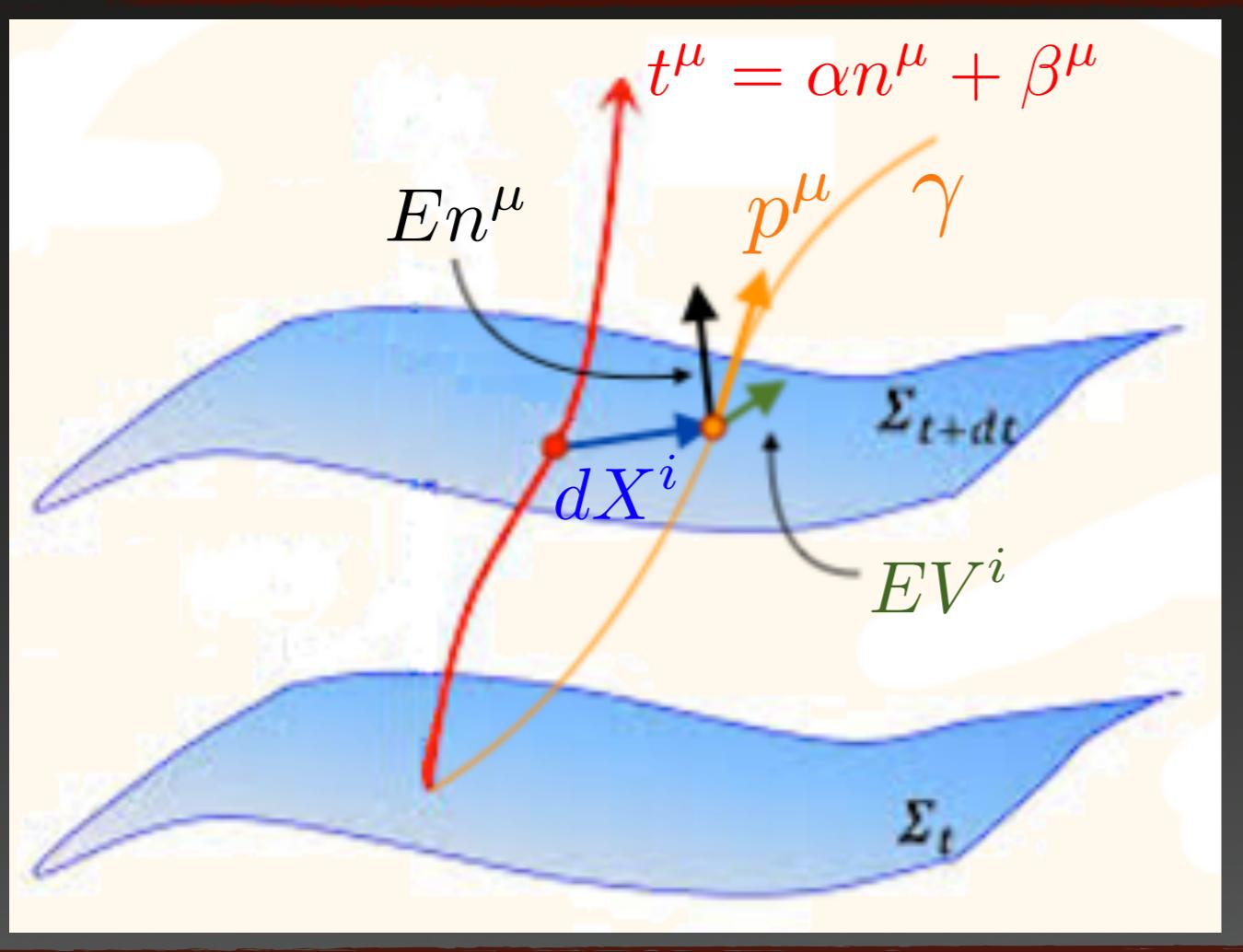
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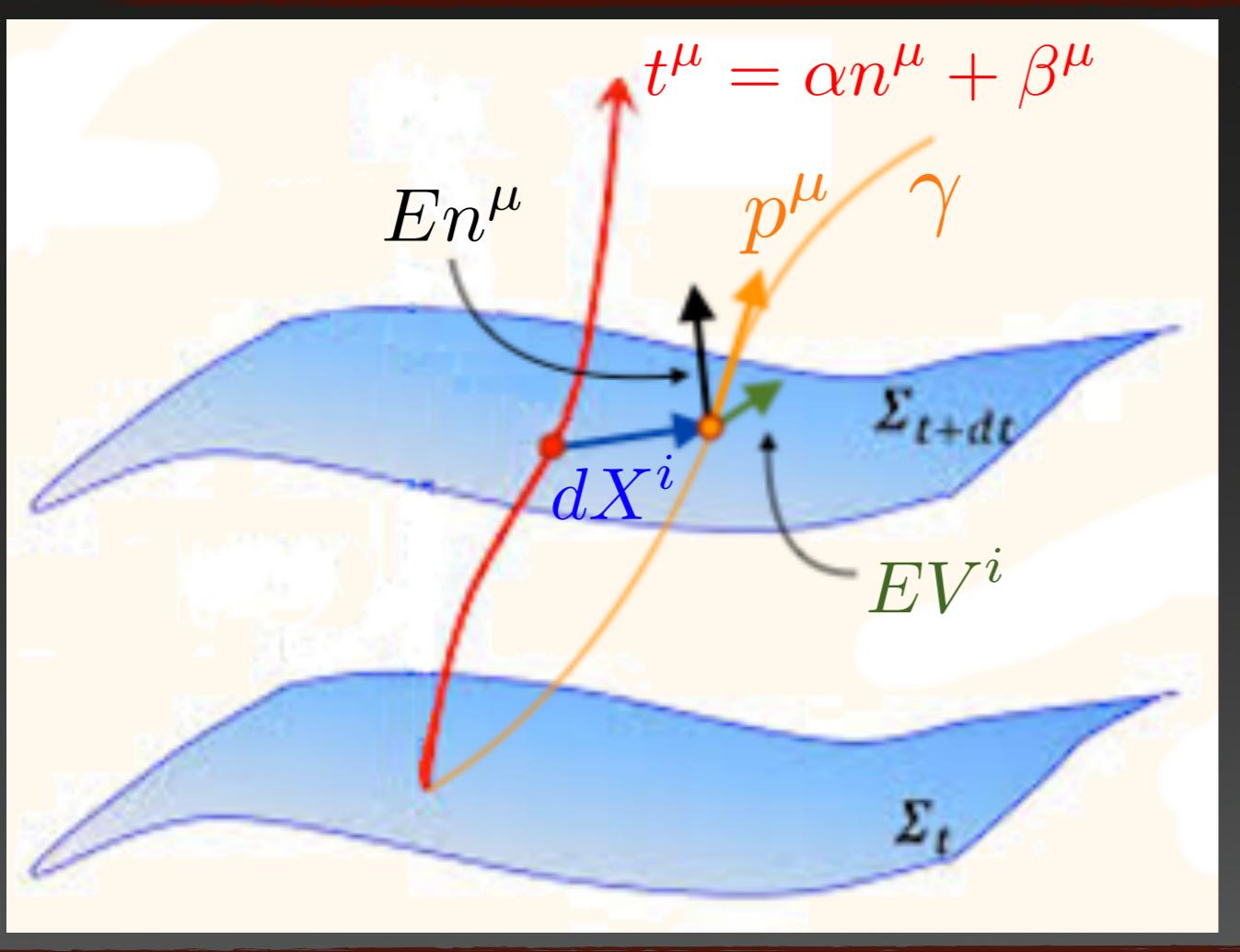
$$p^\mu \nabla_\mu p^\nu = 0$$



