



# Special and General Relativity

Boud Roukema

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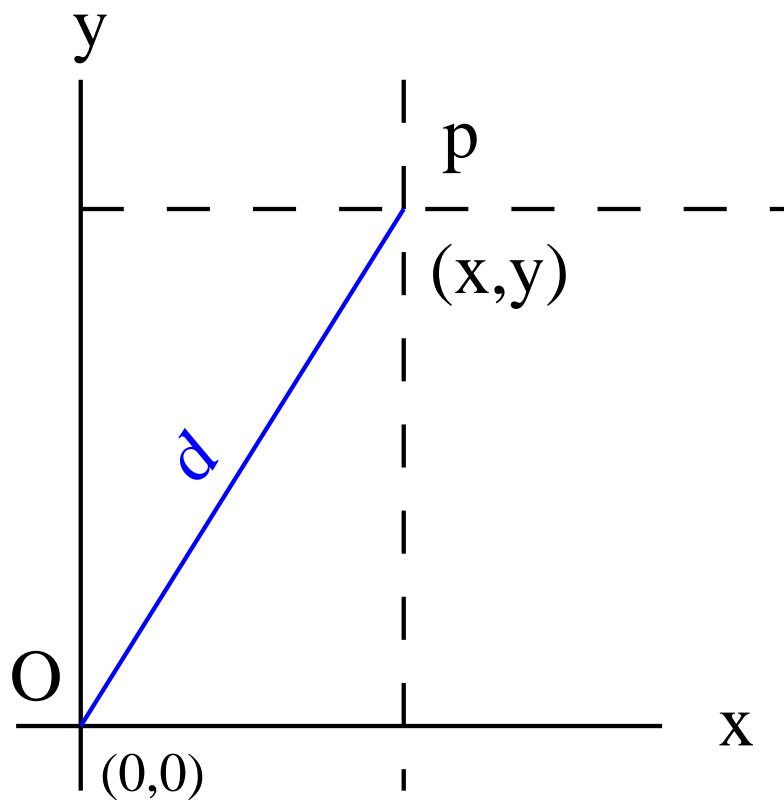
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  - $\rightarrow$  differentiable 4-(pseudo-)manifold
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- SR: spacetime = w:Minkowski space
- GR: spacetime = a solution of the  
w:Einstein field equations





# SR: Minkowski spacetime

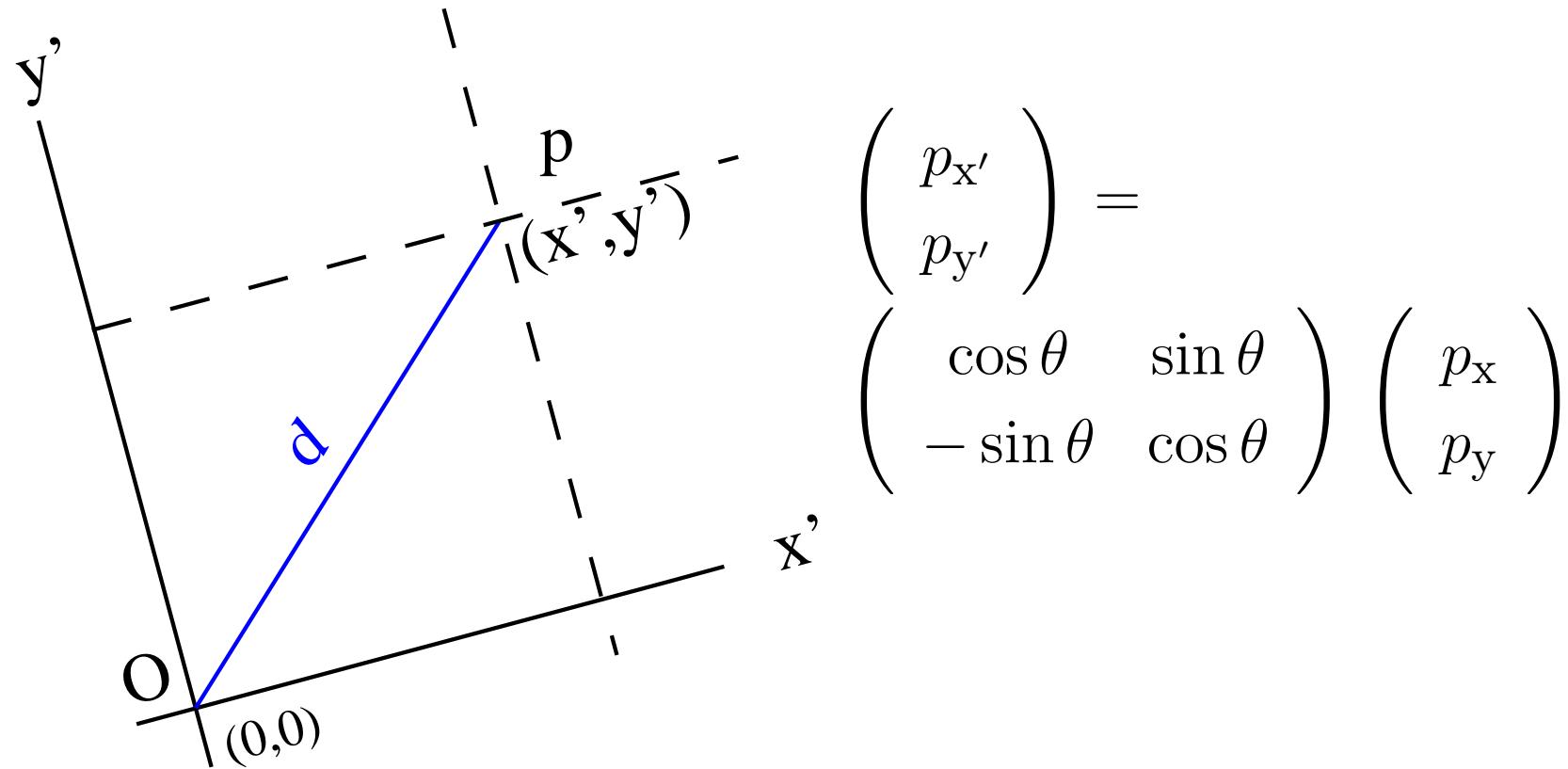


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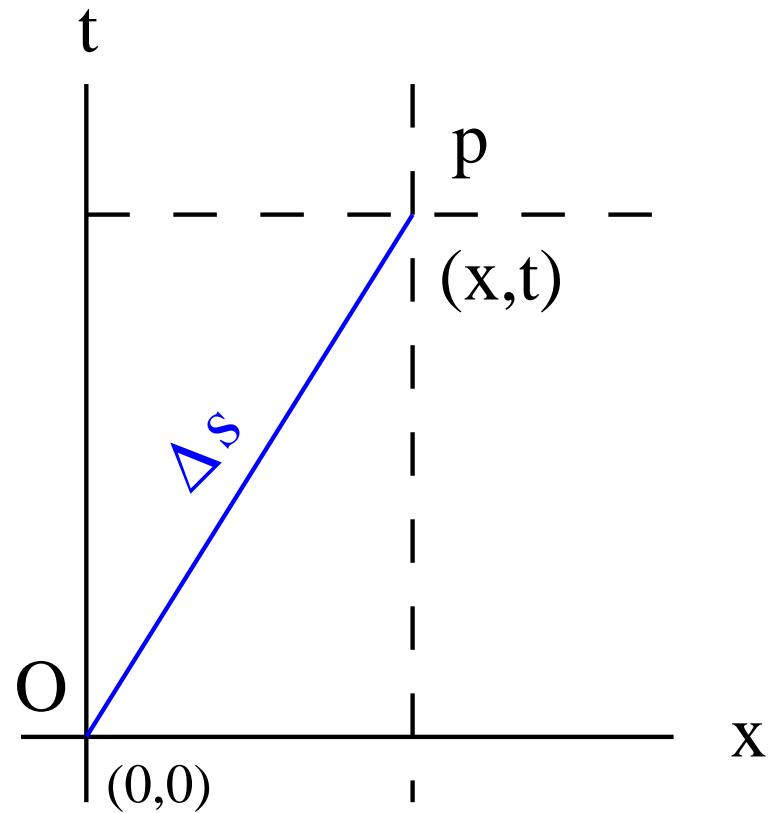


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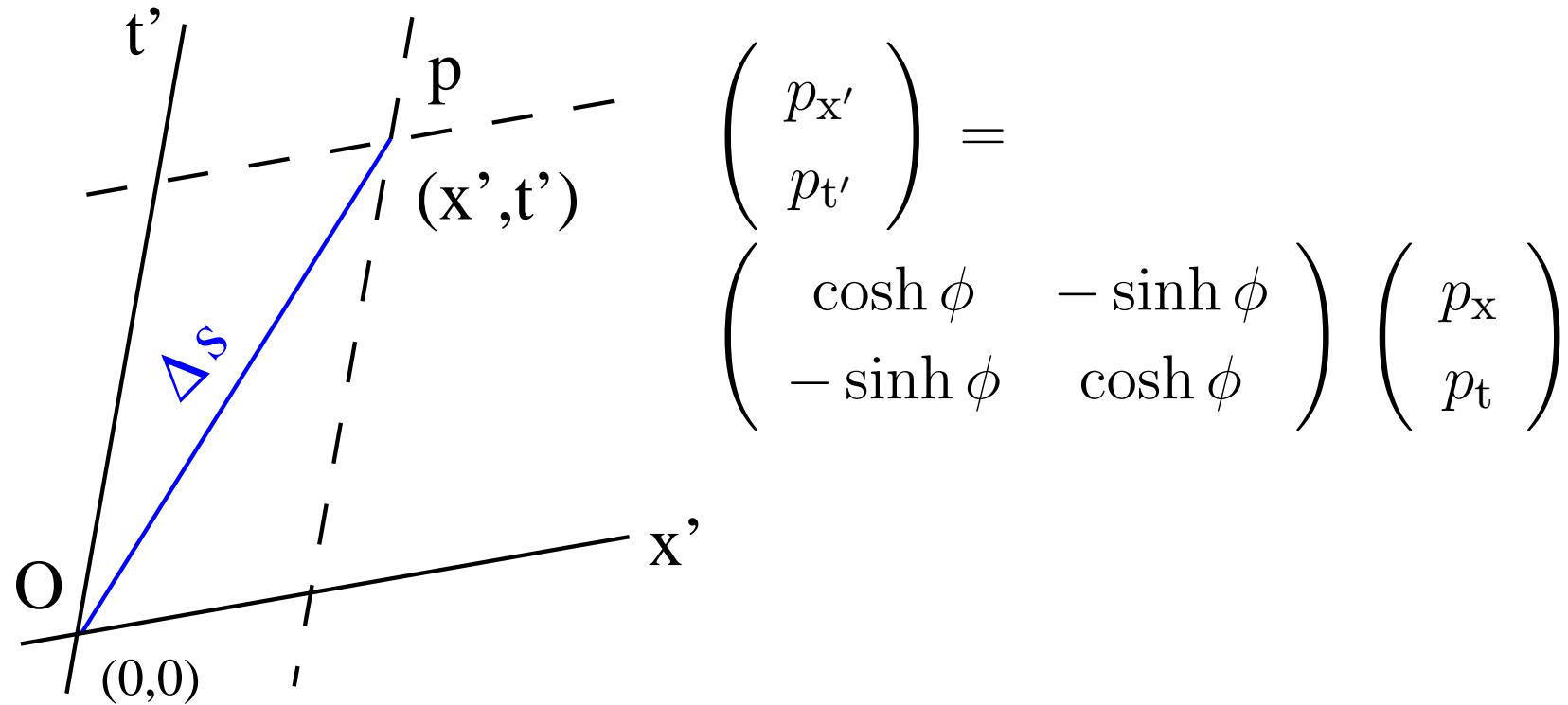


$p$  at  $(x, t)$ , w:invariant interval from observer at  $O$  is  $\Delta s$   
where  $(\Delta s)^2 = -(\Delta t)^2 + (\Delta x)^2$





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$$\Lambda(\phi) := \begin{pmatrix} \cosh \phi & -\sinh \phi \\ -\sinh \phi & \cosh \phi \end{pmatrix}$$

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**L** =  $\frac{1 \text{ s}}{1 \text{ s}} = 1$  (dimensionless)



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where velocity  $\beta := v/c \equiv v = \tanh \phi$





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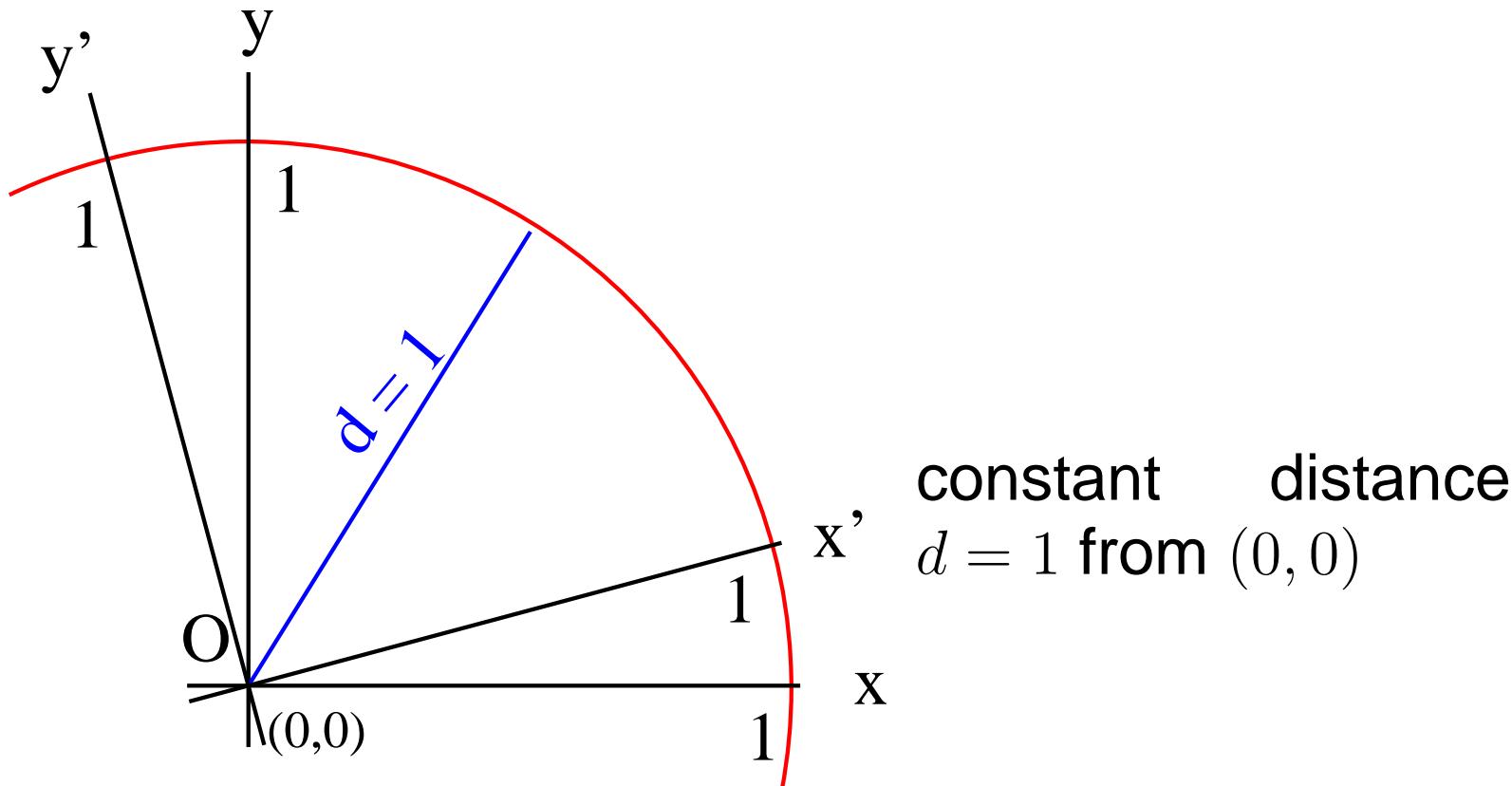




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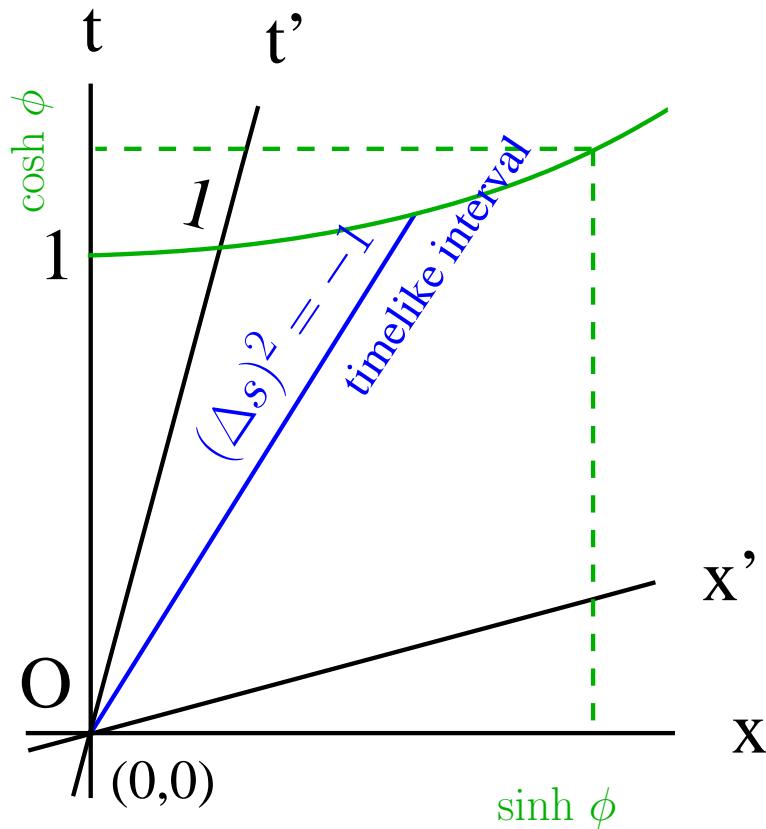




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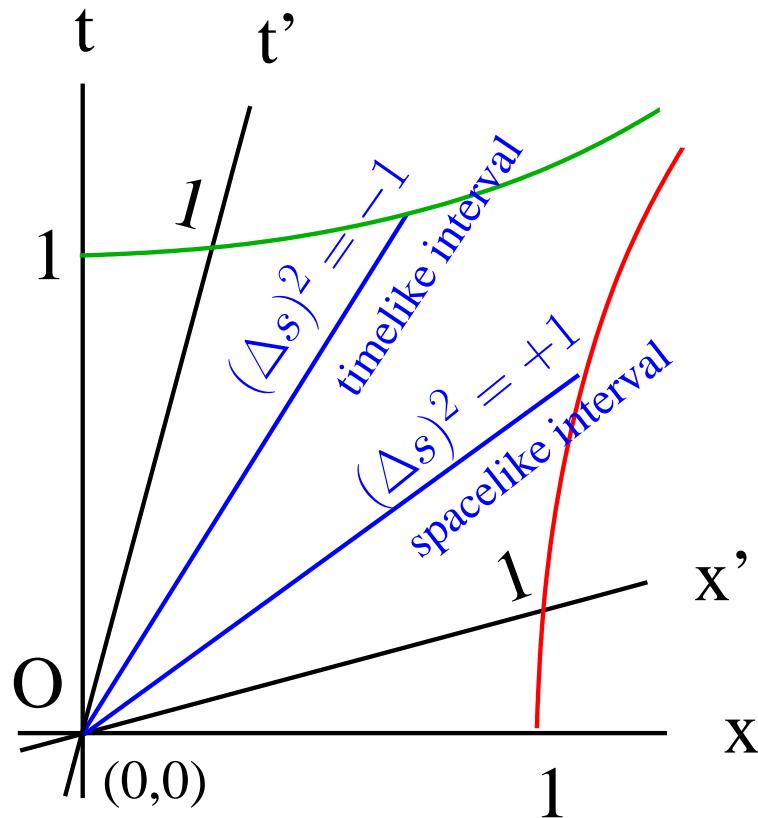




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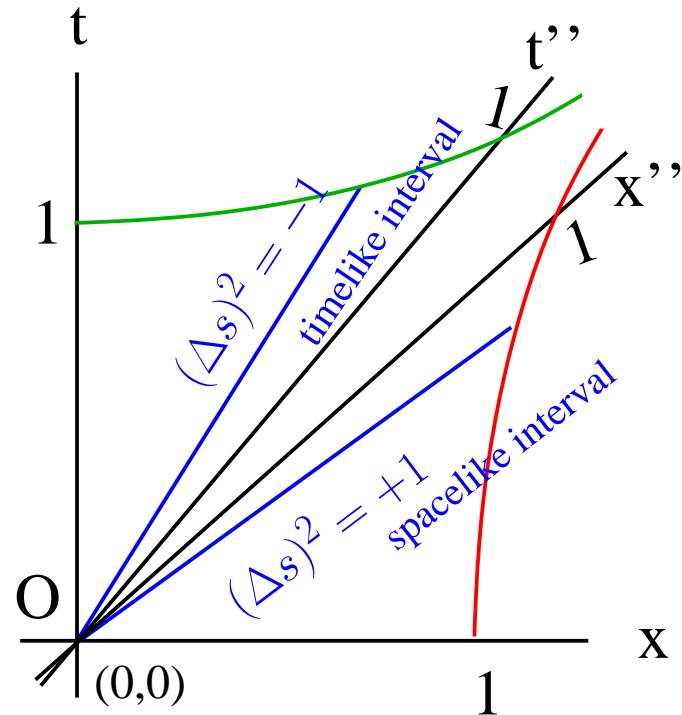


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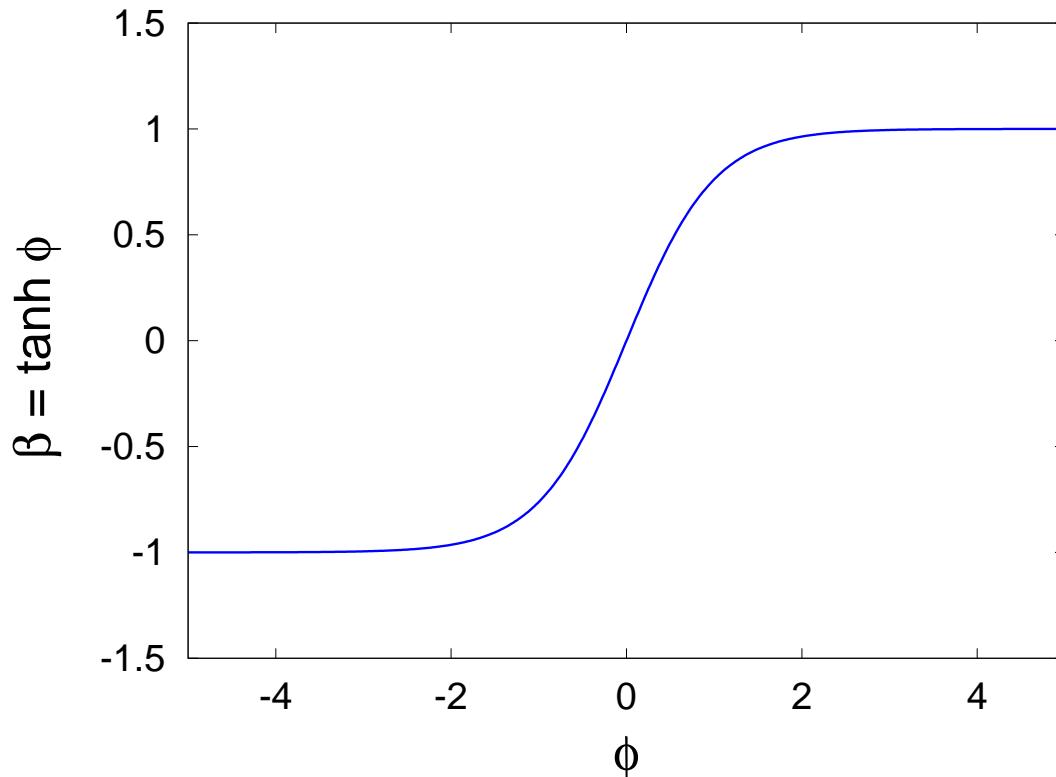


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- w:[Michelson-Morley experiment](#) (1887)





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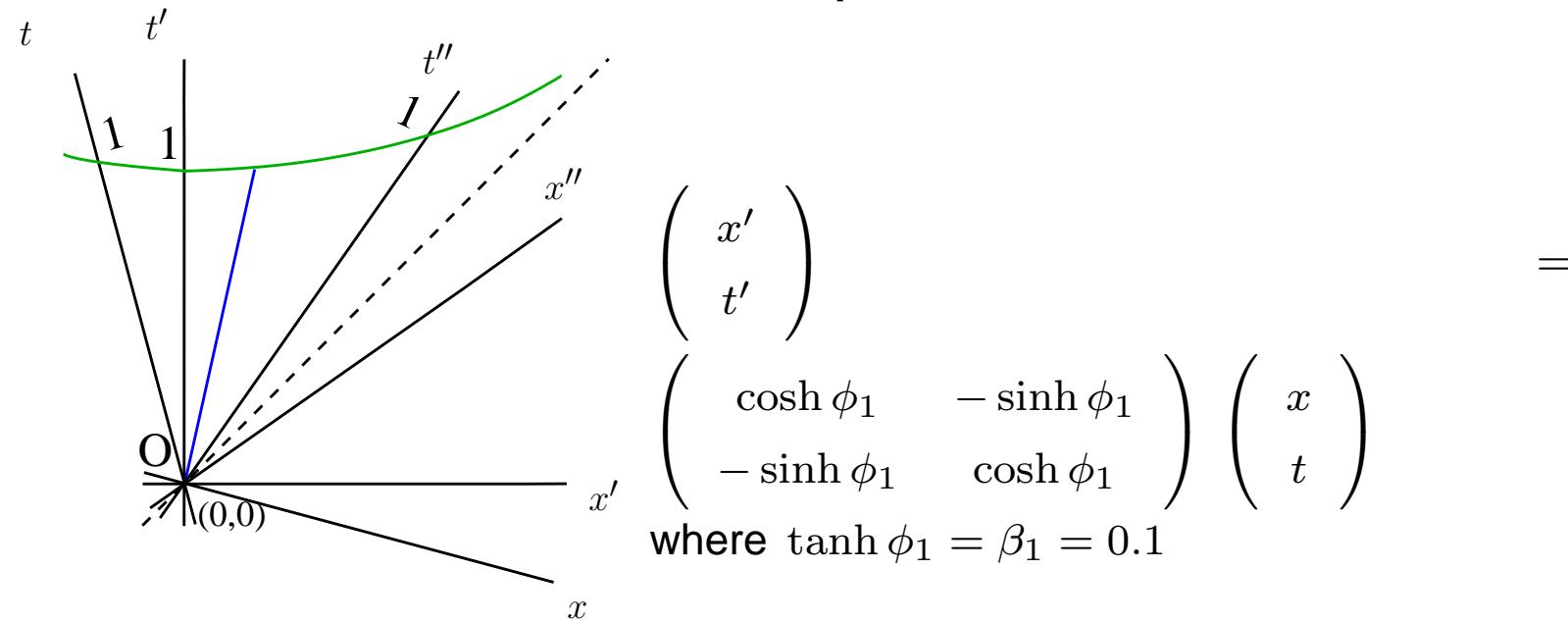
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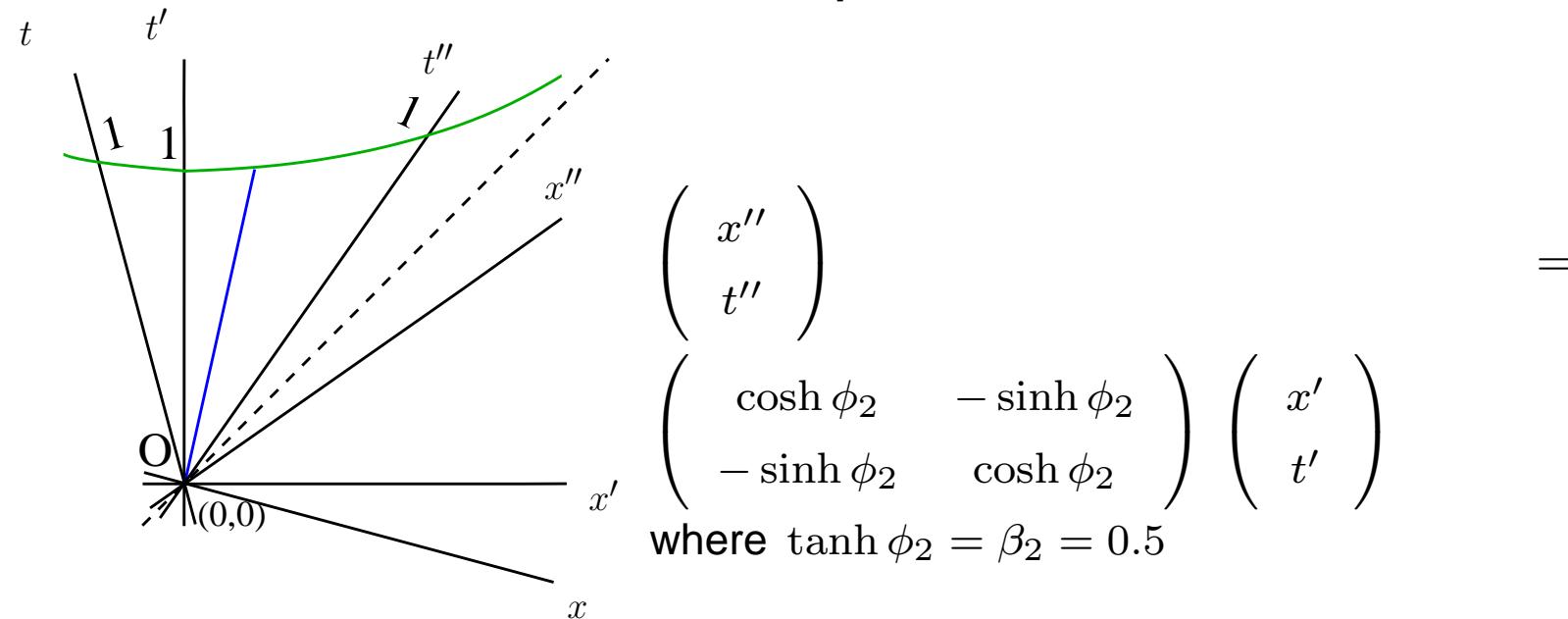
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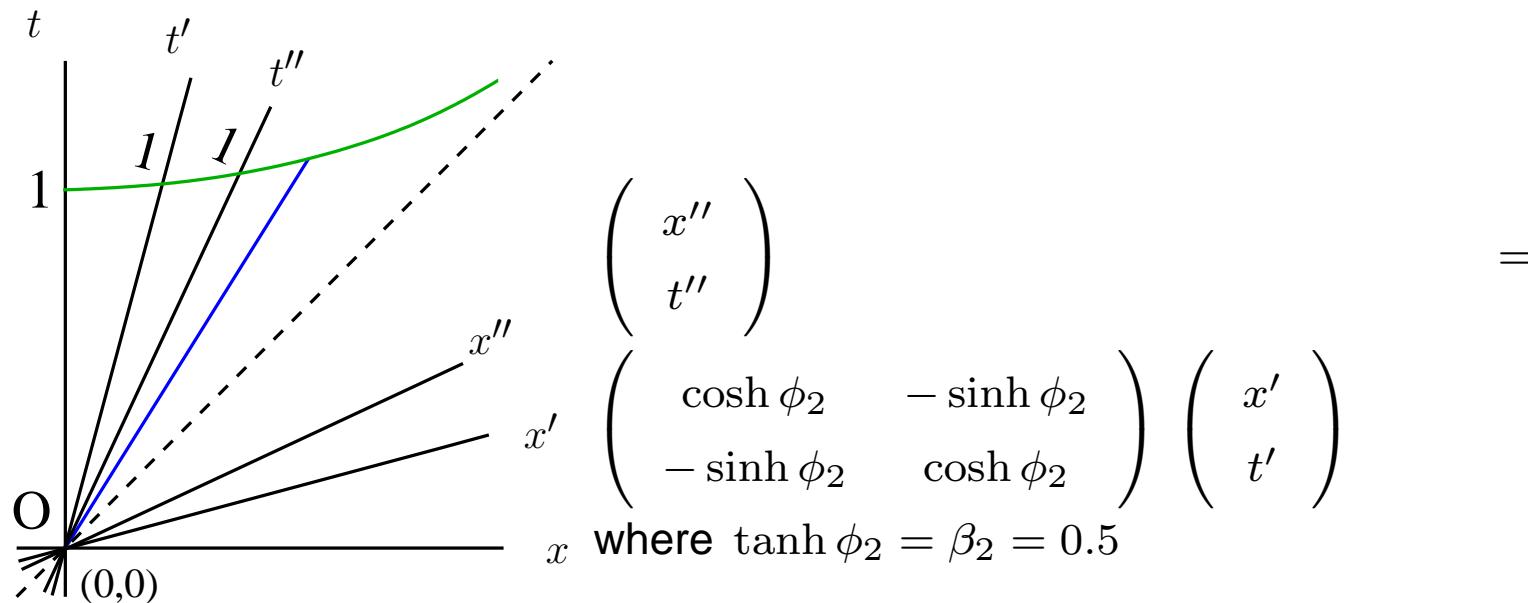
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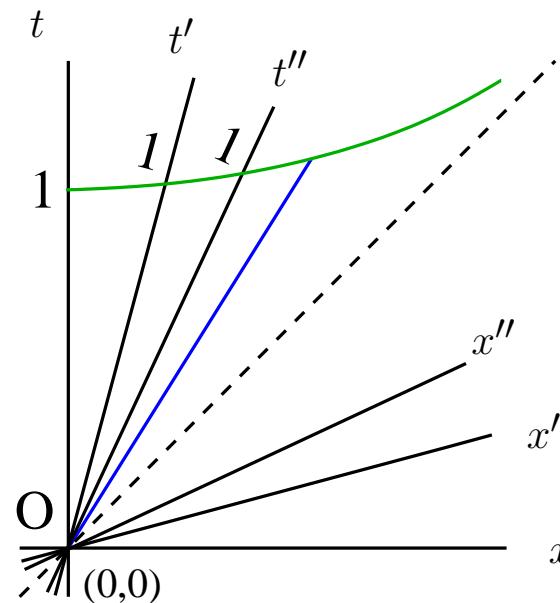
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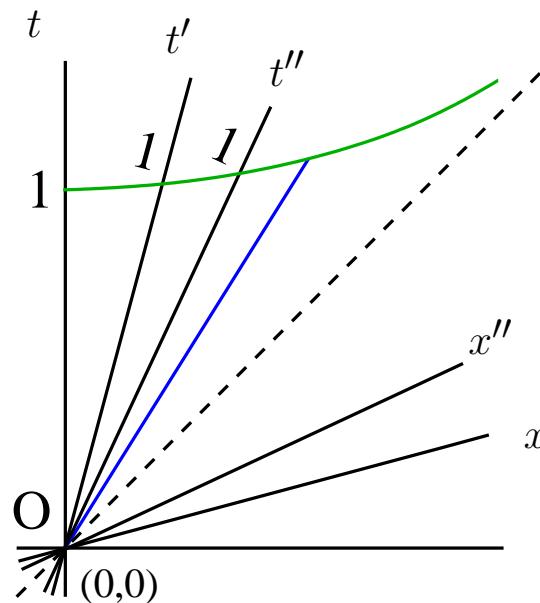
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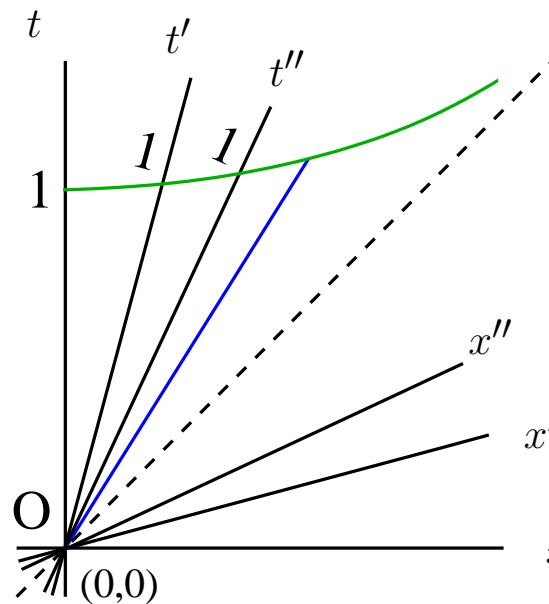
but  $\Lambda(\phi_2)\Lambda(\phi_1) = \Lambda(\phi_1 + \phi_2)$





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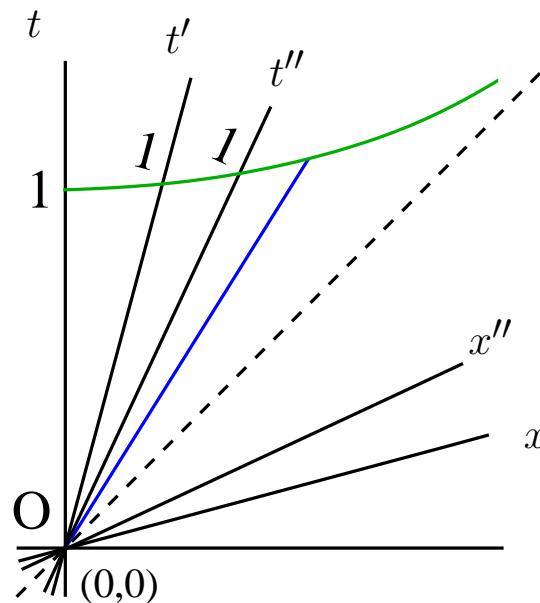
cf. rotation  $\theta_1$  “plus” rotation  $\theta_2$  = rotation  $\theta_1 + \theta_2$





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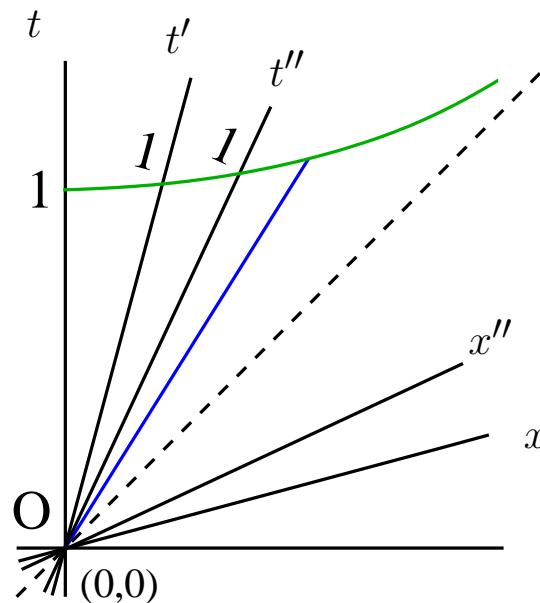
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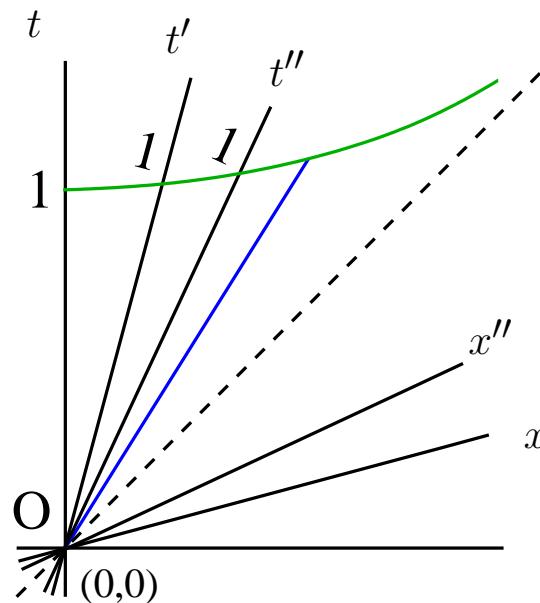
$$\text{so } \beta_3 = \tanh(\phi_1 + \phi_2) = \frac{\tanh \phi_1 + \tanh \phi_2}{1 + \tanh \phi_1 \tanh \phi_2}$$





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interstellar ark travels at  $\beta_1 = 0.1$  from Sun, sends out rocket at  $\beta_2 = 0.5$ ; rocket's speed  $\beta_3$  in Sun frame =?



$$\begin{pmatrix} x'' \\ t'' \end{pmatrix} = \Lambda(\phi_2)\Lambda(\phi_1) \begin{pmatrix} x \\ t \end{pmatrix}$$

but  $\Lambda(\phi_2)\Lambda(\phi_1) = \Lambda(\phi_1 + \phi_2)$

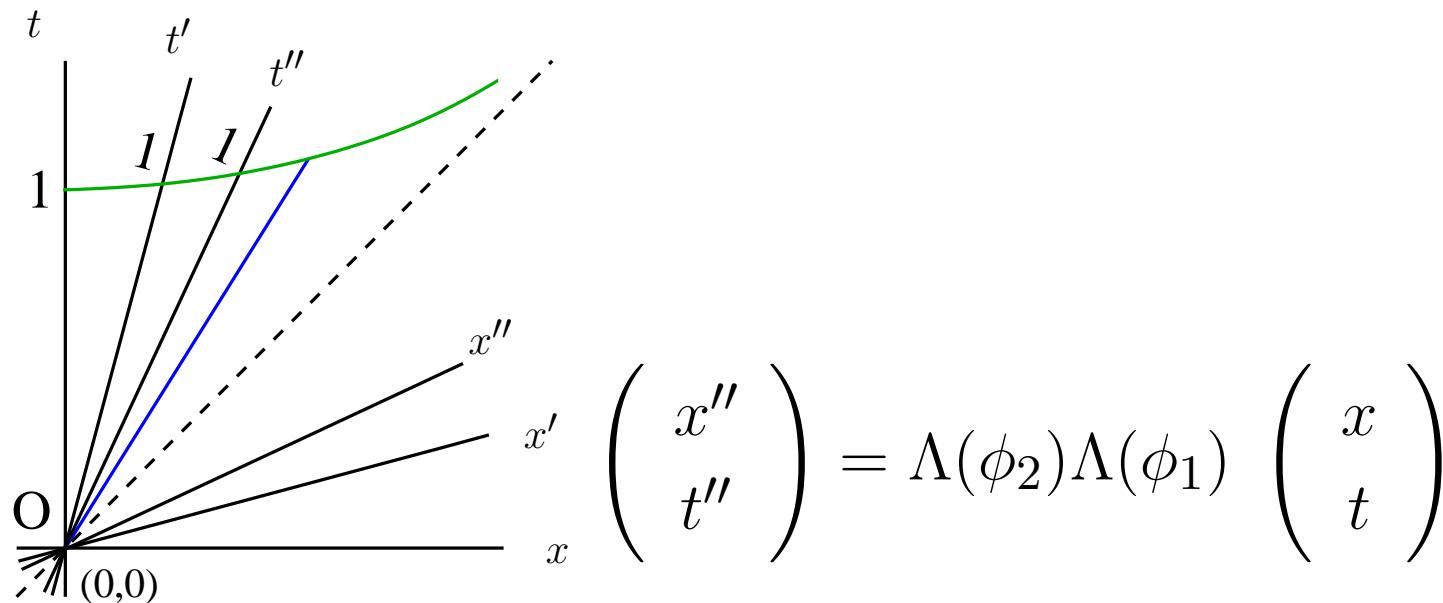
$$\text{so } \beta_3 = \tanh(\phi_1 + \phi_2) = \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2}$$