$$\frac{d\lambda}{dt} = \frac{\alpha}{E}$$

$$\frac{dE}{dt} = E(\alpha K_{jk} V^j V^k - V^j \partial_j \alpha)$$

The boiling water: the line of sight



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$$p^\mu = E(n^\mu + V^\mu)$$

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$$\frac{d\lambda}{dt} = \frac{\alpha}{E}$$

$$\frac{dE}{dt} = E(\alpha K_{jk} V^j V^k - V^j \partial_j \alpha)$$

$$\left\{ \begin{array}{l} \frac{dX^i}{dt} = \alpha V^i - \beta^i \\ \frac{dV^i}{dt} = \alpha V^j [V^i (\partial_i \log \alpha - K_{jk} V^k) + 2K^i{}_j - {}^{(3)}\Gamma^i_{jk} V^k] - \gamma^{ij} \partial_j \alpha - V^j \partial_j \beta^i \end{array} \right.$$

The salt: Parallel transported SNF

$$\nabla_p {}^{(4)}e^\mu = 0$$

$${}^{(4)}e^\mu = Cn^\mu + e^\mu$$

$$\alpha^{-1} \frac{dC}{dt} + e^i \partial_i \log \alpha - K_{ij} V^i e^j = 0$$

$$\alpha^{-1} \left(\frac{de^i}{dt} + e^j \partial_j \beta^i \right) + {}^{(3)}\Gamma^i{}_{jk} V^j e^k - K^i{}_j e^j + C \gamma^{ij} \partial_j \log \alpha - CK^i{}_j V^j = 0$$

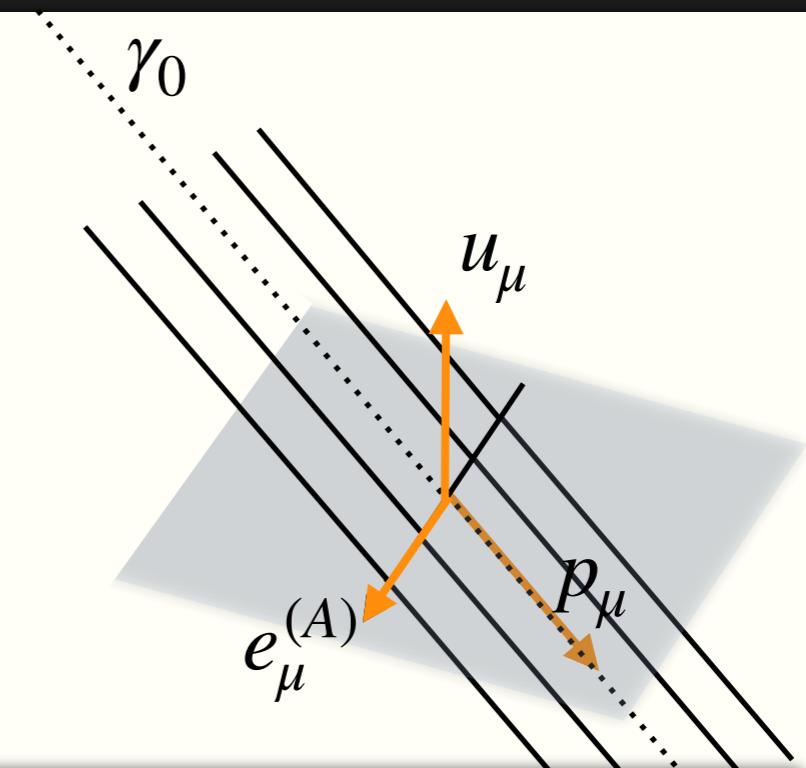
$$\mathbf{f}_{(a)} = (u^\mu, f_1^\mu, f_2^\mu, p^\mu) \quad \mathbf{e}^{(a)} = (u_\mu, e_\mu^1, e_\mu^2, p_\mu)$$

$$u^\mu u_\mu = -1 ; \; f_{(A)}^\mu e_\mu^{(B)} = \delta_{(A)}^{(B)} ; \; p^\mu p_\mu = 0 ; \; p^\mu u_\mu = Q > 0$$