# SR: adding velocities

interstellar ark travels at  $\beta_1 = 0.1$  from Sun, sends out rocket at  $\beta_2 = 0.5$ ; rocket's speed  $\beta_3$  in Sun frame =?



$$\text{but } \Lambda(\phi_2)\Lambda(\phi_1) = \Lambda(\phi_1 + \phi_2)$$

$$\text{so } \beta_3 = \tanh(\phi_1 + \phi_2) = \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} = \frac{0.1 + 0.5}{1 + 0.1 \times 0.5} \approx 0.57$$





# SR: Lorentz factor

$\Lambda$ : alternative to hyperbolic trig functions





# SR: Lorentz factor

$\Lambda$ : alternative to hyperbolic trig functions

$$\Lambda(\phi) := \begin{pmatrix} \cosh \phi & -\sinh \phi \\ -\sinh \phi & \cosh \phi \end{pmatrix}$$

$$\cosh \phi := \frac{e^\phi + e^{-\phi}}{2}$$

$$\sinh \phi := \frac{e^\phi - e^{-\phi}}{2}$$

w:**hyperbolic function**





# SR: Lorentz factor

$\Lambda$ : alternative to hyperbolic trig functions

$$\Lambda(\beta) := \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix}$$

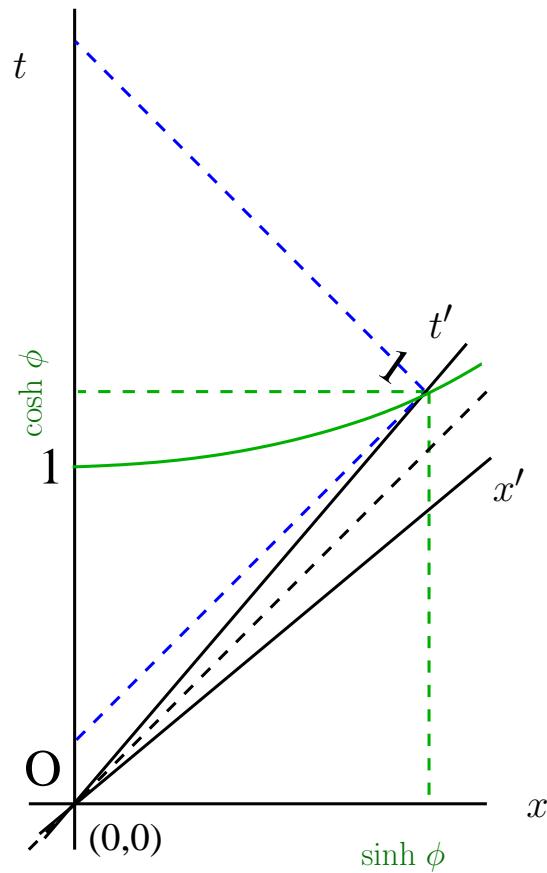
$$\begin{aligned}\beta &= \tanh \phi \\ \gamma &:= (1 - \beta^2)^{-1/2} = \\ &\text{Lorentz factor}\end{aligned}$$

$$\begin{aligned}\gamma &= \cosh \phi \\ \beta\gamma &= \sinh \phi\end{aligned}$$



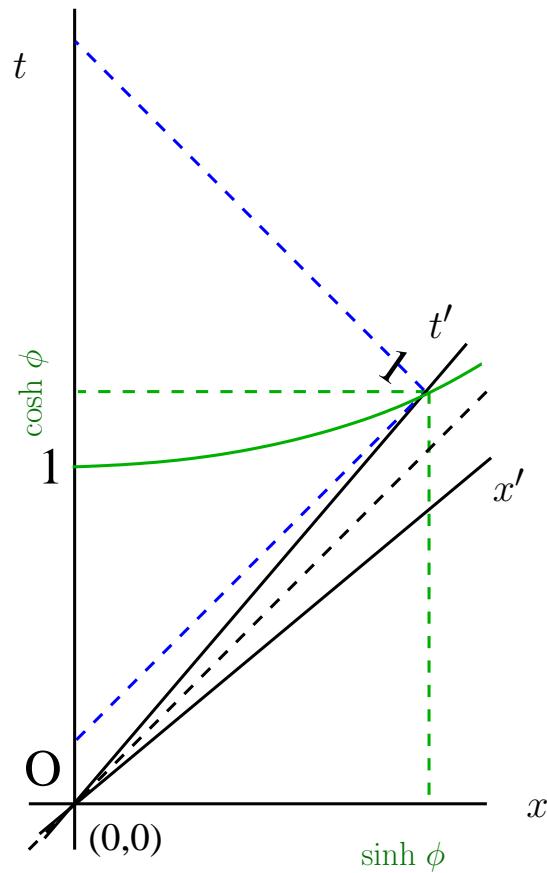


# SR: worldline time dilation





# SR: worldline time dilation

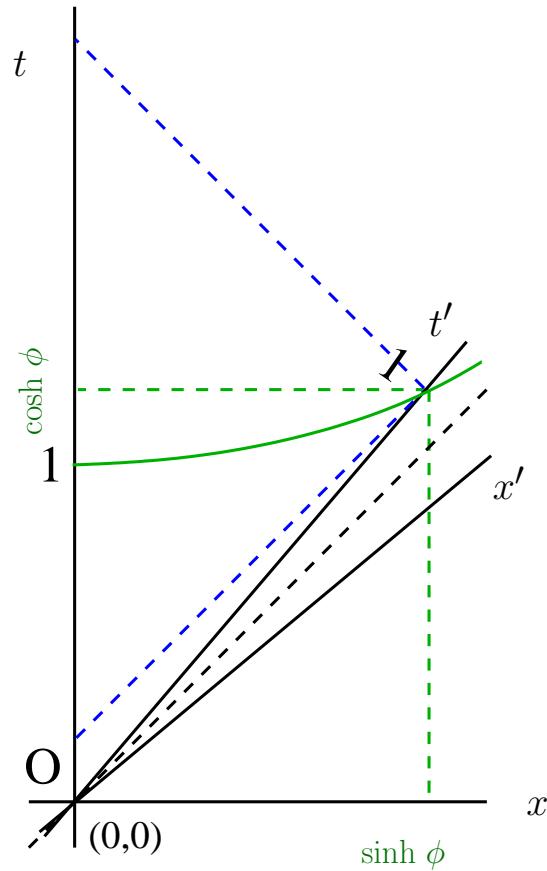


$$\cosh \phi \equiv \gamma$$





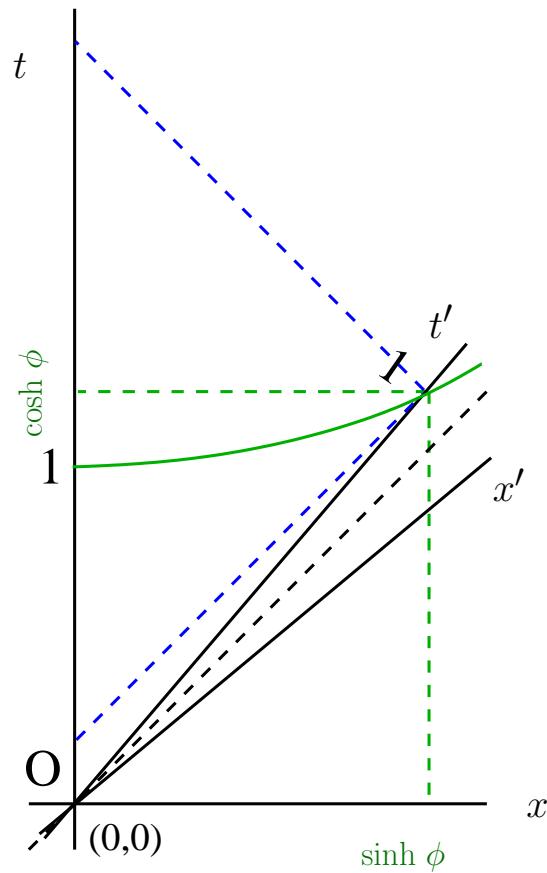
# SR: worldline time dilation



$$\cosh \phi \equiv \gamma \equiv \frac{1}{\sqrt{1-\beta^2}}$$



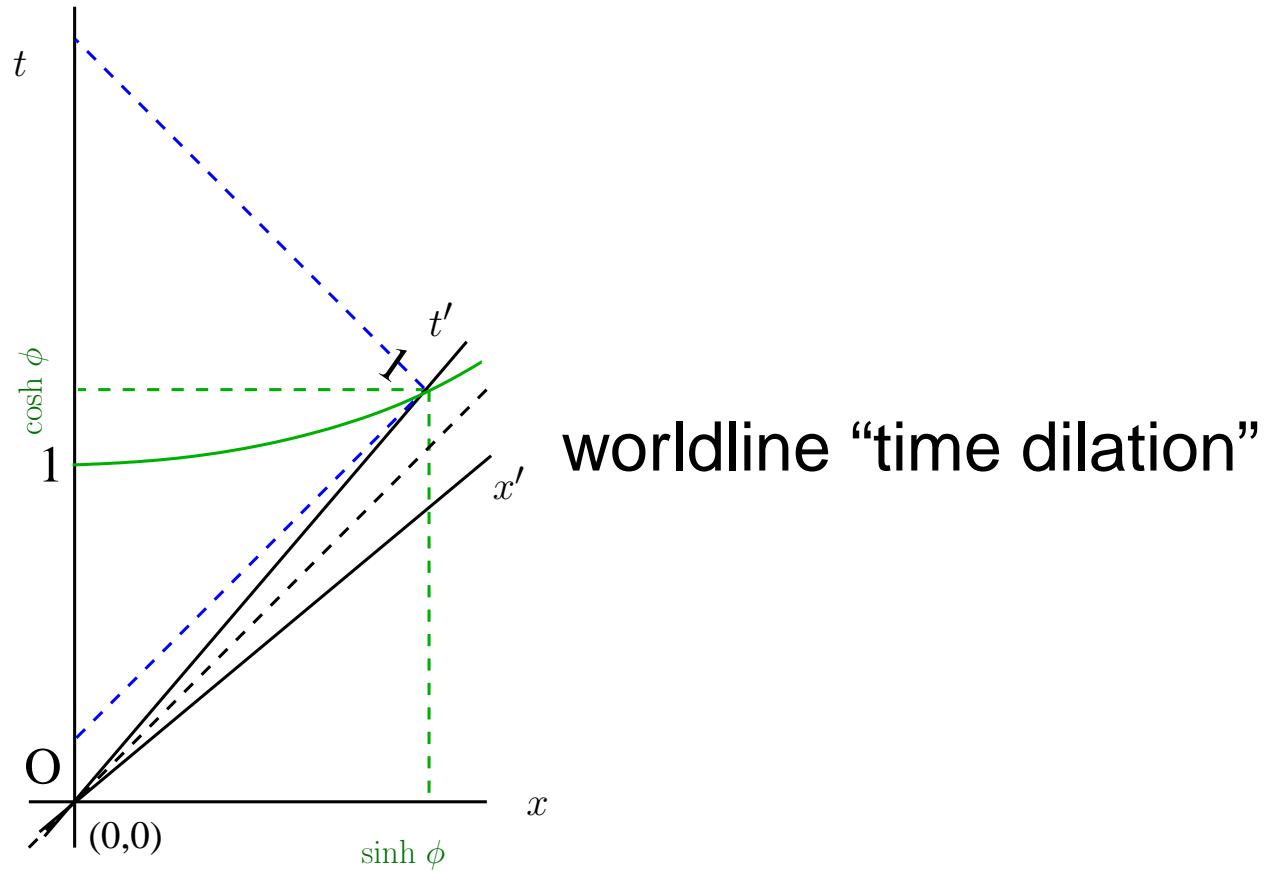
# SR: worldline time dilation



$$\cosh \phi \equiv \gamma \equiv \frac{1}{\sqrt{1-\beta^2}} > 1$$



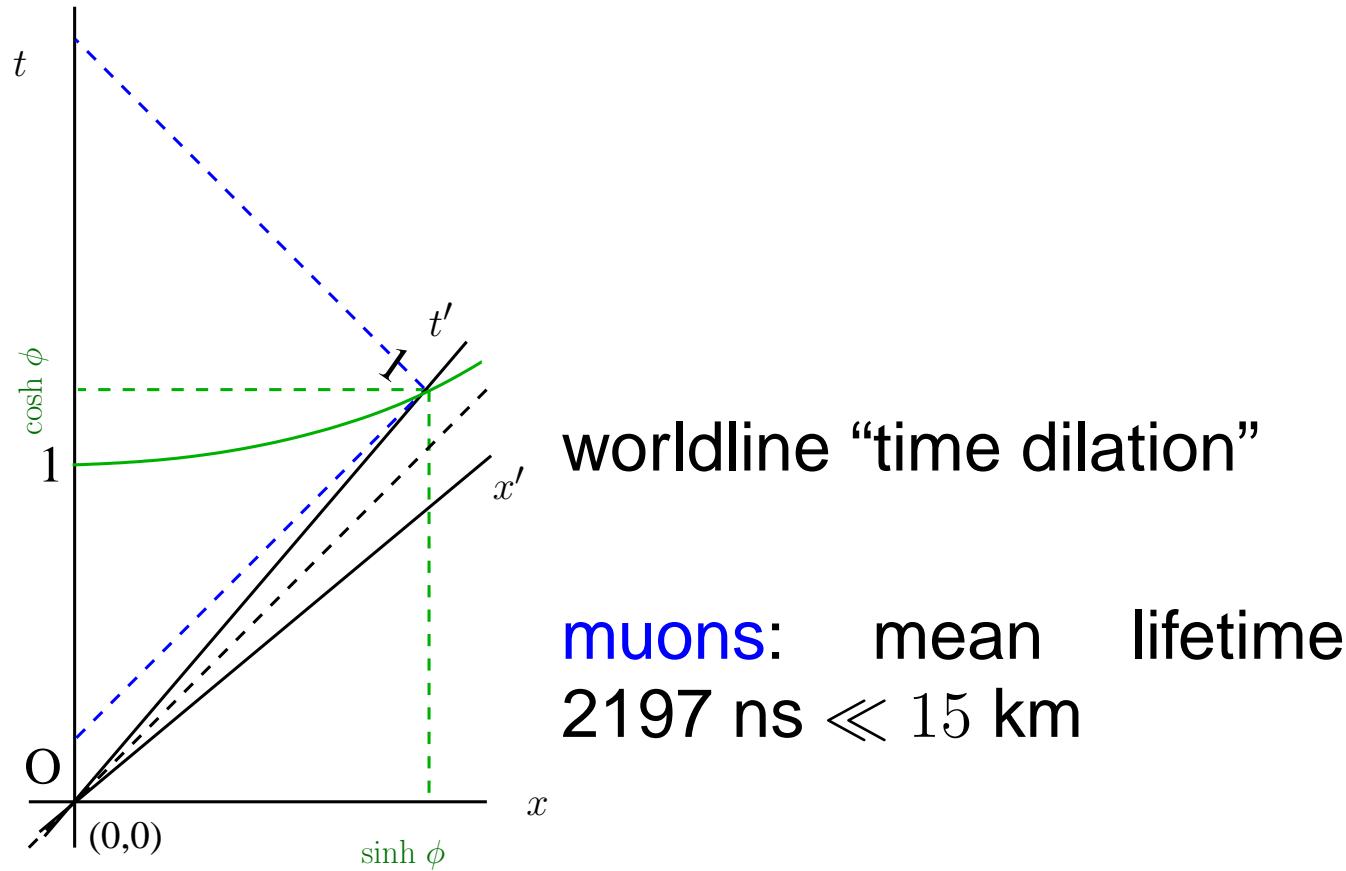
# SR: worldline time dilation



$$\cosh \phi \equiv \gamma \equiv \frac{1}{\sqrt{1-\beta^2}} > 1$$

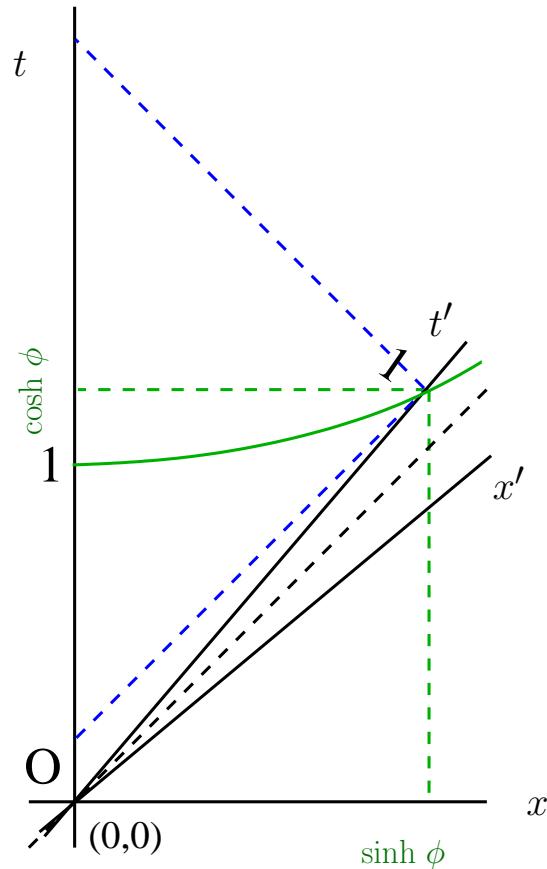


# SR: worldline time dilation



$$\cosh \phi \equiv \gamma \equiv \frac{1}{\sqrt{1-\beta^2}} > 1$$

# SR: worldline time dilation



worldline “time dilation”

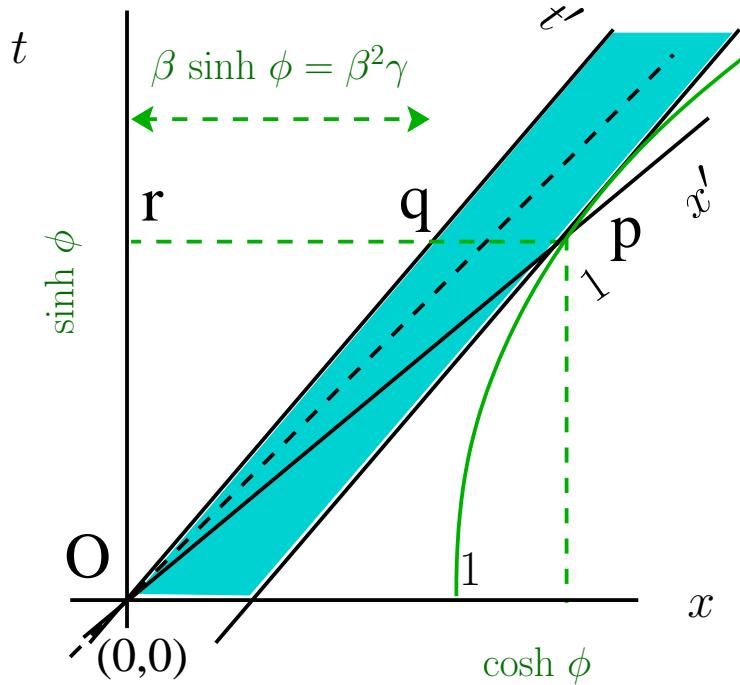
**muons:** mean lifetime  
 $2197 \text{ ns} \ll 15 \text{ km}$

time dilation  $\Rightarrow$  muons  
 can hit the ground

$$\cosh \phi \equiv \gamma \equiv \frac{1}{\sqrt{1-\beta^2}} > 1$$

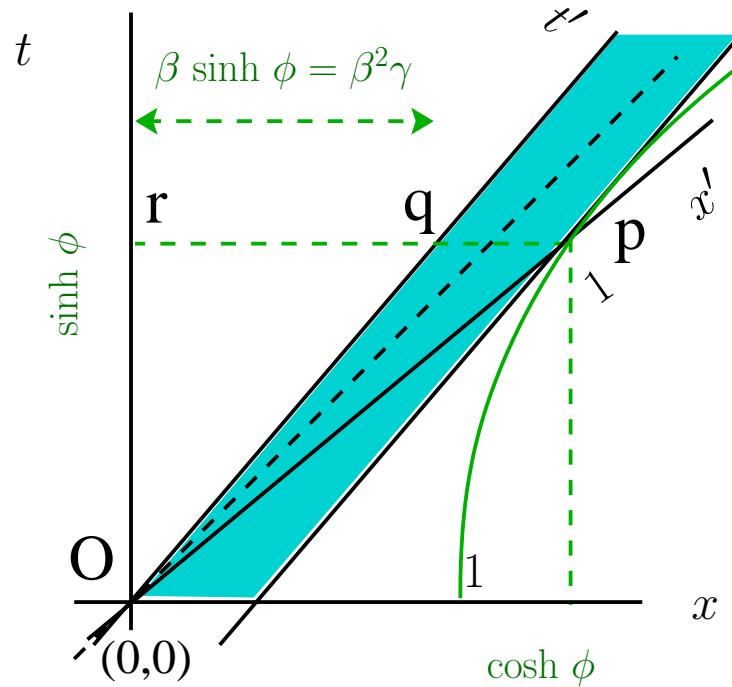


# SR: worldsheet space contraction





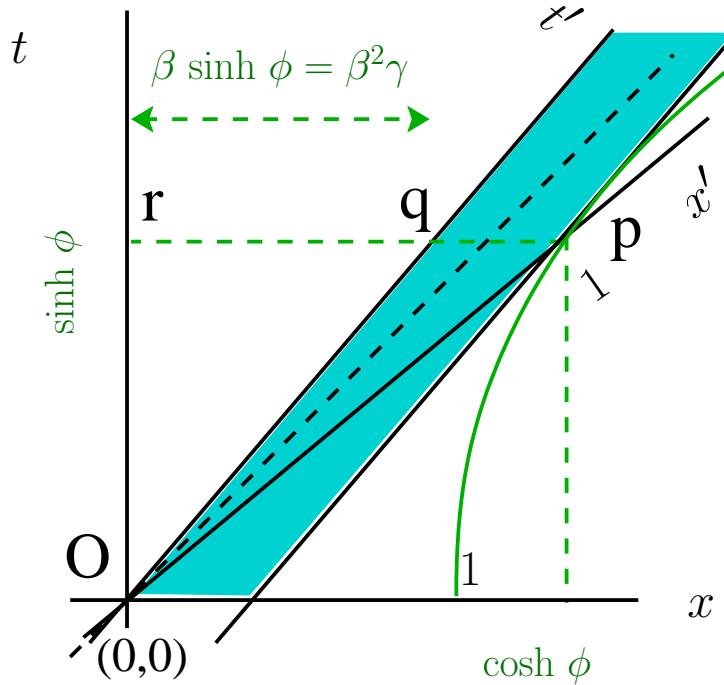
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} =$$



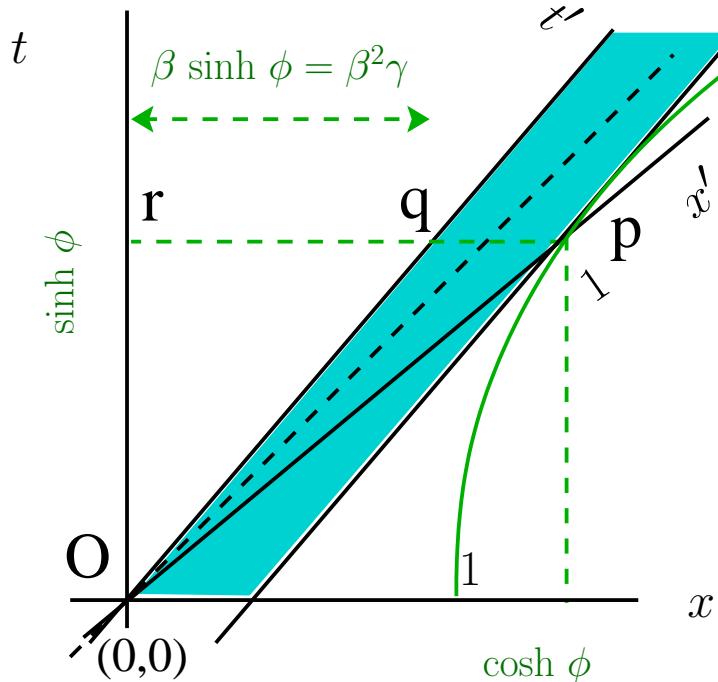
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = \cosh \phi - \beta \sinh \phi$$



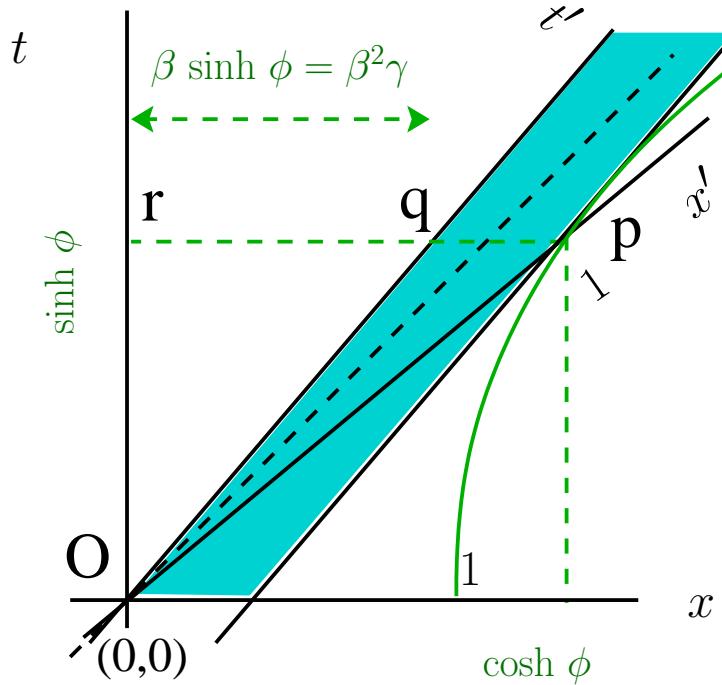
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = \gamma - \beta \beta \gamma$$



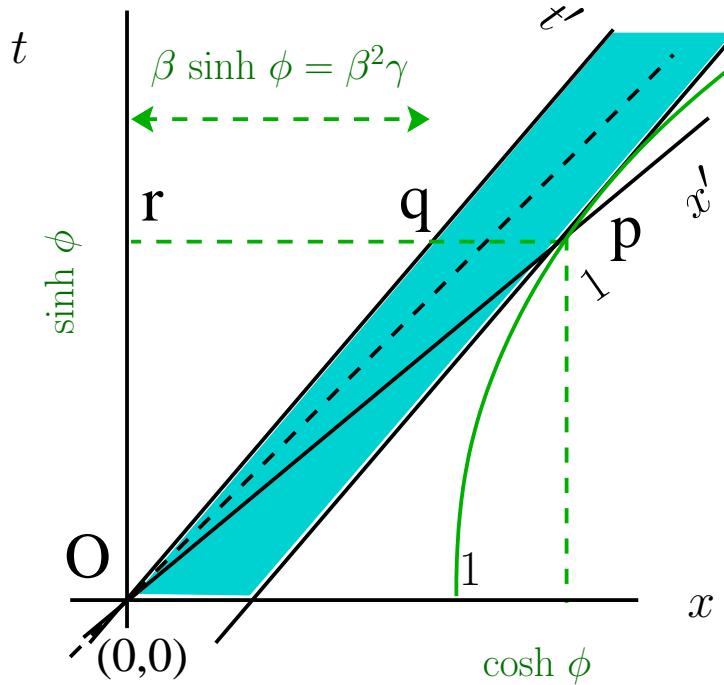
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = (1 - \beta^2)\gamma$$



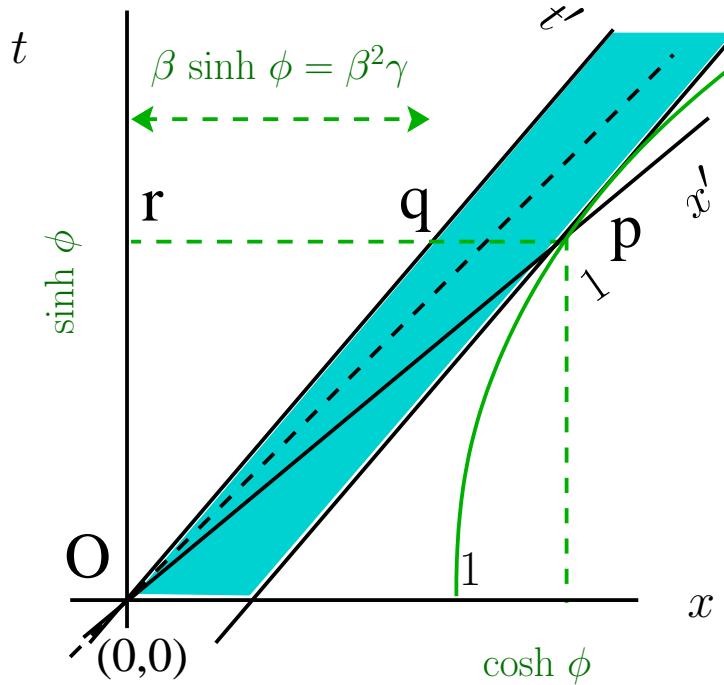
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = (1 - \beta^2)^{1+(-1/2)}$$



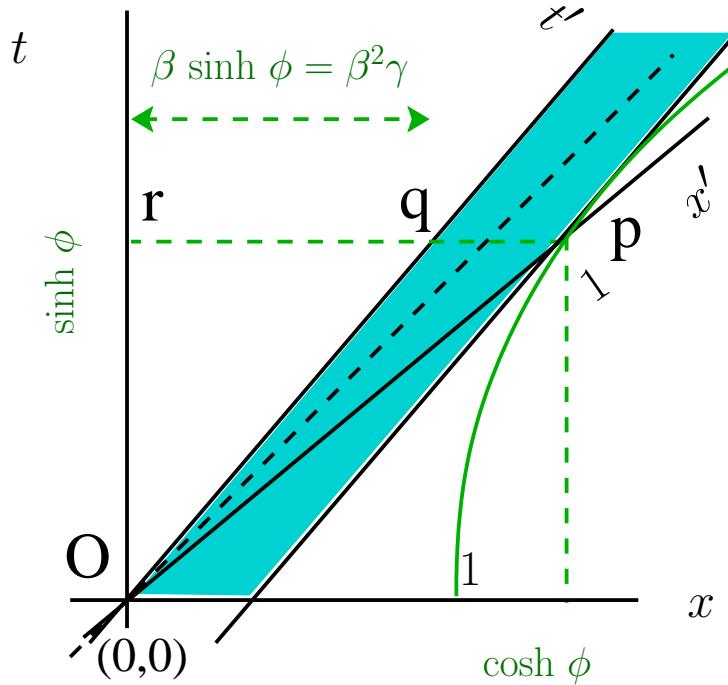
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = (1 - \beta^2)^{1/2}$$



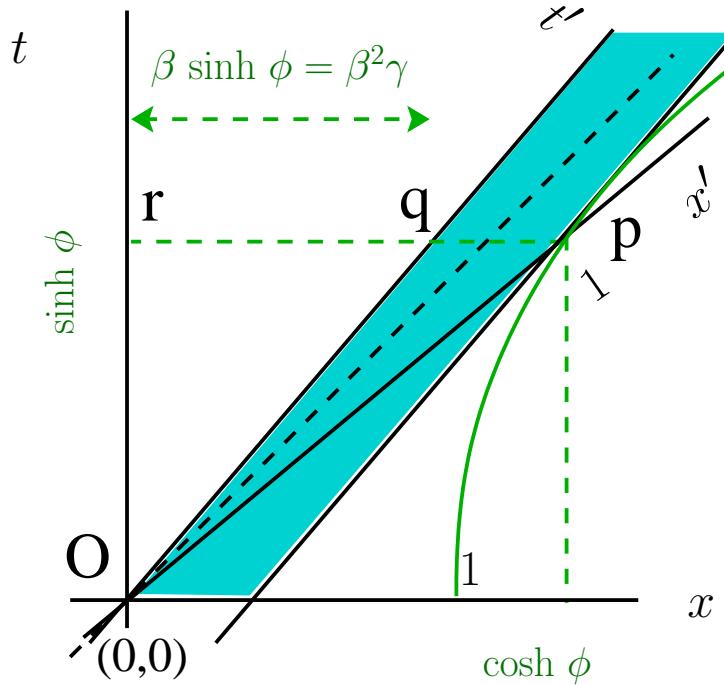
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = \gamma^{-1} < 1$$

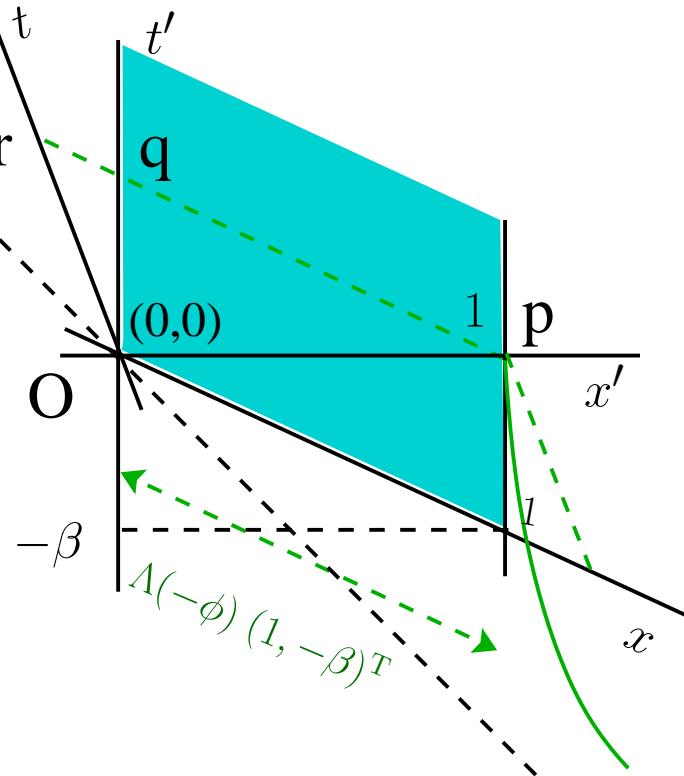


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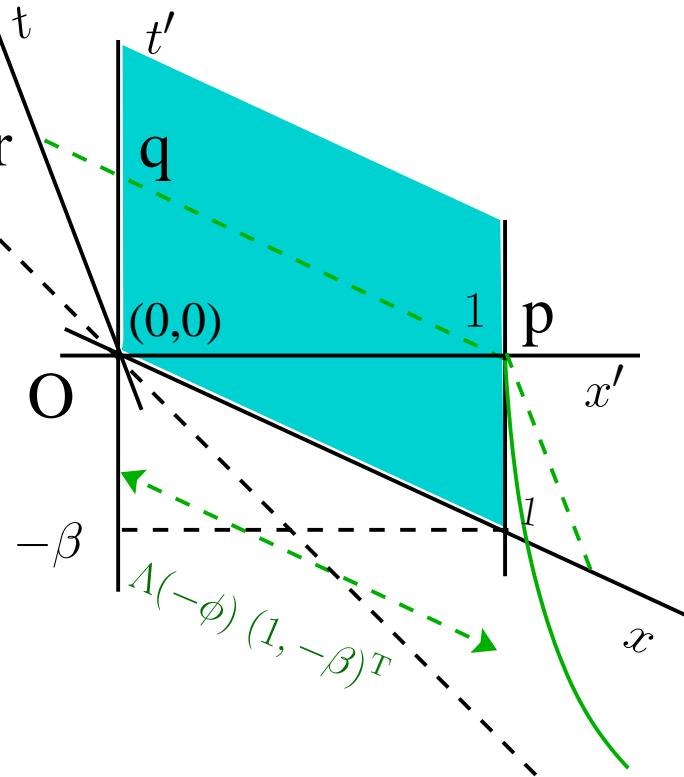


$$\sqrt{\Delta s^2(q, p)} = \gamma^{-1} < 1 \quad \text{worldsheet "space contraction"}$$

# SR: worldsheet space contraction

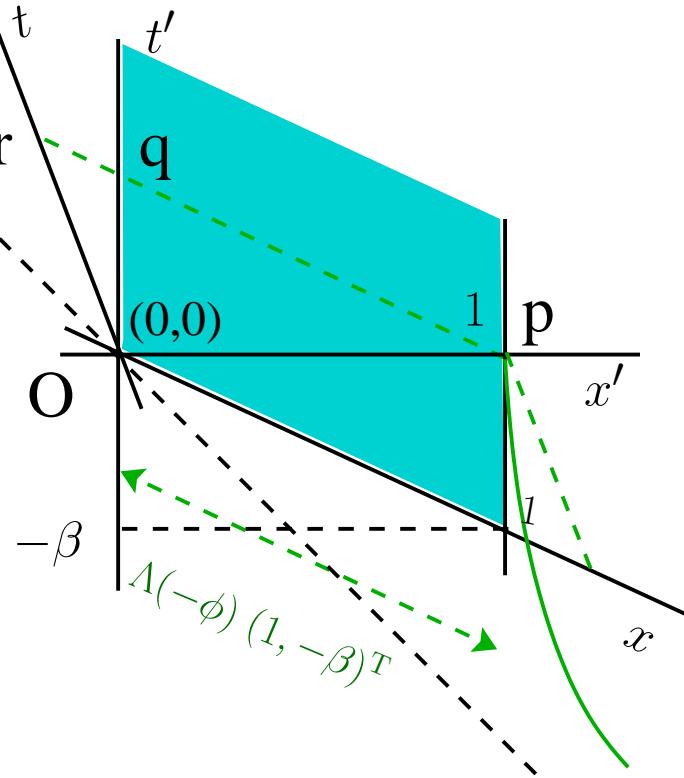


# SR: worldsheet space contraction



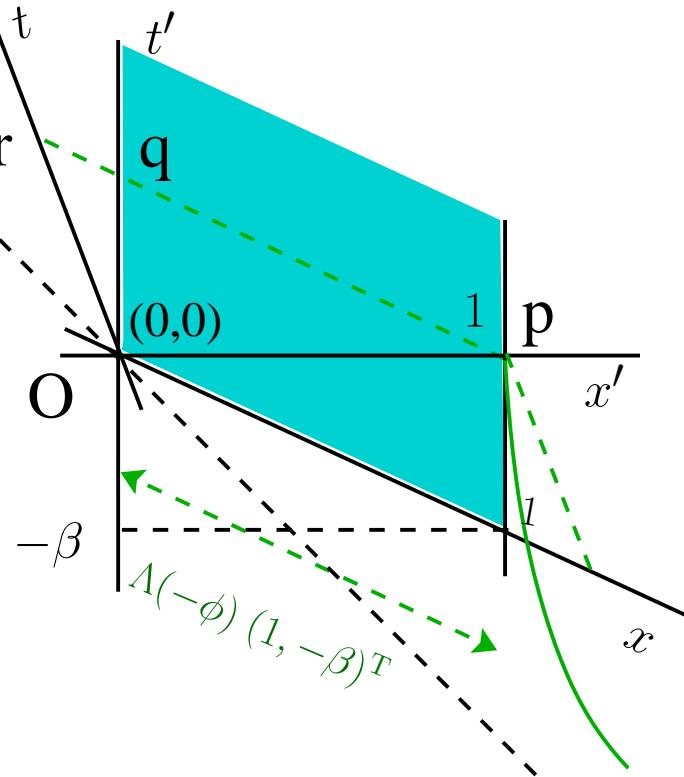
$$\Lambda^{-1} \begin{pmatrix} 1 \\ -\beta \end{pmatrix} = \begin{pmatrix} \cosh \phi - \beta \sinh \phi \\ \sinh \phi - \beta \cosh \phi \end{pmatrix}$$

# SR: worldsheet space contraction



$$\Lambda^{-1} \begin{pmatrix} 1 \\ -\beta \end{pmatrix} = \begin{pmatrix} \cosh \phi (1 - \beta^2) \\ \cosh \phi (\tanh \phi - \tanh \phi) \end{pmatrix}$$

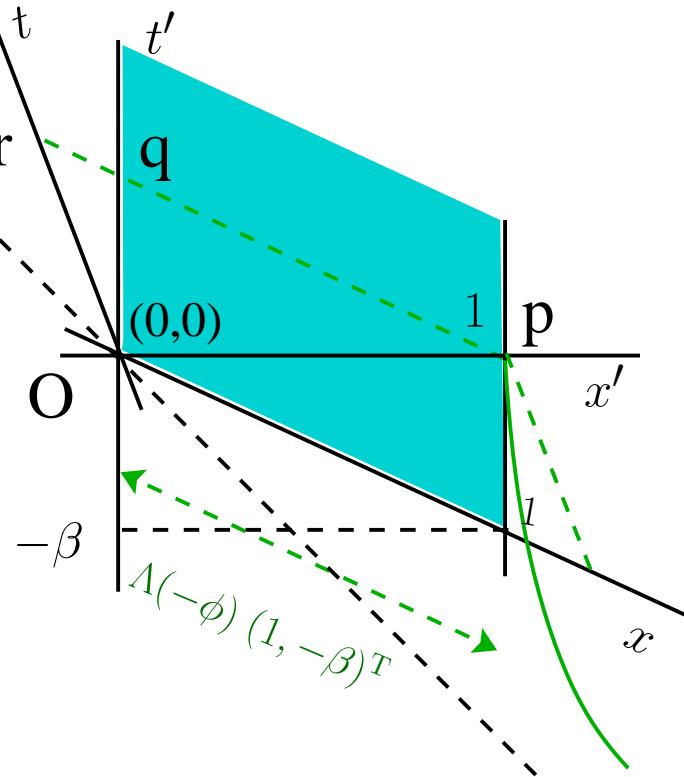
# SR: worldsheet space contraction



$$\Lambda^{-1} \begin{pmatrix} 1 \\ -\beta \end{pmatrix} = \begin{pmatrix} \gamma(1 - \beta^2) \\ 0 \end{pmatrix}$$

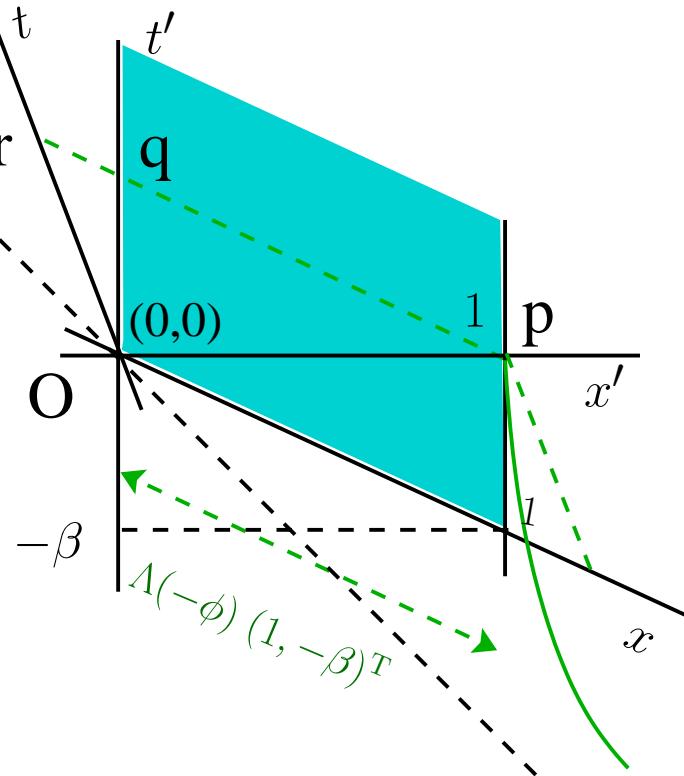


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$$\Lambda^{-1} \begin{pmatrix} 1 \\ -\beta \end{pmatrix} = \begin{pmatrix} \gamma^{-1} \\ 0 \end{pmatrix}$$