GDE in parallel transported frame

GDE

$$\nabla_p \nabla_p \xi^\mu - {}^{(4)}R^\mu{}_{\alpha\beta\nu} p^\alpha p^\beta \xi^\nu = 0$$



$$\mathbf{e}^{(a)} = (u_\mu, e_\mu^1, e_\mu^2, p_\mu)$$

$$\frac{d^2 \xi^{(a)}}{dt^2} - {}^{(4)}R^{(a)}{}_{\alpha\beta(b)} p^\alpha p^\beta \xi^{(b)} + f(t) \frac{d\xi^{(a)}}{dt} = 0$$

$$\xi^{(a)} = \xi^\mu e_\mu^{(a)}$$

The raw pasta: the curvature

$$\mathcal{G}^\mu_{\alpha\beta\nu} = \gamma_\rho^\mu \gamma_\alpha^\phi \gamma_\beta^\theta \gamma_\nu^\chi \quad {}^{(4)}R^\rho_{\phi\theta\chi} = R^\mu{}_{\alpha\beta\nu} + K^\mu{}_\beta K_{\alpha\nu} - K^\mu{}_\alpha K_{\beta\nu}$$

$$\mathcal{C}^\mu{}_{\alpha\nu} = \gamma_\rho^\mu n^\phi \gamma_\alpha^\theta \gamma_\nu^\chi \quad {}^{(4)}R^\rho_{\phi\theta\chi} = D_\nu K^\mu{}_\alpha - D_\alpha K^\mu{}_\nu$$

$$\mathcal{R}^\mu{}_\nu = \gamma_\rho^\mu n^\phi \gamma_\nu^\theta n^\chi \quad {}^{(4)}R^\rho_{\phi\theta\chi} = \frac{1}{\alpha} \mathcal{L}_{\alpha n} K^\mu{}_\nu + \frac{1}{\alpha} \gamma^{\mu\sigma} D_\sigma D_\nu \alpha + K^\mu{}_\sigma K^\sigma{}_\nu$$

$${}^{(4)}R^\mu{}_{\alpha\beta\nu} = \gamma_\rho^\mu \gamma_\alpha^\delta \gamma_\beta^\eta \gamma_\nu^\sigma \quad {}^{(4)}R^\rho{}_{\delta\eta\sigma} - 2\gamma_\rho^\mu \gamma_\alpha^\delta \gamma_{[\beta}^\eta n_{\nu]} n^\sigma \quad {}^{(4)}R^\rho{}_{\delta\eta\sigma}$$

$$- 2\gamma_{\beta\rho} \gamma_\nu^\delta \gamma^{\eta[\mu} n_{\alpha]} n^\sigma \quad {}^{(4)}R^\rho{}_{\delta\eta\sigma} + 2\gamma_\rho^\mu \gamma_{[\beta}^\eta n_{\nu]} n_\alpha n^\delta n^\sigma \quad {}^{(4)}R^\rho{}_{\delta\eta\sigma}$$

$$- 2\gamma_{\alpha\rho} \gamma_{[\beta}^\eta n_{\nu]} n^\mu n^\delta n^\sigma \quad {}^{(4)}R^\rho{}_{\delta\eta\sigma}$$

APPLICATIONS: Schwarzschild metric as toy model

Initial conditions:

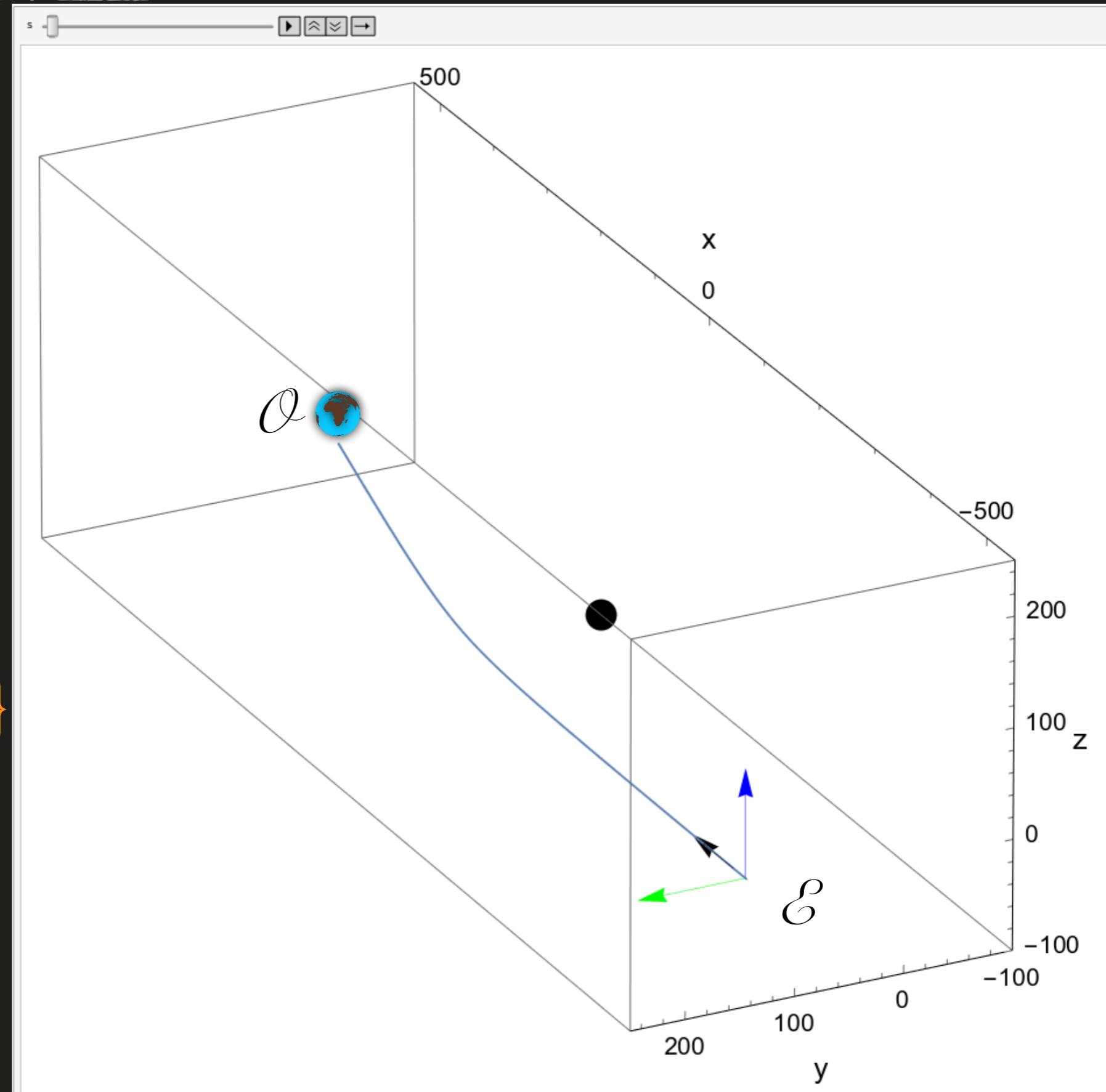
(i) $x_{in} = 40 r_S$

(ii) $y_{in} = 10 r_S$

(iii) $z_{in} = 0$

(iv) $k_{in}^\mu = \{1, 1, 0, 0\}$

(v) Static
emitter and
observer



APPLICATIONS: Schwarzschild metric as toy model

Initial conditions:

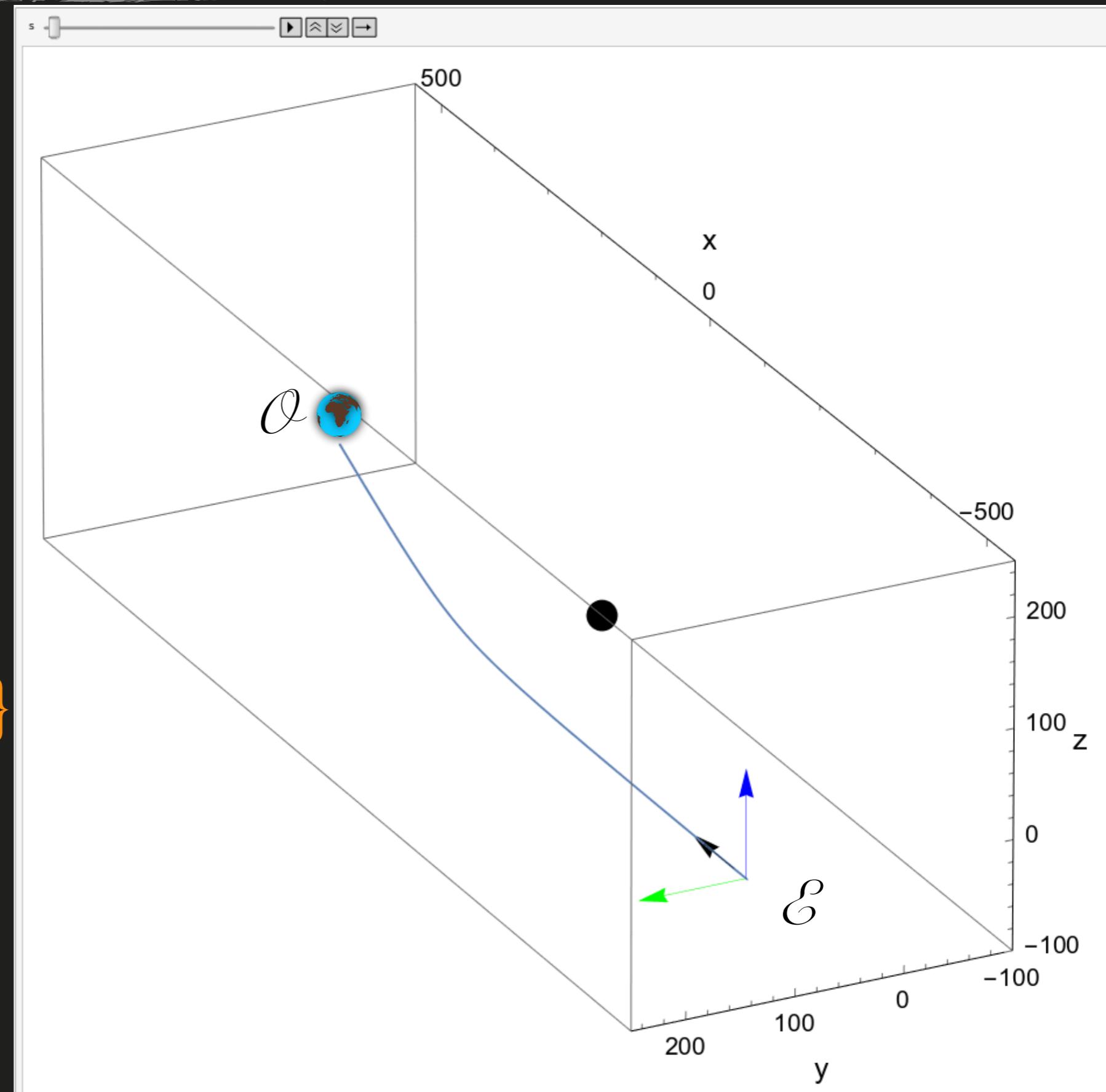
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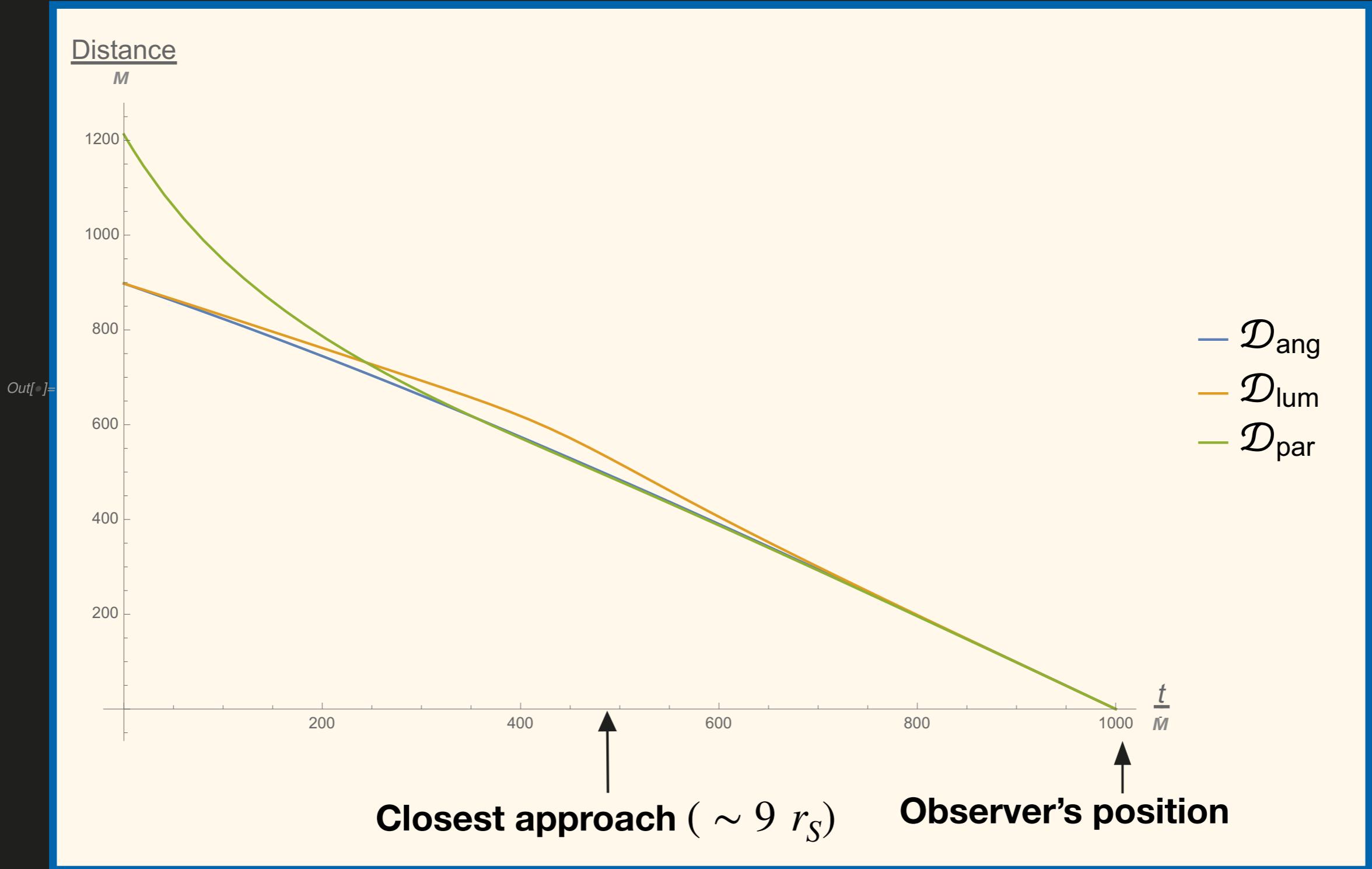
(iv) $k_{in}^\mu = \{1, 1, 0, 0\}$

(v) Static
emitter and
observer



From the BGO we get \mathcal{D}_{ang} and \mathcal{D}_{par}

From the BGO + redshift we get \mathcal{D}_{lum}



APPLICATIONS: isolating non linearities in light propagation in inhomogeneous cosmology

Our toy-model: plane-parallel dynamics in the PN approximation

- Perturbations depend on x and time only
- Leading order in the $1/c^{2n}$ expansion of Einstein equations