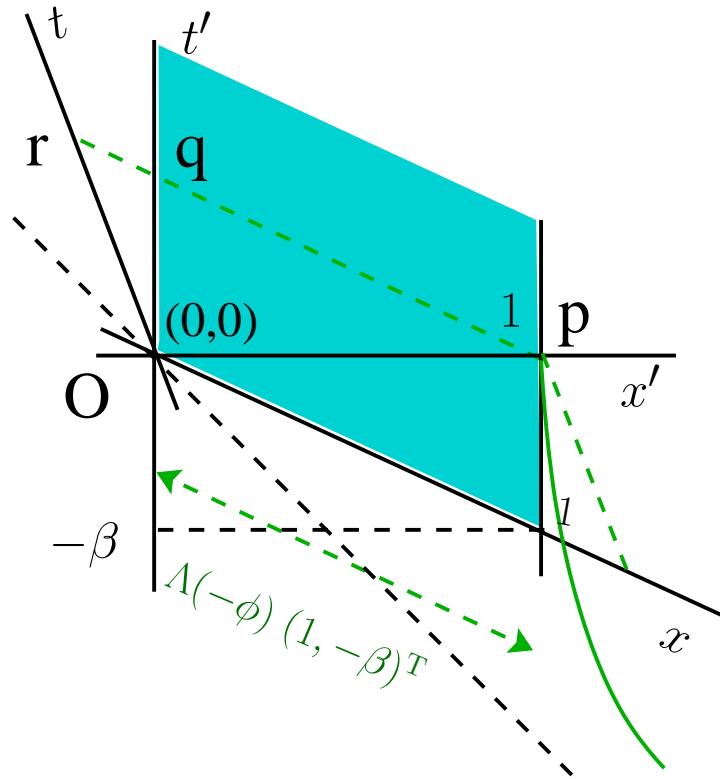
# SR: worldsheet space contraction



$$\sqrt{\Delta s^2(q, p)} = \gamma^{-1} < 1$$



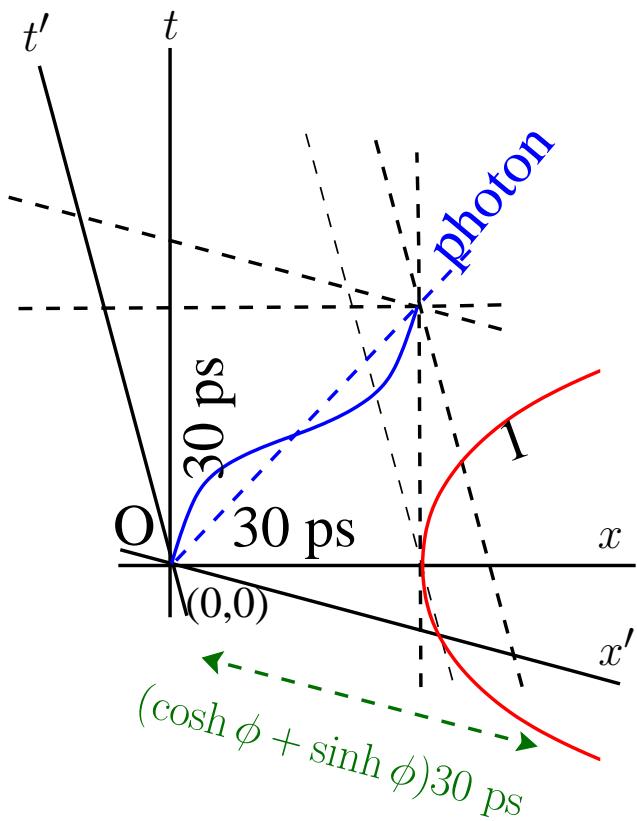
# SR: worldsheet space contraction



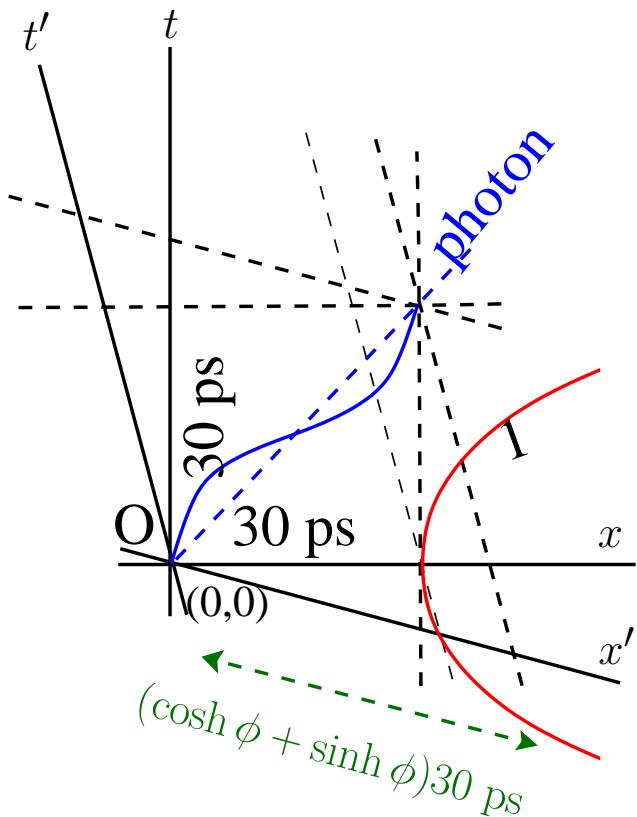
$$\sqrt{\Delta s^2(q, p)} = \gamma^{-1} < 1 \quad \text{worldsheet "space contraction"}$$



# SR: Doppler shift

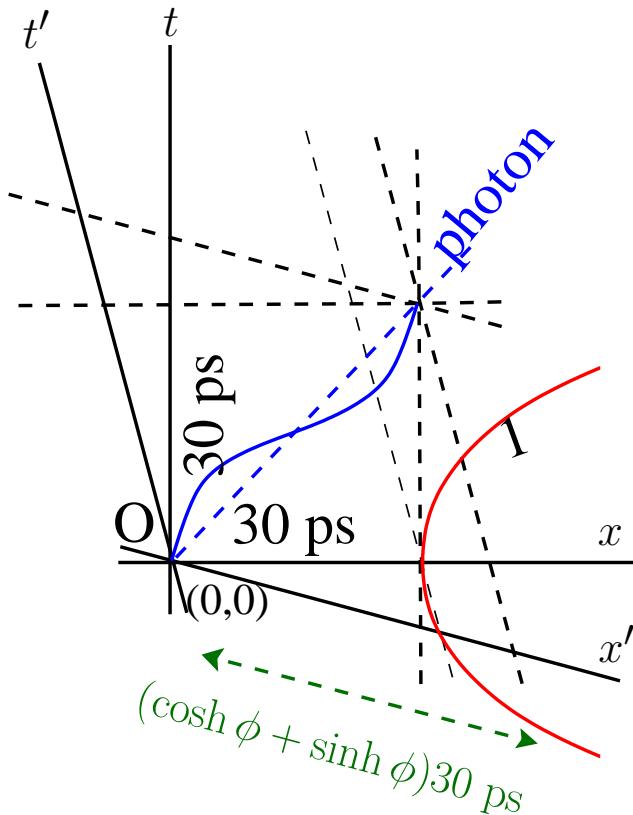


# SR: Doppler shift



see photon worldline calculation

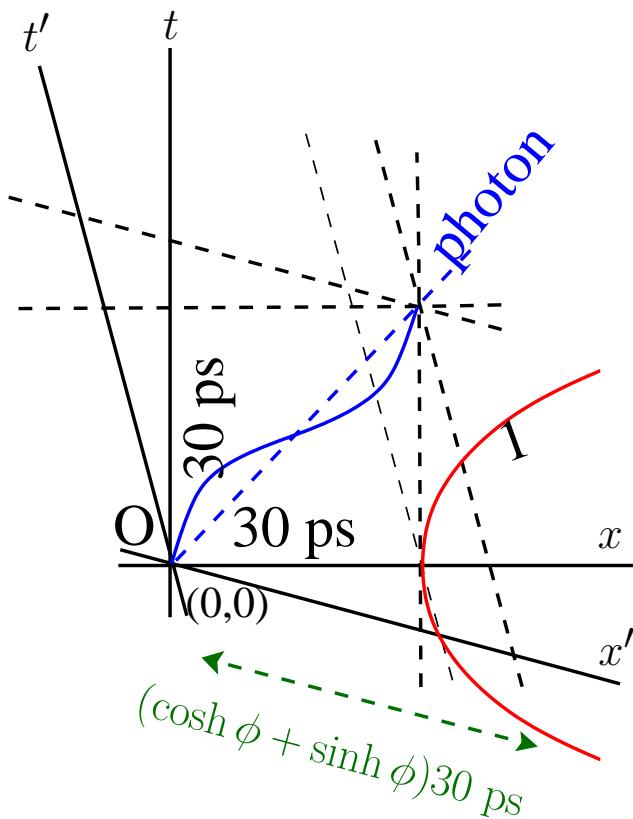
# SR: Doppler shift



see photon worldline calculation

$$x' = (\cosh \phi + \sinh \phi)t$$

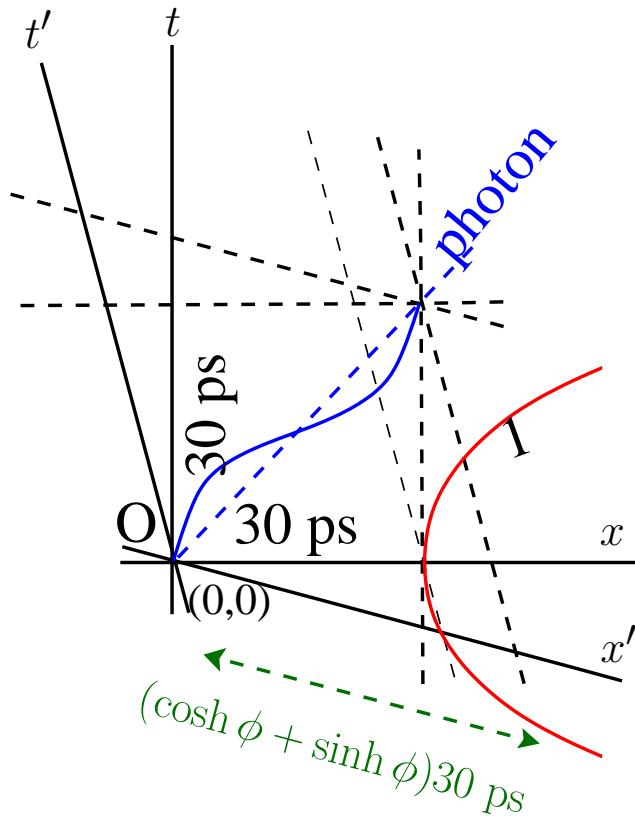
# SR: Doppler shift



see photon worldline calculation

$$x' = (\cosh \phi + \sinh \phi)x$$

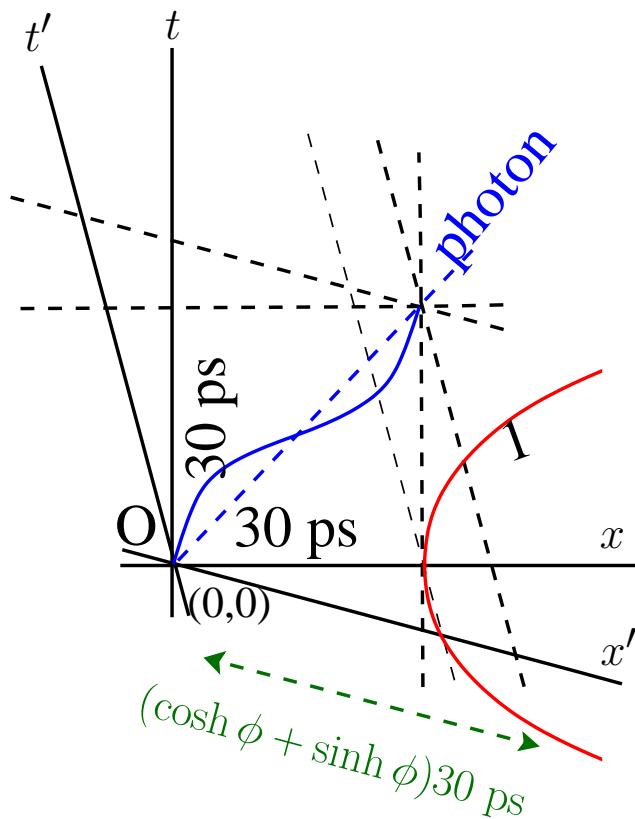
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \cosh \phi + \sinh \phi$$

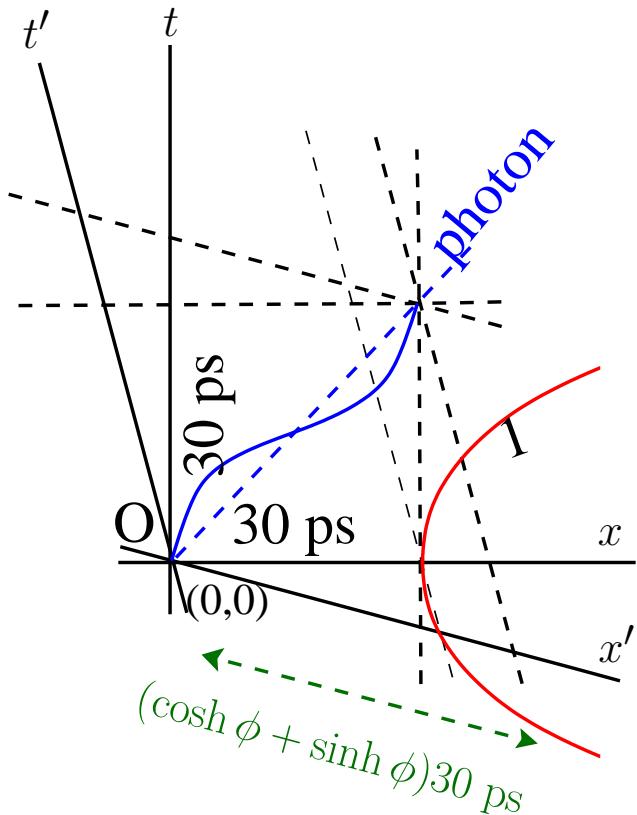
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \gamma + \beta\gamma$$

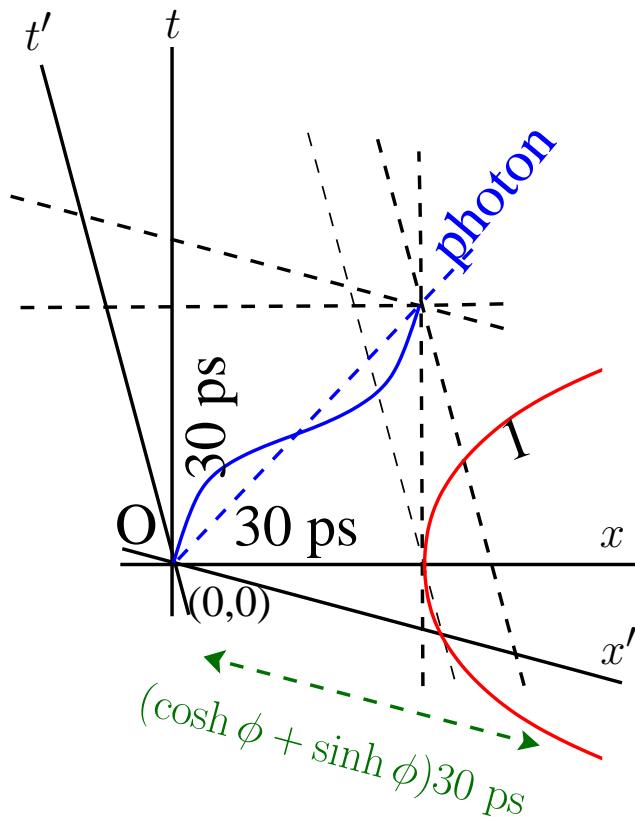
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \gamma(1 + \beta)$$

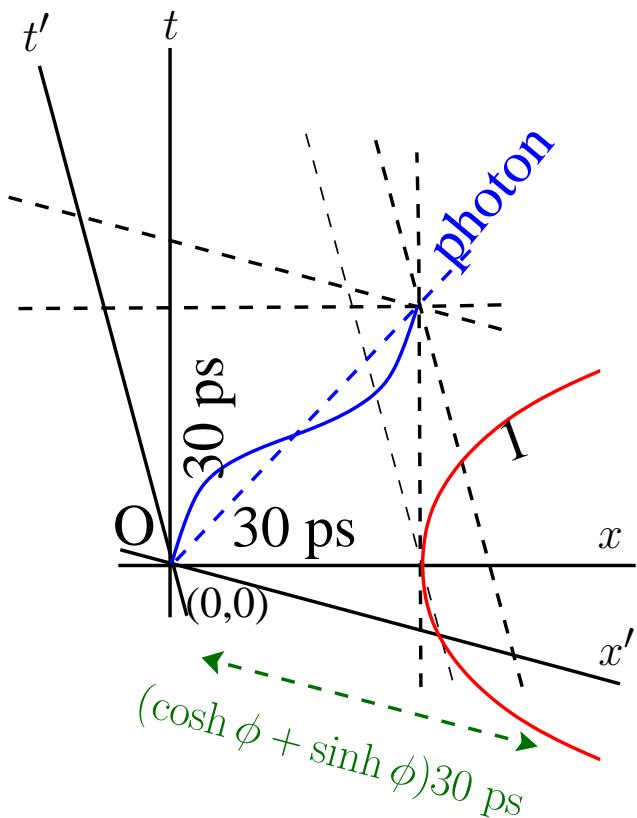
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \gamma(1 + \beta) = \frac{1+\beta}{\sqrt{1-\beta^2}}$$

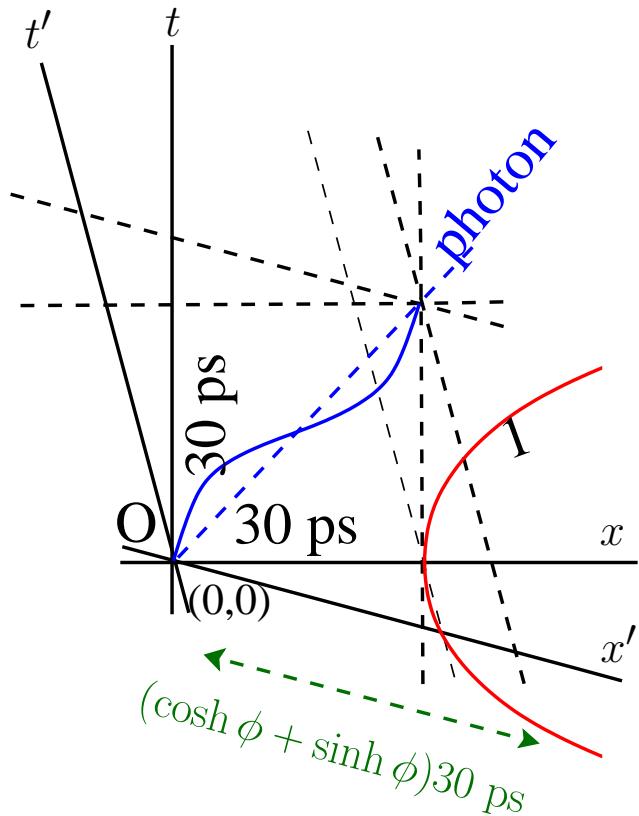
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \gamma(1 + \beta) = \frac{\sqrt{(1+\beta)^2}}{\sqrt{(1-\beta)(1+\beta)}}$$

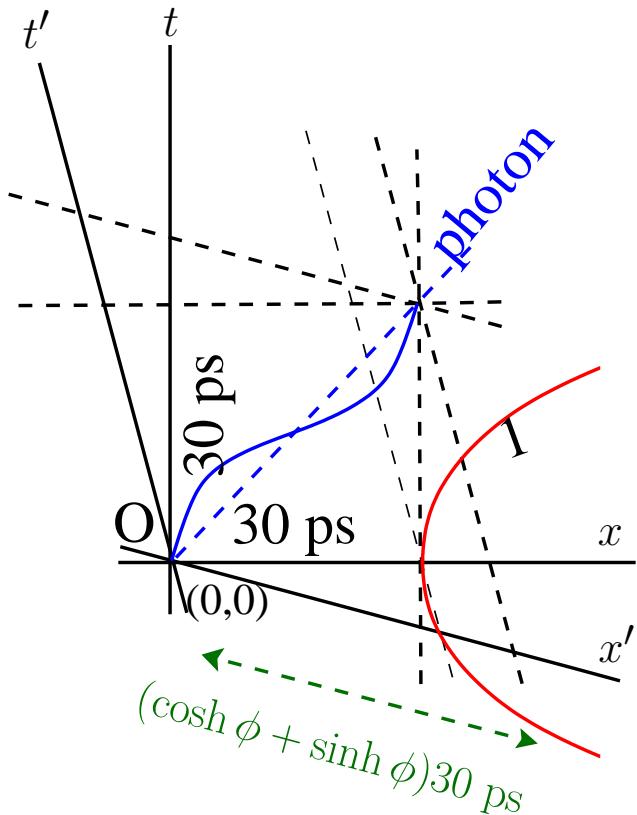
# SR: Doppler shift



see photon worldline calculation

$$x'/x = \gamma(1 + \beta) = \sqrt{\frac{1+\beta}{1-\beta}}$$

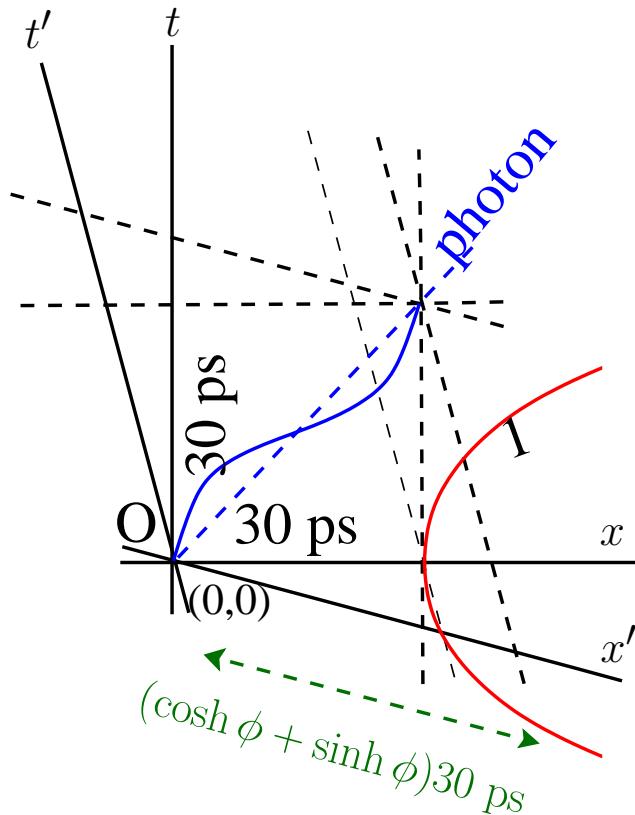
# SR: Doppler shift



see photon worldline calculation

$$1 + z := \lambda'/\lambda = \gamma(1 + \beta) \equiv \sqrt{\frac{1+\beta}{1-\beta}}$$
redshift