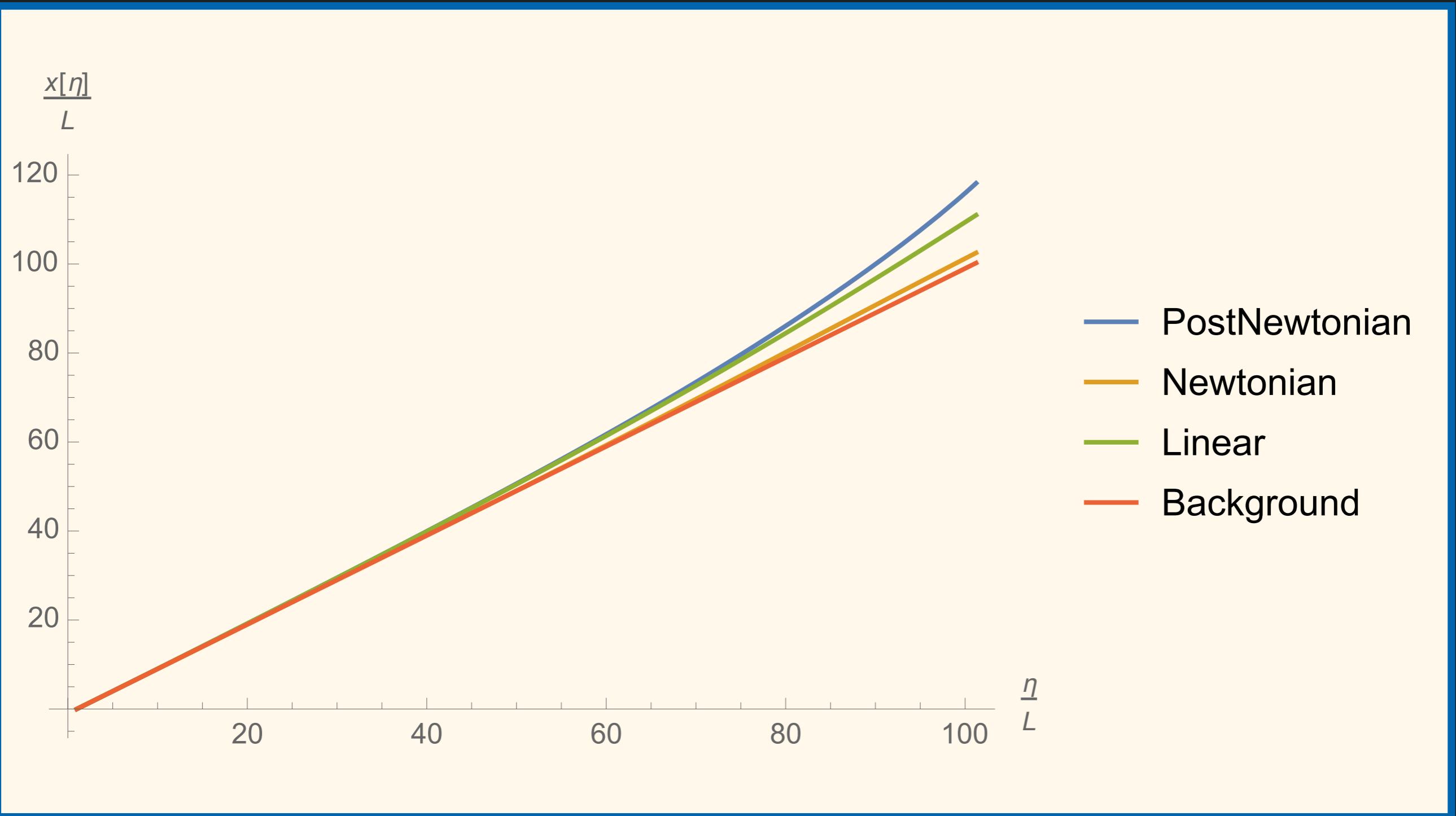
$$ds^2 = a^2(\eta) \left\{ -d\eta^2 + [\gamma_{11}^{\text{Nwt}}(\eta, x) + \gamma_{11}^{\text{PN}}(\eta, x)] dx^2 + [1 + \gamma_{22}^{\text{PN}}(\eta, x)] (dy^2 + dz^2) \right\}$$

APPLICATIONS: isolating non linearities in light propagation in inhomogeneous cosmology

Comparison of cosmological observables in three cases:

- (I) **Newtonian** dynamics + **exact** relativistic light propagation
- (II) **PN** dynamics + **exact** relativistic light propagation
- (III) **Standard PT at first order** for both

APPLICATIONS: isolating non linearities in light propagation in inhomogeneous cosmology