# SR: Doppler shift



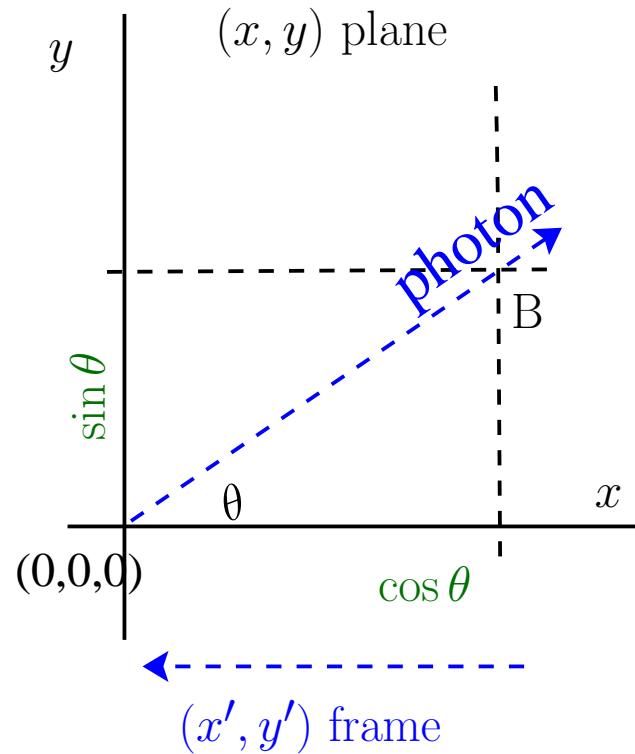
see photon worldline calculation

$$1 + z := \lambda'/\lambda = \gamma(1 + \beta) \equiv \sqrt{\frac{1+\beta}{1-\beta}} \quad \text{redshift}$$

$\Rightarrow$  when  $\beta \ll 1$ ,  $z \approx \beta$

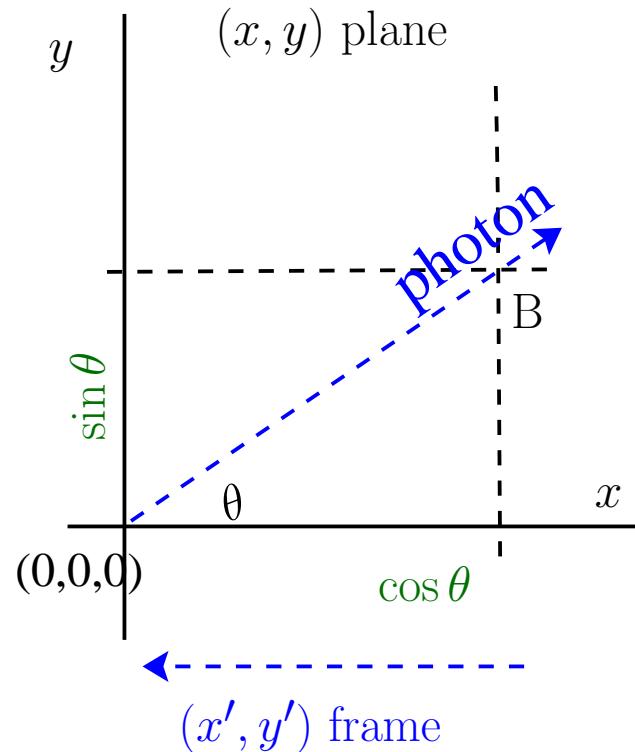


# SR: relativistic aberration





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event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$





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# SR: relativistic aberration

event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$

$$\Lambda^{-1} \begin{pmatrix} \cos \theta \\ \sin \theta \\ 1 \end{pmatrix} = \begin{pmatrix} \gamma & 0 & \beta\gamma \\ 0 & 1 & 0 \\ \beta\gamma & 0 & \gamma \end{pmatrix} \begin{pmatrix} \cos \theta \\ \sin \theta \\ 1 \end{pmatrix}$$





# SR: relativistic aberration

event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$

$$\Lambda^{-1} \begin{pmatrix} \cos \theta \\ \sin \theta \\ 1 \end{pmatrix} = \begin{pmatrix} \gamma \cos \theta + \beta \gamma \\ \sin \theta \\ \beta \gamma \cos \theta + \gamma \end{pmatrix}$$





# SR: relativistic aberration

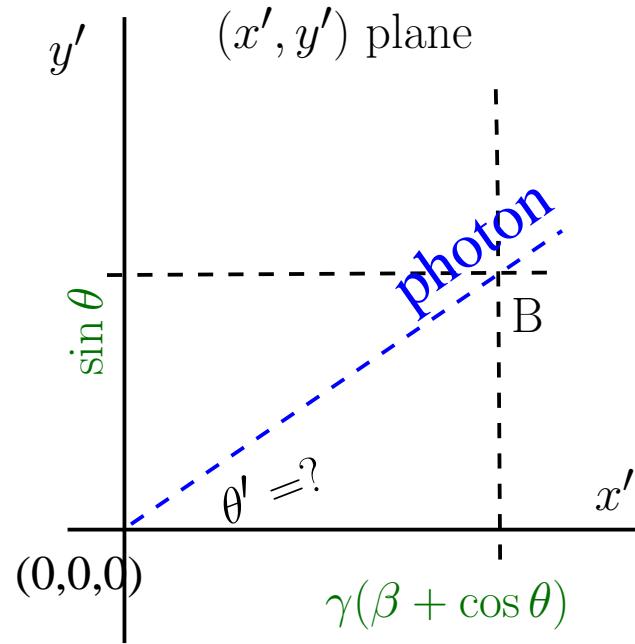
event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$

$$\Lambda^{-1} \begin{pmatrix} \cos \theta \\ \sin \theta \\ 1 \end{pmatrix} = \begin{pmatrix} \gamma(\cos \theta + \beta) \\ \sin \theta \\ \gamma(1 + \beta \cos \theta) \end{pmatrix}$$



# SR: relativistic aberration

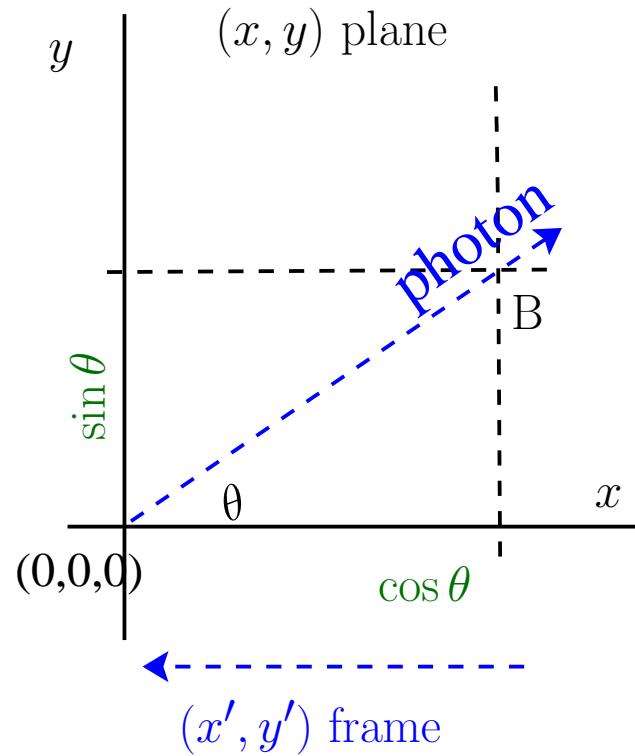
event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$





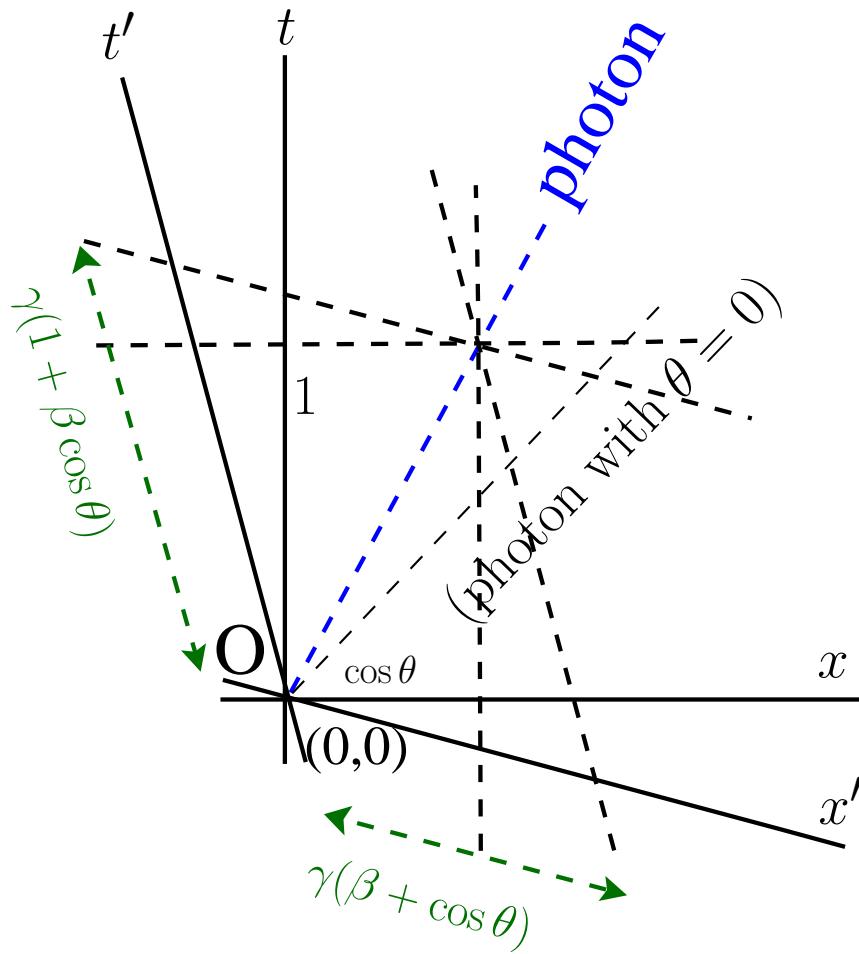
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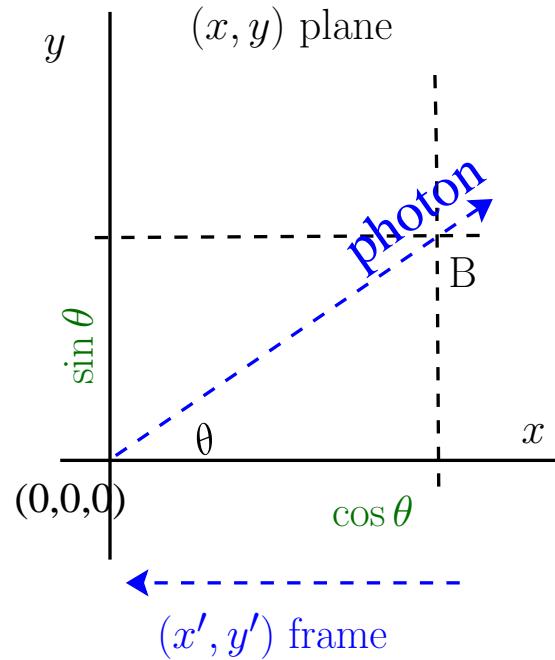
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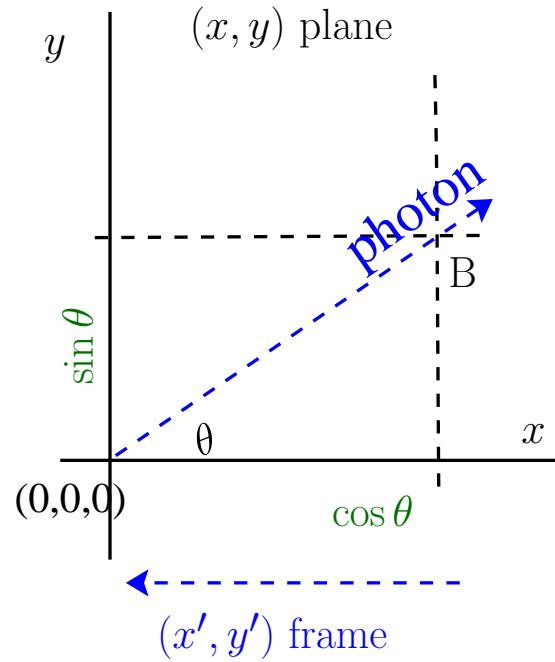
$$\tan \theta' = \frac{\sin \theta}{\gamma(\beta + \cos \theta)}$$





# SR: relativistic aberration

event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$



$$\tan \theta' = \frac{\sin \theta}{\gamma(\beta + \cos \theta)} < \tan \theta \text{ if } 0 < \beta < 1$$

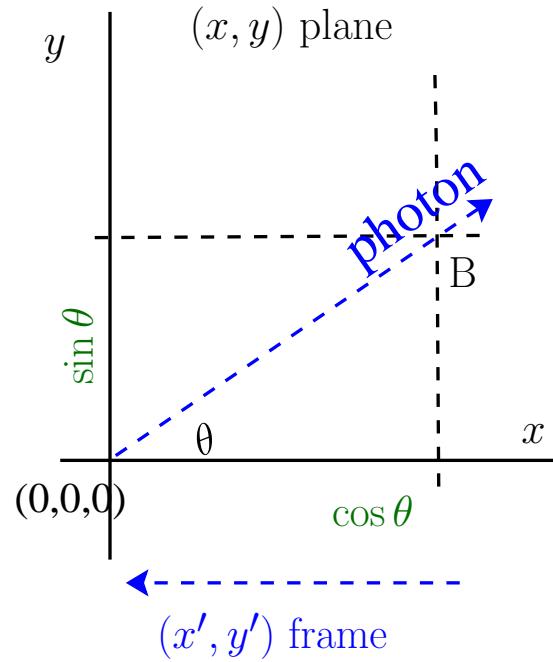
w:Relativistic aberration (2011-02-22: quality=weak)





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event B:  $(x, y, t) = (\cos \theta, \sin \theta, 1)$



$$\tan \theta' = \frac{\sin \theta}{\gamma(\beta + \cos \theta)} < \tan \theta \text{ if } 0 < \beta < 1$$

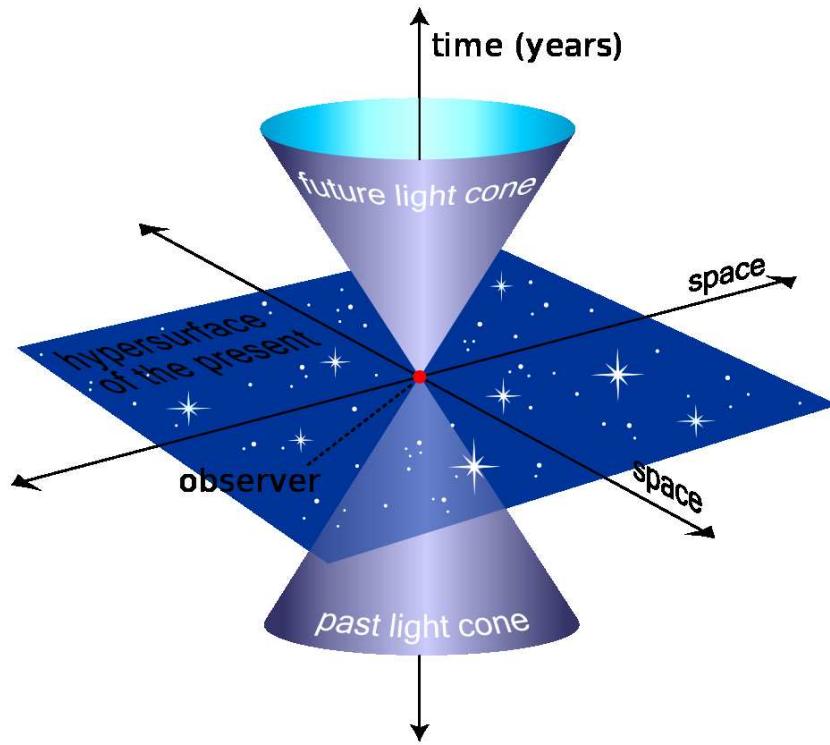
w:Relativistic aberration (2011-02-22: quality=weak)

⇒ relativistic beaming, e.g. AGN jets



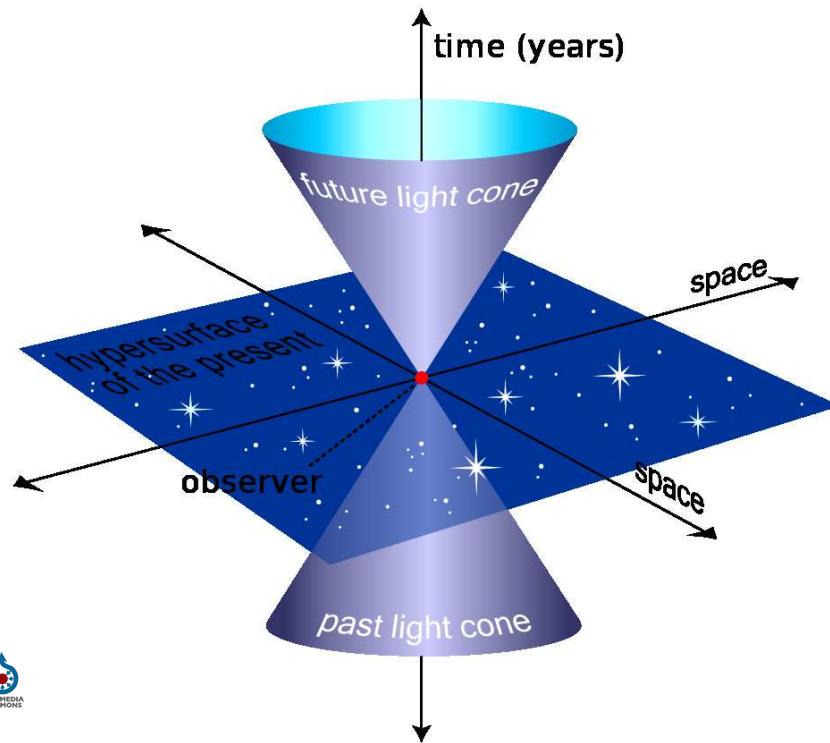


# SR: world line





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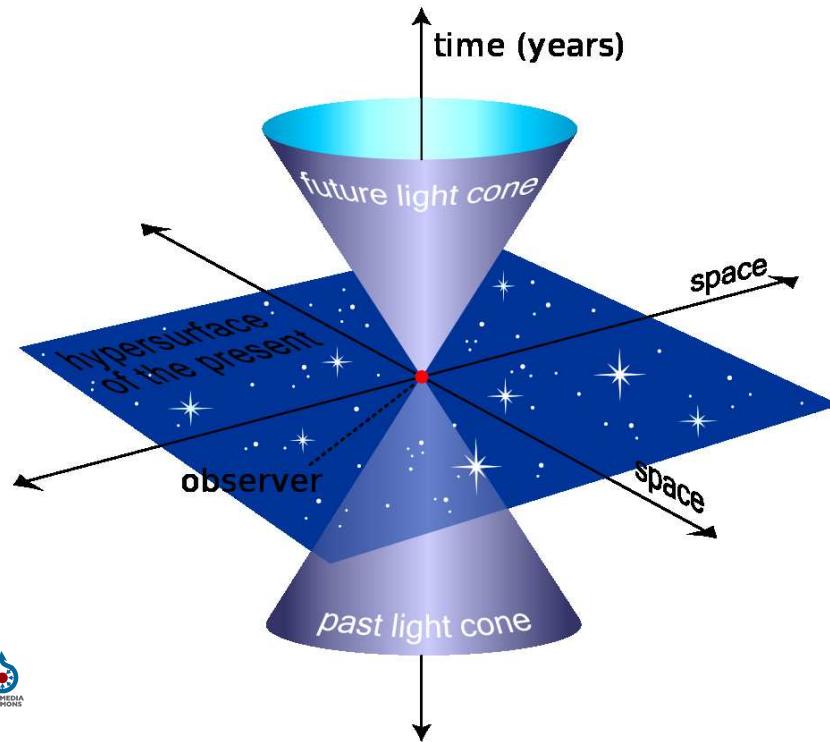