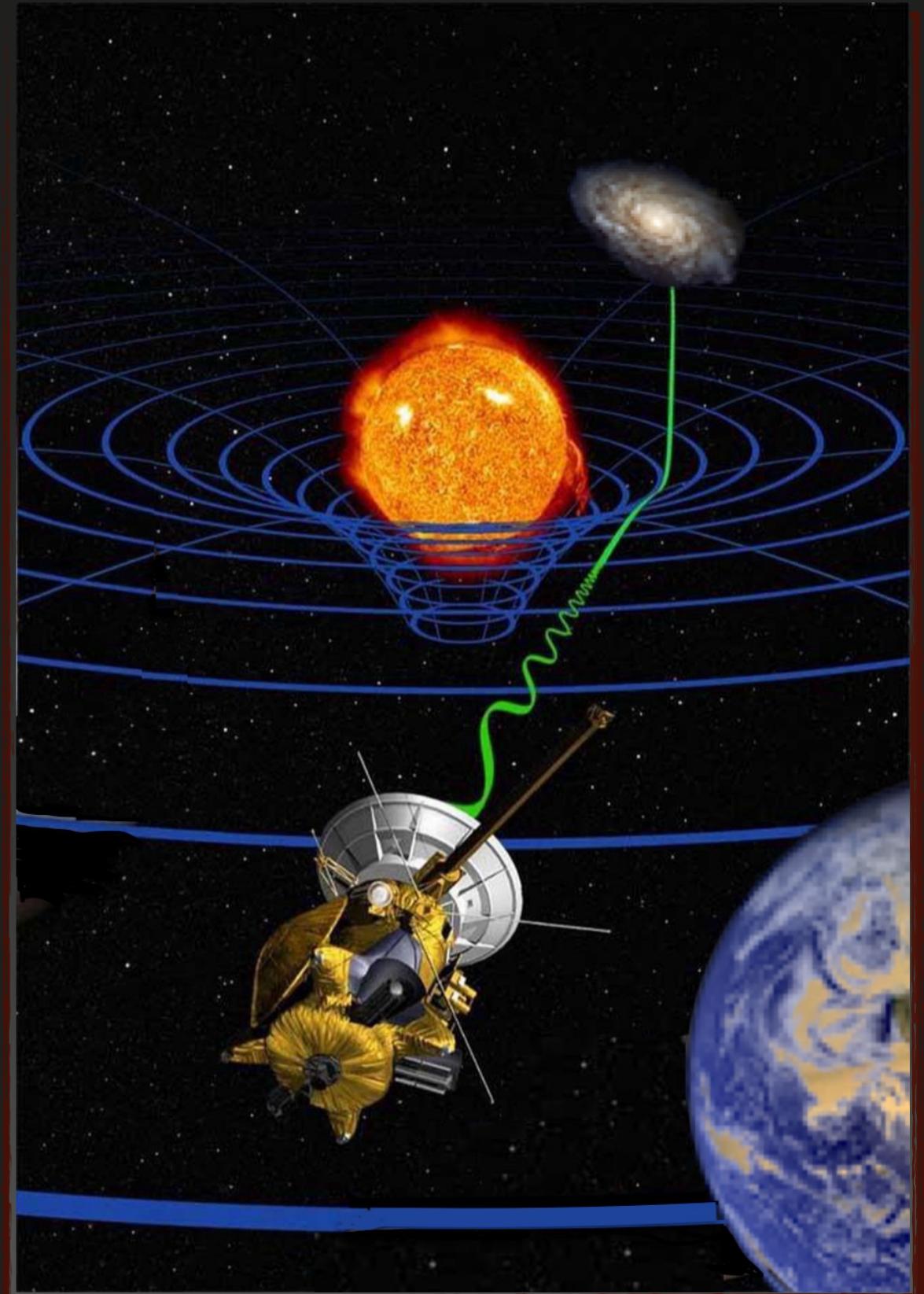
Summary

Take-home message:

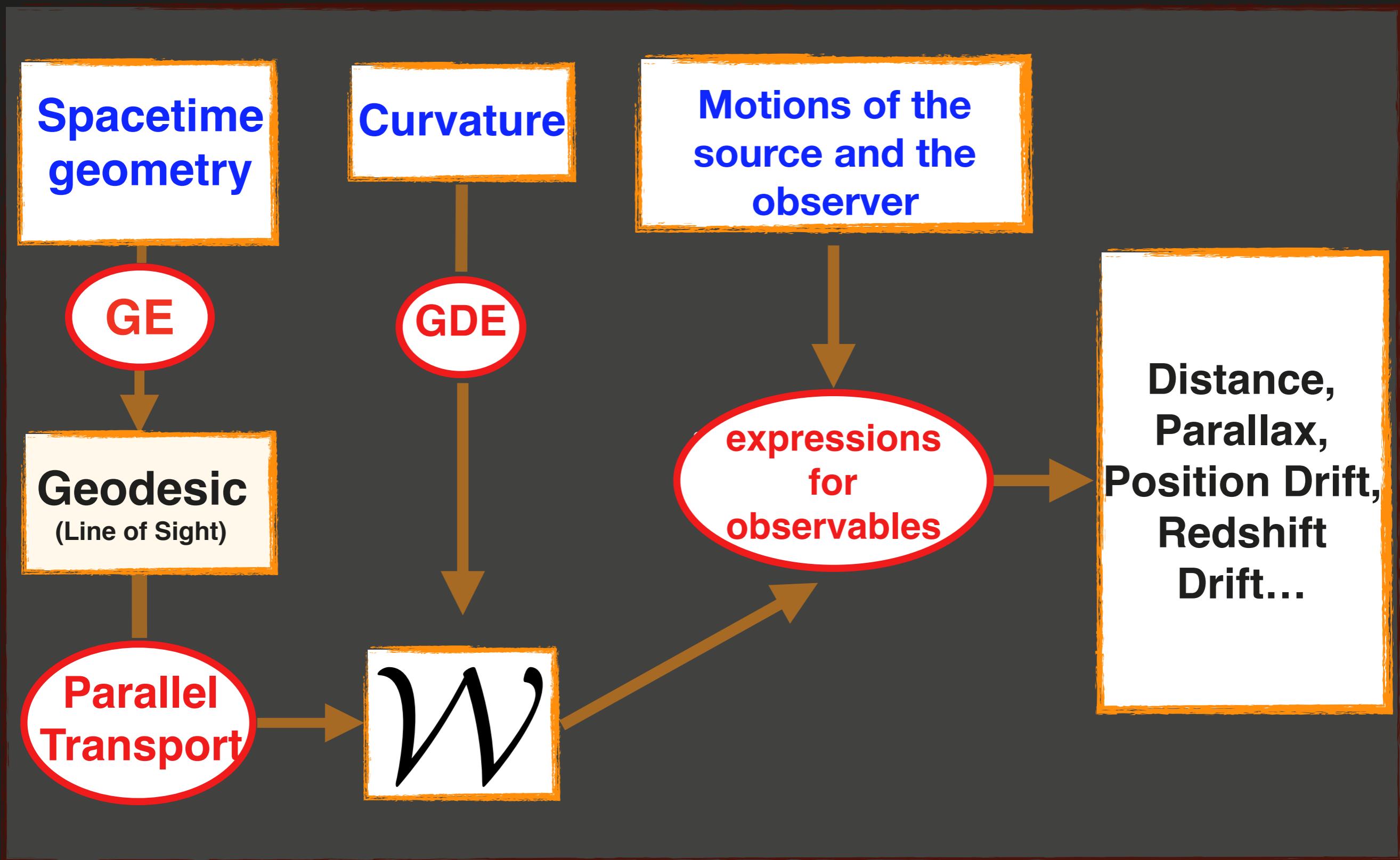
- BGOs as unified framework for **all** the observables and **all** scales

Work in progress:

- investigating non-linearities
- the full numerical tool for \mathcal{W} within the EINSTEIN TOOLKIT



APPENDIX: Recipe for computing observables using BGO's



APPENDIX: Relation between BGO and standard geometric optics

$$\begin{pmatrix} \delta x_{\mathcal{E}} \\ \Delta l_{\mathcal{E}} \end{pmatrix} = \mathcal{W} \begin{pmatrix} \delta x_{\mathcal{O}} \\ \Delta l_{\mathcal{O}} \end{pmatrix}$$

$$\delta x_{\mathcal{E}}^{\mu} = (W_{XX})^{\mu}_{\nu} \delta x_{\mathcal{O}}^{\nu} + (W_{XL})^{\mu}_{\nu} \Delta l_{\mathcal{O}}^{\nu}$$

$$\Delta l_{\mathcal{E}}^{\mu} = (W_{LX})^{\mu}_{\nu} \delta x_{\mathcal{O}}^{\nu} + (W_{LL})^{\mu}_{\nu} \Delta l_{\mathcal{O}}^{\nu}$$

$$(W_{XL})^{(a)}_{(b)} = \begin{pmatrix} X & X & X & X \\ X & D^A_B & X & X \\ X & & X & X \\ X & x & x & X \end{pmatrix} \quad (W_{XX})^{(a)}_{(b)} = \begin{pmatrix} X & X & X & X \\ X & (\delta + m)^A_B & X & X \\ X & & X & X \\ X & x & x & X \end{pmatrix}$$

APPENDIX: Relation between BGO and different notions of distances

$$\mathcal{D}_{\text{ang}} = (p_{\mathcal{O} \mu} u_{\mathcal{O}}^{\mu}) \left| \det D_{(B)}^{(A)} \right|^{\frac{1}{2}}$$

$$\mathcal{D}_{\text{lum}} = (1+z)^2 \mathcal{D}_{\text{ang}}$$

$$\mathcal{D}_{\text{par}} = (p_{\mathcal{O} \mu} u_{\mathcal{O}}^{\mu}) \left| \det D_{(B)}^{(A)} \right|^{\frac{1}{2}} \left| \det \left(\delta_{(B)}^{(A)} + m_{(B)}^{(A)} \right) \right|^{-\frac{1}{2}}$$

APPENDIX: Plane-parallel metric

$$ds^2 = a^2(\eta) \left\{ -c^2 d\eta^2 + \gamma_{11}(\eta, x) dx^2 + \gamma_{22}(\eta, x) dy^2 + \gamma_{33}(\eta, x) dz^2 \right\}$$

- $\gamma_{11} = \left(1 - \frac{2}{3} \frac{D \partial_1^2 \phi_0}{\mathcal{H}_0^2 \Omega_{m0}} \right)^2 + \frac{1}{c^2} \left[-\frac{10}{3} \phi_0 + (4a_{nl} - 5) \frac{10}{9} \frac{D(\partial_1 \phi_0)^2}{\mathcal{H}_0^2 \Omega_{m0}} + (a_{nl} - 1) \frac{40}{9} \frac{D \phi_0 \partial_1^2 \phi_0}{\mathcal{H}_0^2 \Omega_{m0}} \right.$

$$\left. - (32a_{nl} - 49) \frac{5}{54} \frac{D^2 (\partial_1 \phi_0)^2 \partial_1^2 \phi_0}{(\mathcal{H}_0^2 \Omega_{m0})^2} + \left(1 - \frac{2}{3} a_{nl} \right) \frac{40}{9} \frac{D^2 \phi_0 (\partial_1^2 \phi_0)^2}{(\mathcal{H}_0^2 \Omega_{m0})^2} - \frac{5}{9} \frac{D^3 (\partial_1 \phi_0)^2 (\partial_1^2 \phi_0)^2}{(\mathcal{H}_0^2 \Omega_{m0})^3} \right]$$
- $\gamma_{22} = \gamma_{33} = 1 + \frac{1}{c^2} \left[\frac{10}{9} \left(\frac{D(\partial_1 \phi_0)^2}{\mathcal{H}_0^2 \Omega_{m0}} - 3\phi_0 \right) \right].$