lightlike interval = null interval:  $(\Delta s)^2 = 0$   
spacetime =





# SR: world line



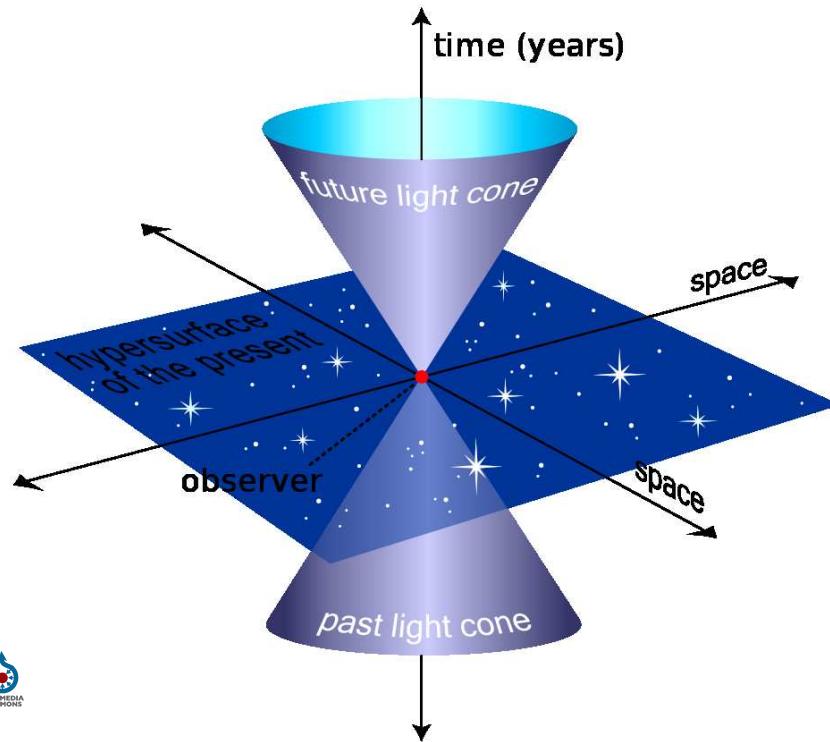
lightlike interval = null interval:  $(\Delta s)^2 = 0$

spacetime = on past w:light cone + inside past light cone





# SR: world line



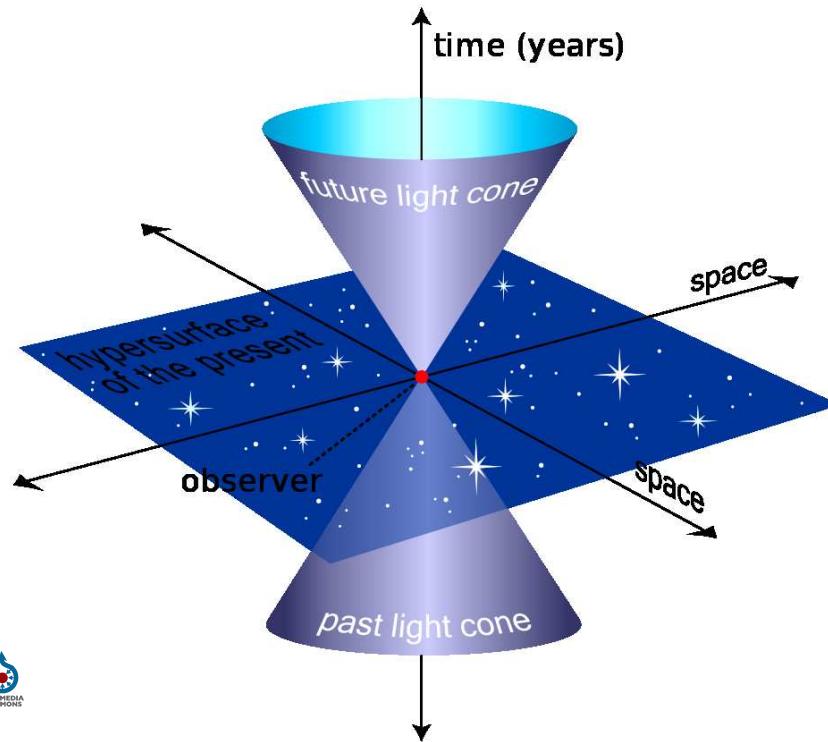
lightlike interval = null interval:  $(\Delta s)^2 = 0$

spacetime = on past w:light cone + inside past light cone  
+ on future light cone + inside future light cone





# SR: world line



lightlike interval = null interval:  $(\Delta s)^2 = 0$

spacetime = on past w:light cone + inside past light cone  
+ on future light cone + inside future light cone  
+ elsewhere





# SR: world line

Lorentz transform of world line





# SR: world line

## Lorentz transform of world line





# SR: world line

Lorentz transform of world line



- coordinate time in spacetime model  $\neq$  time in your brain (thinking)





# SR: world line

Lorentz transform of world line



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- $\frac{dt}{dt_{\text{thinking}}}$  can be positive or negative





# SR: world line

Lorentz transform of world line



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Lorentz transform of world line



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- “elsewhere” spacetime events can change from past to future even though  $\frac{dt}{d\lambda} > 0$





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Lorentz transform of world line



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# SR: world line

Lorentz transform of world line

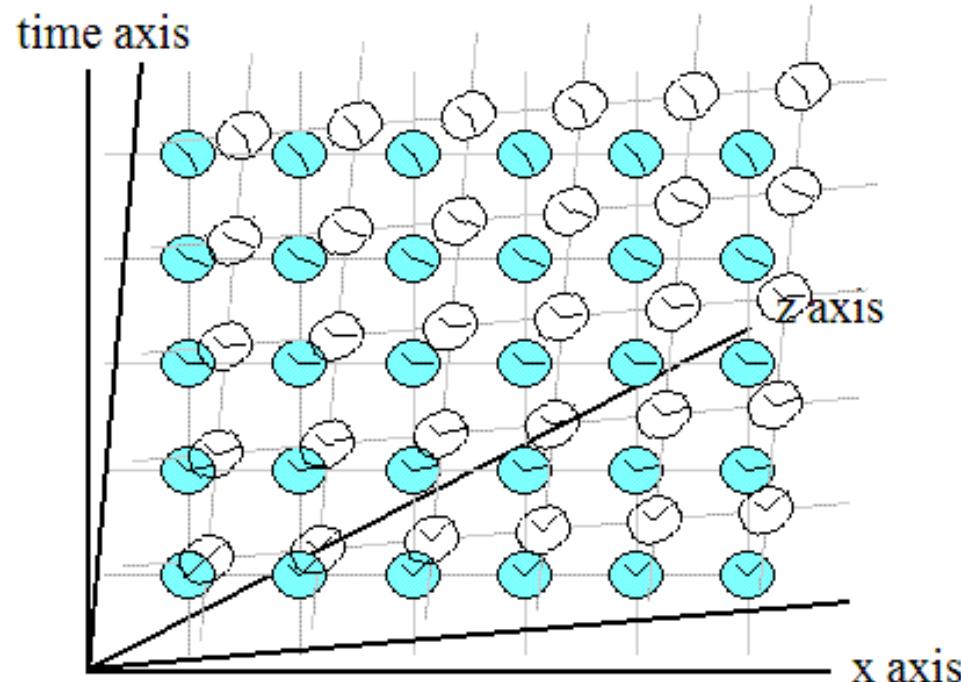


- coordinate time in spacetime model  $\neq$  time in your brain (thinking)
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- “elsewhere” spacetime events can change from past to future even though  $\frac{dt}{d\lambda} > 0$
- w:proper time  $\tau :=$  time along a worldline measured by clock following that worldline
- often  $d\tau$  is useful for integrating





# SR: Rietdijk–Putnam–Penrose p.

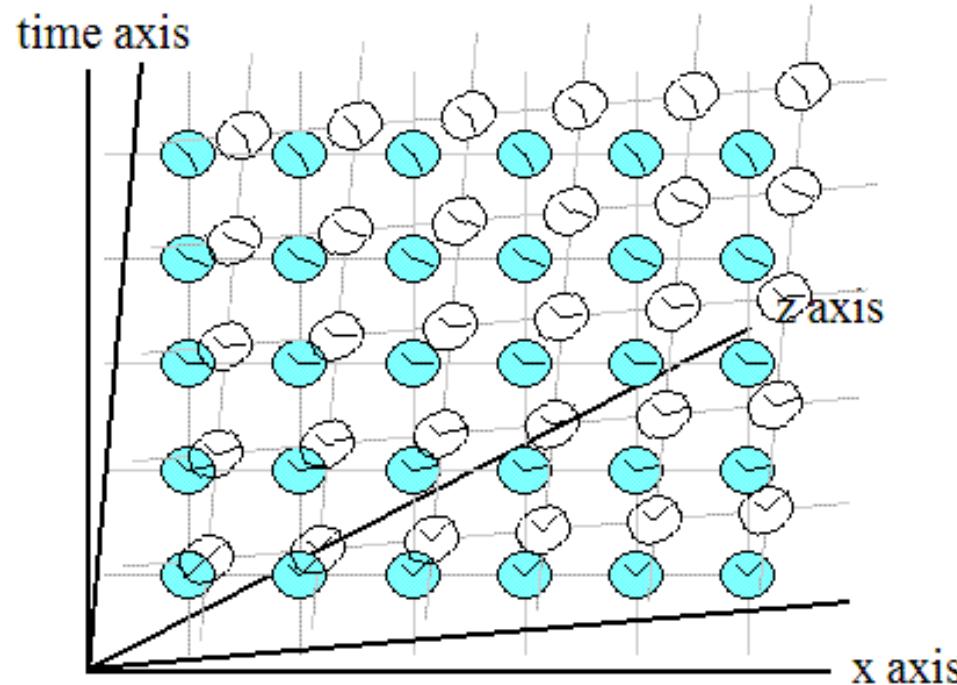


Relativity shows that the inertial frames of reference of relatively moving objects do not overlap each other.





# SR: Rietdijk–Putnam–Penrose p.



Relativity shows that the inertial frames of reference of relatively moving objects do not overlap each other.

b:Inertialoverlay.GIF

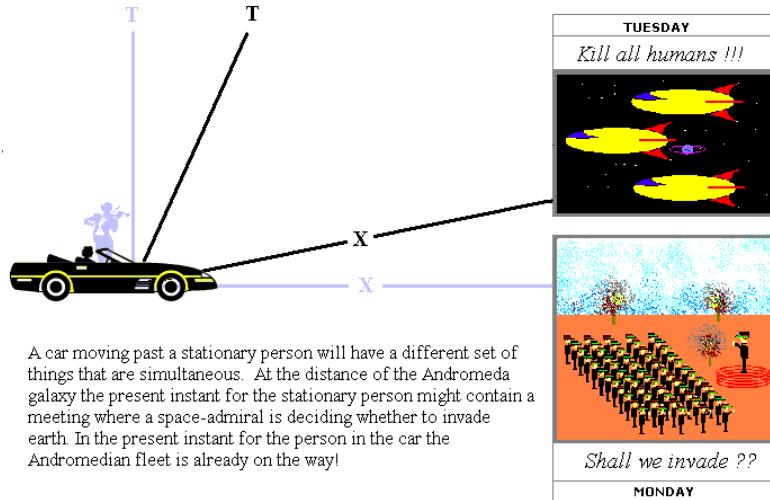
- each observer can synchronise clocks + rods



# SR: Rietdijk–Putnam–Penrose p.



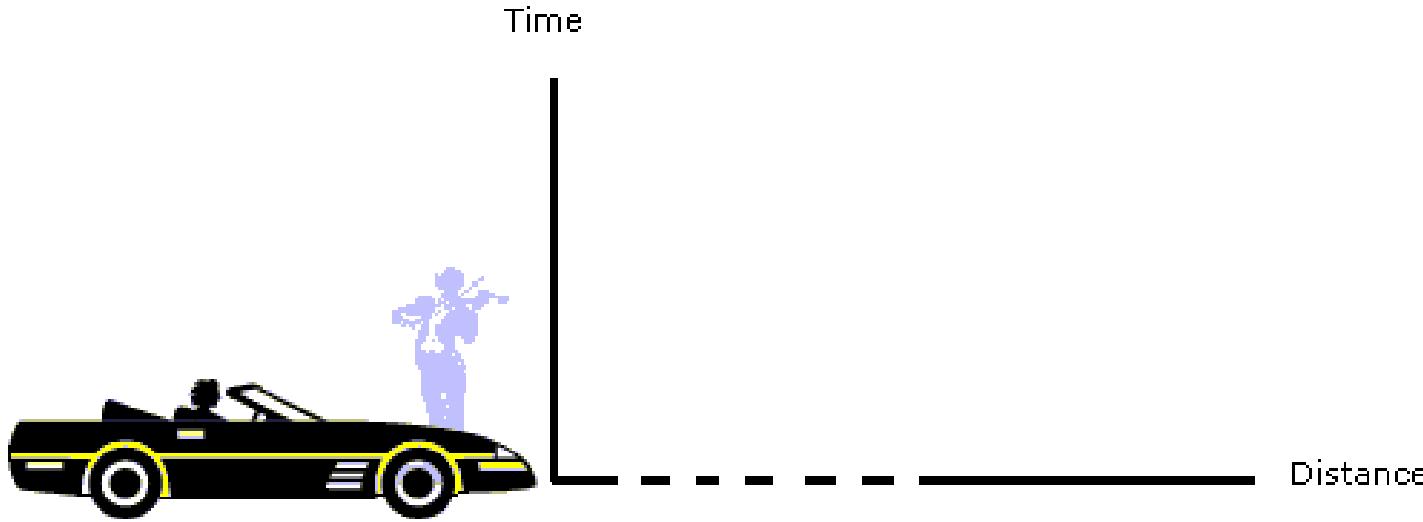
## The Andromeda Paradox



w:Rietdijk-Putnam argument b:Rel2.gif



# SR: Rietdijk–Putnam–Penrose p.

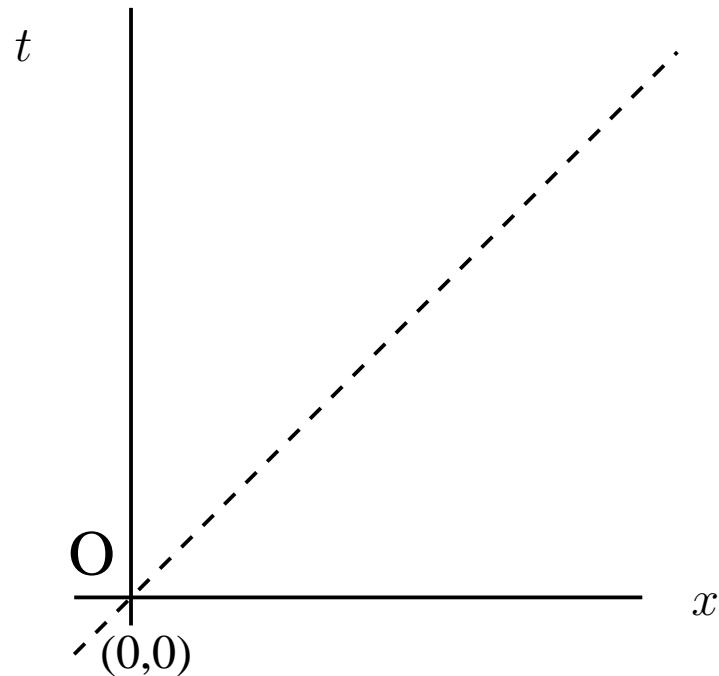


[w:Rietdijk–Putnam argument](#) [b:Rel3.gif](#)





# SR: tachyons and causality

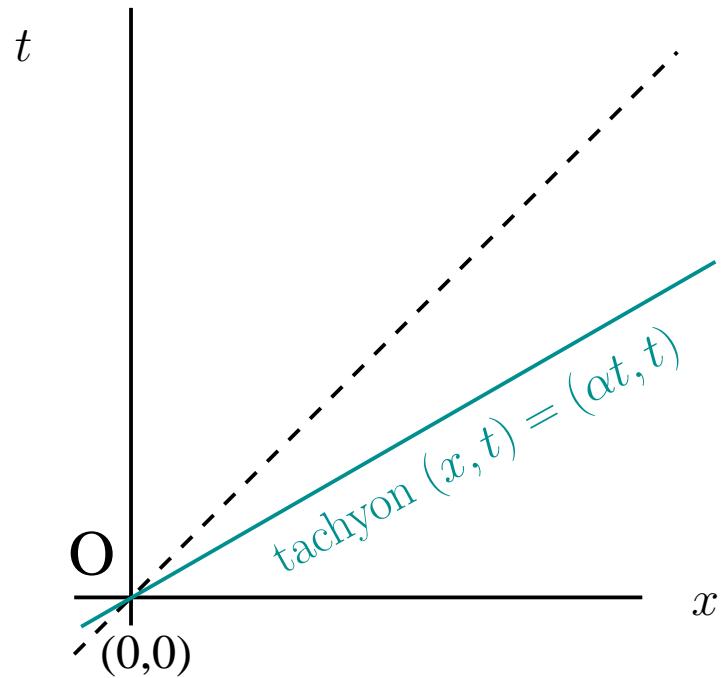


observer “at rest”





# SR: tachyons and causality

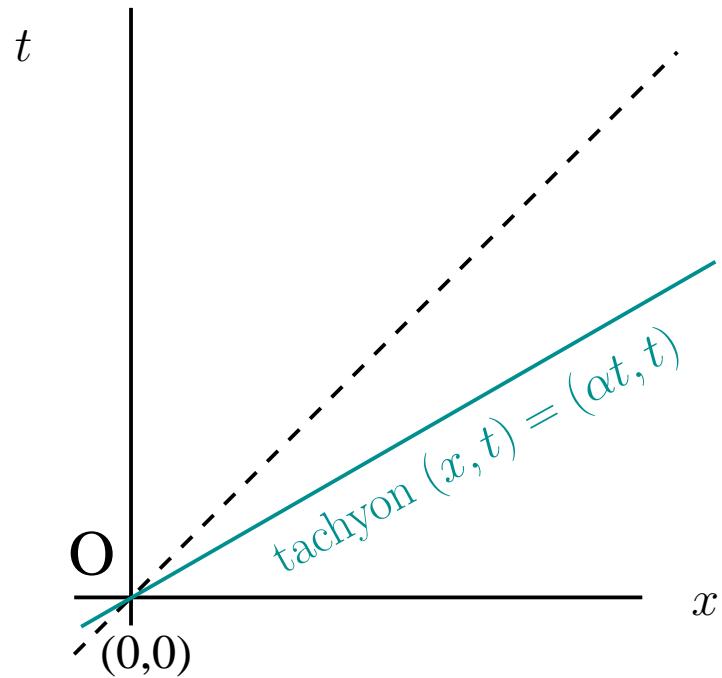


add a tachyon with speed  $\alpha > 1$





# SR: tachyons and causality



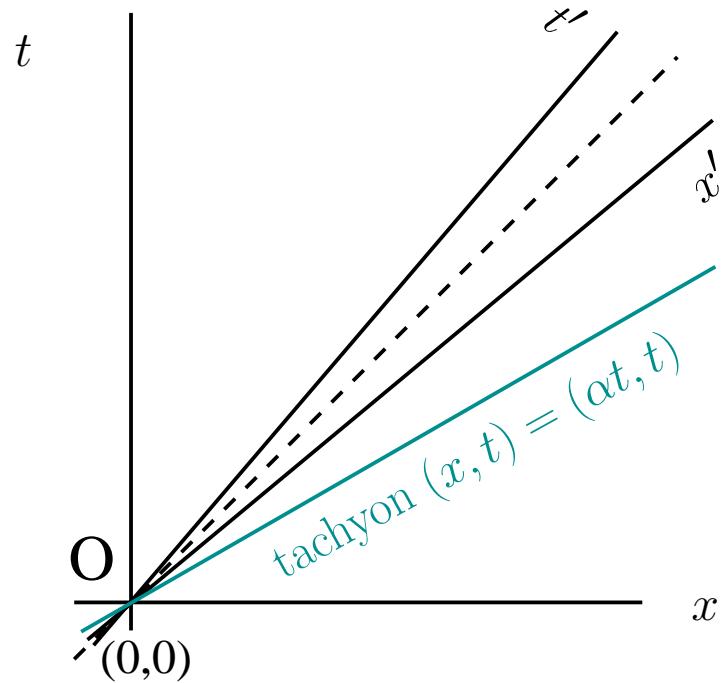
add a tachyon with speed  $\alpha > 1$

choose rocket at speed  $\beta$  with  $1/\alpha < \beta < 1$





# SR: tachyons and causality



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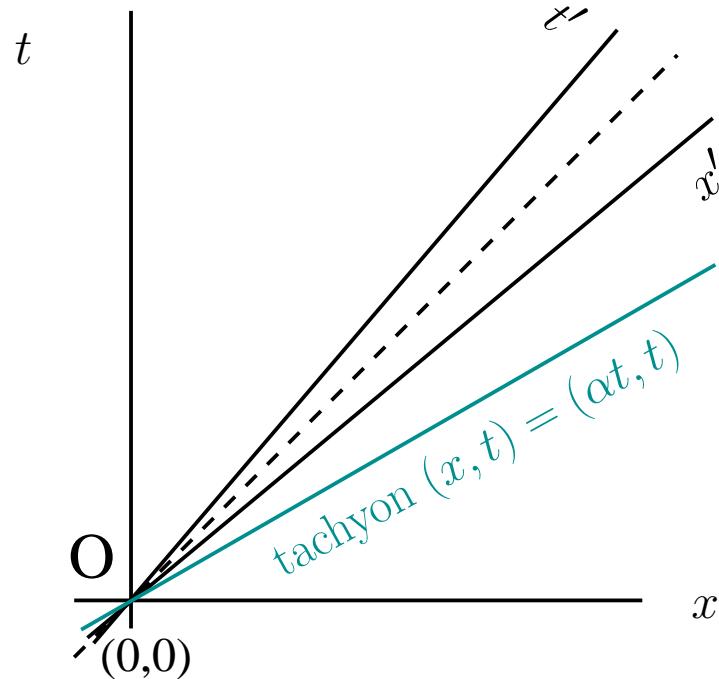
choose rocket at speed  $\beta$  with  $1/\alpha < \beta < 1$

add axes  $x', t'$  for the rocket





# SR: tachyons and causality



add a tachyon with speed  $\alpha > 1$

choose rocket at speed  $\beta$  with  $1/\alpha < \beta < 1$

add axes  $x'$ ,  $t'$  for the rocket

rocket frame:  $(\alpha t, t)$  becomes  $\Lambda (\alpha t, t)^T$





# SR: tachyons and causality

$$\begin{pmatrix} x' \\ t' \end{pmatrix} = \Lambda \begin{pmatrix} \alpha t \\ t \end{pmatrix} =$$





# SR: tachyons and causality

$$\begin{pmatrix} x' \\ t' \end{pmatrix} = \Lambda \begin{pmatrix} \alpha t \\ t \end{pmatrix} = \begin{pmatrix} \gamma\alpha t - \beta\gamma t \\ -\alpha\beta\gamma t + \gamma t \end{pmatrix}$$





# SR: tachyons and causality

$$\begin{pmatrix} x' \\ t' \end{pmatrix} = \Lambda \begin{pmatrix} \alpha t \\ t \end{pmatrix} = \gamma t \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$





# SR: tachyons and causality

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$$x' = \gamma t(\alpha - \beta) > 0 \text{ since } \alpha > 1 > \beta$$





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$dt'/dt < 0$

same sequence of spacetime events = tachyon worldline:

$t$  increases for observer “at rest”,

$t'$  decreases for rocket observer (with  $\beta > 1/\alpha$ )





# SR: tachyons and causality

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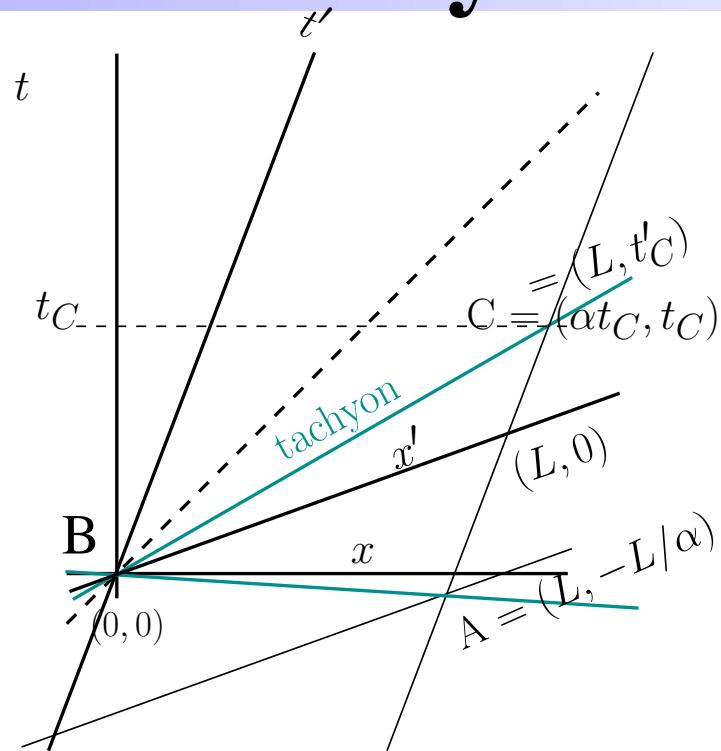
$t'$  decreases for rocket observer (with  $\beta > 1/\alpha$ )

- observer at rest: tachyon emitted at origin
- rocket: tachyon absorbed at origin



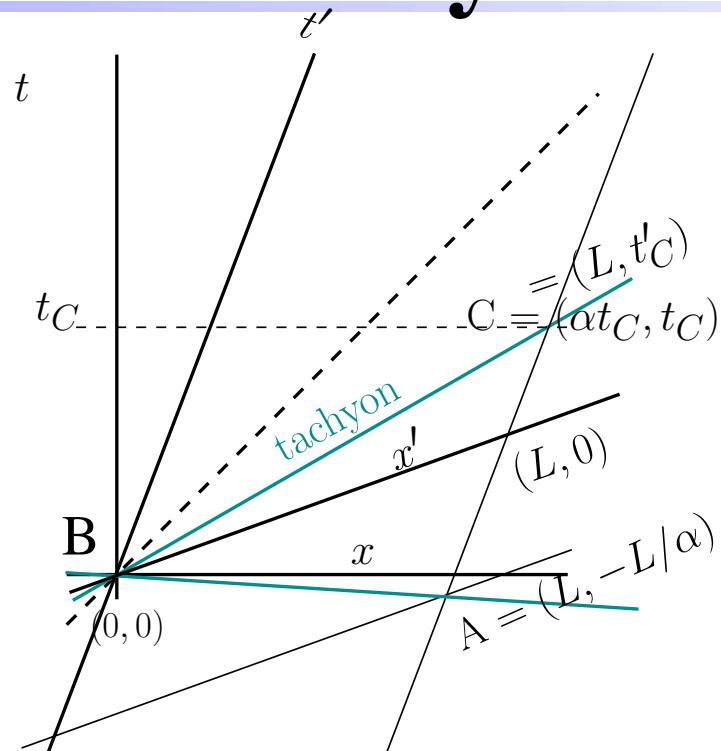


# SR: tachyonic antitelephone





# SR: tachyonic antitelephone

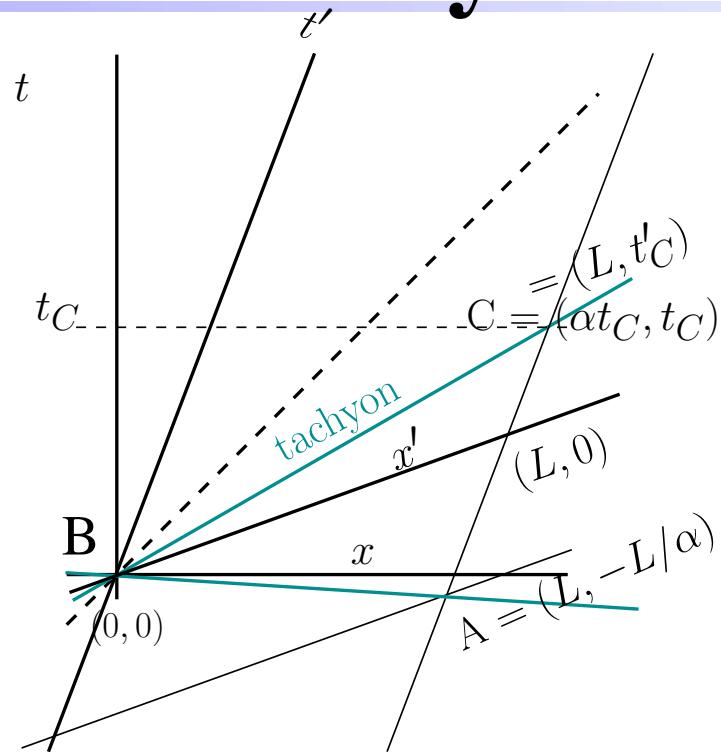


B stationary:  $(x, t)$   
frame





# SR: tachyonic antitelephone



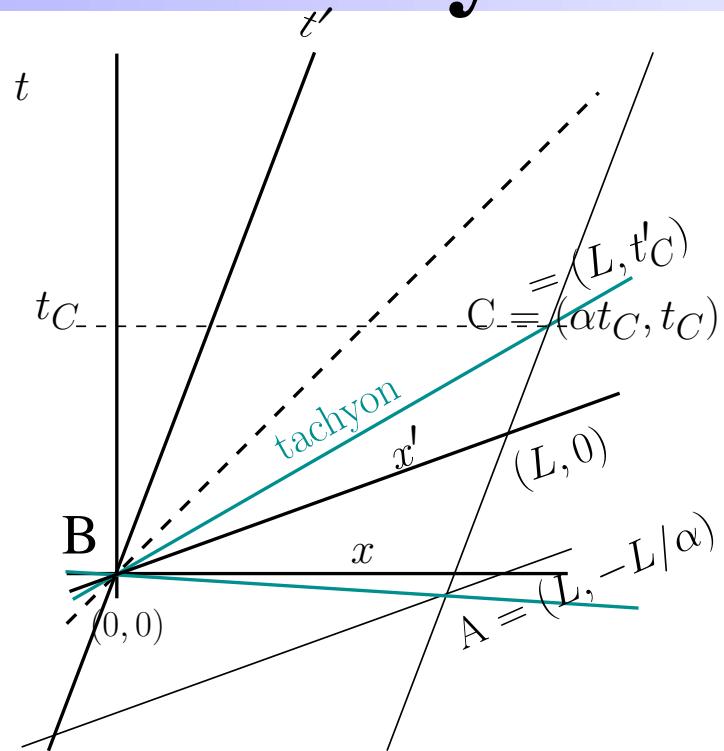
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



A: tachyon at  $\alpha > 1$  to B

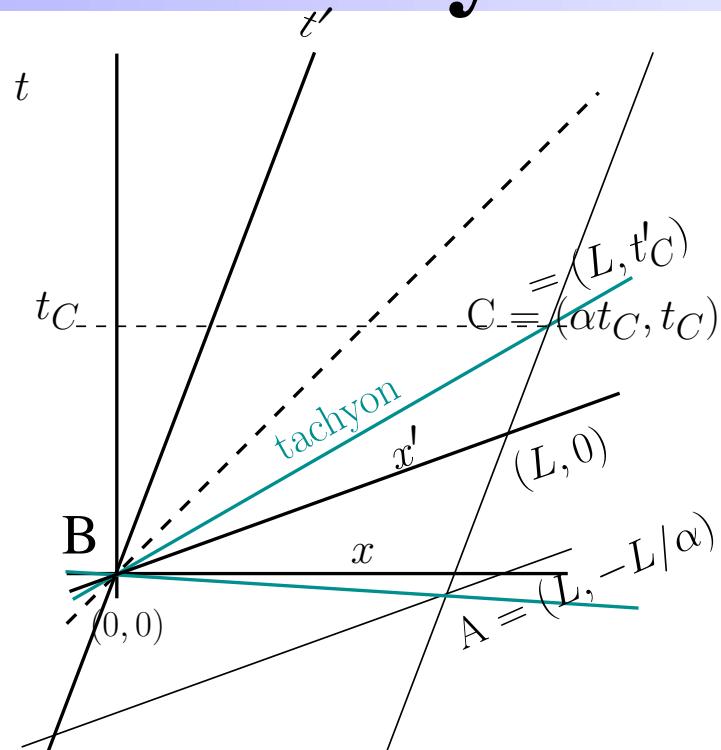
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



B: tachyon at  $\alpha > 1$  to C

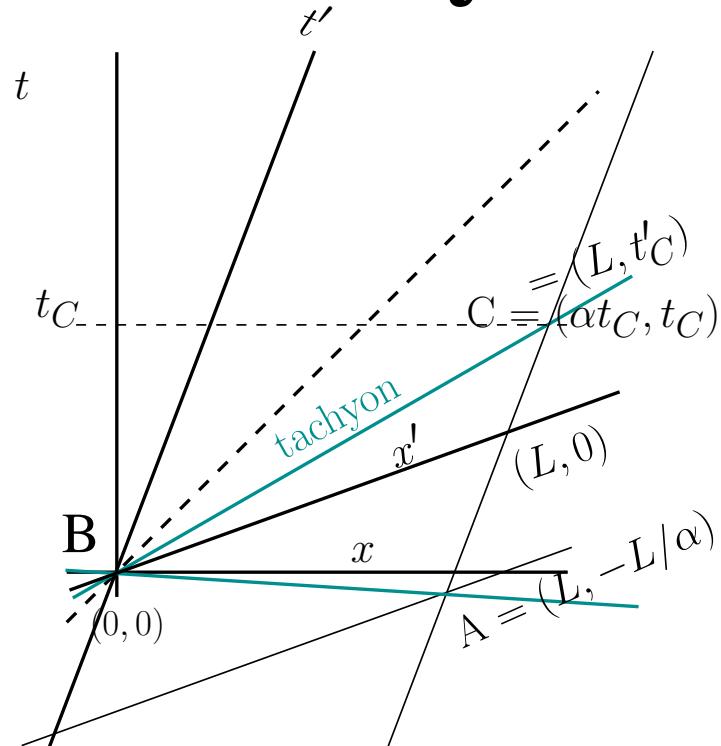
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



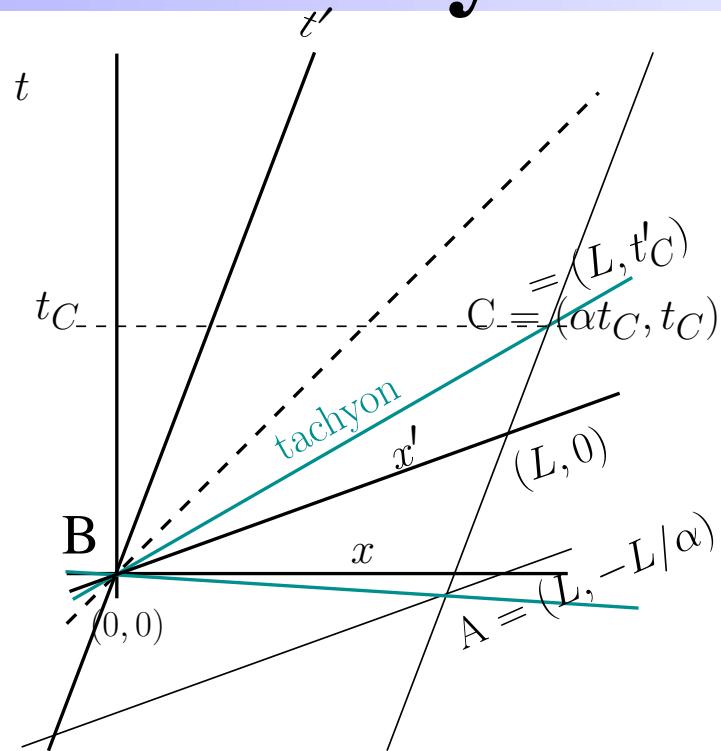
$$C: \begin{pmatrix} L \\ t'_C \end{pmatrix} = \Lambda \begin{pmatrix} \alpha t_C \\ t_C \end{pmatrix}$$

B stationary:  $(x, t)$   
frame  
A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



$$\begin{pmatrix} L \\ t'_C \end{pmatrix} = \gamma t_C \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$

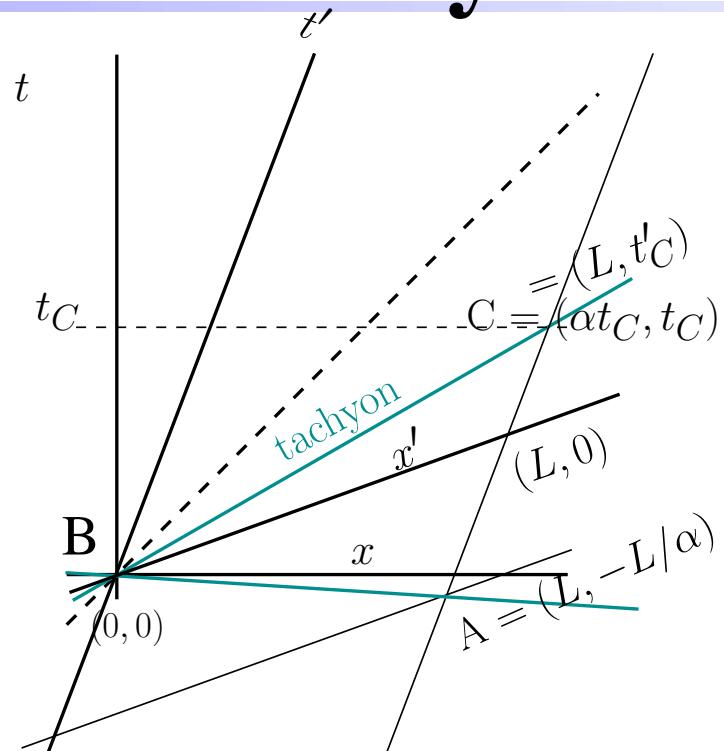
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



$$\begin{pmatrix} L \\ t'_C \end{pmatrix} = \gamma t_C \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$
$$t'_C = \gamma t_C (1 - \alpha\beta)$$

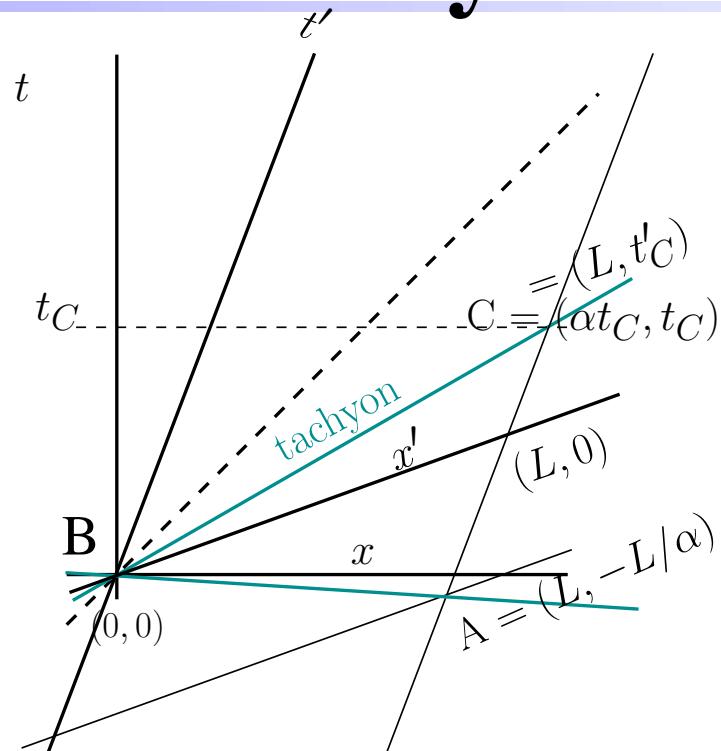
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



$$\begin{pmatrix} L \\ t'_C \end{pmatrix} = \gamma t_C \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$

$$t'_C = \gamma \frac{L}{\gamma(\alpha-\beta)}(1 - \alpha\beta)$$

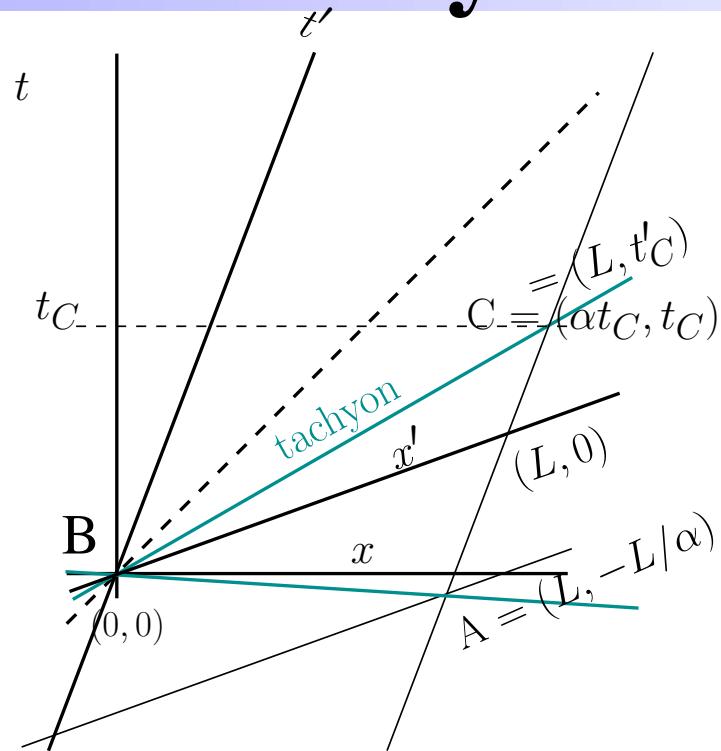
B stationary:  $(x, t)$   
frame

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 $(x', t')$  frame





# SR: tachyonic antitelephone



$$\begin{pmatrix} L \\ t'_C \end{pmatrix} = \gamma t_C \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$

$$t'_C = L \frac{1 - \alpha\beta}{\alpha - \beta}$$

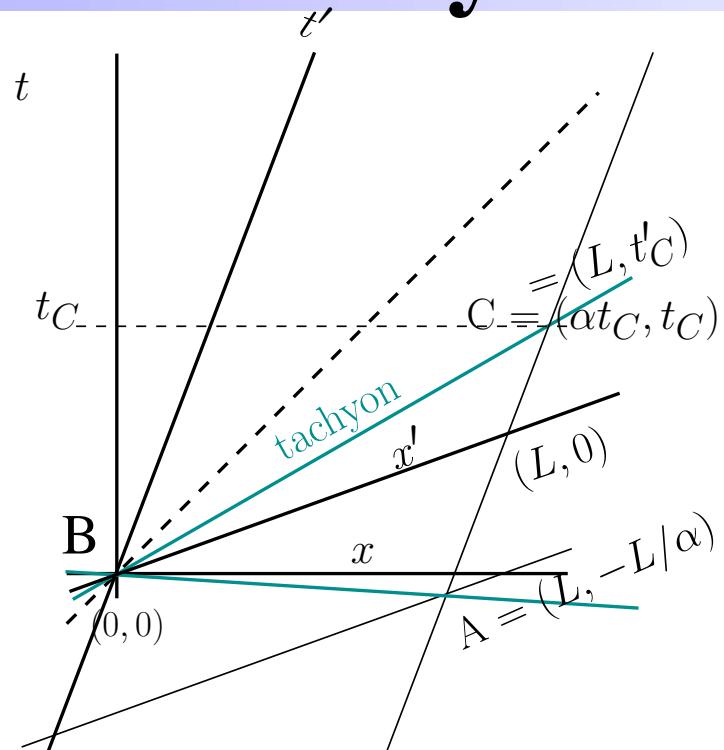
B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame





# SR: tachyonic antitelephone



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frame

A moving at speed  $\beta$ :  
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$$\begin{pmatrix} L \\ t'_C \end{pmatrix} = \gamma t_C \begin{pmatrix} \alpha - \beta \\ -\alpha\beta + 1 \end{pmatrix}$$

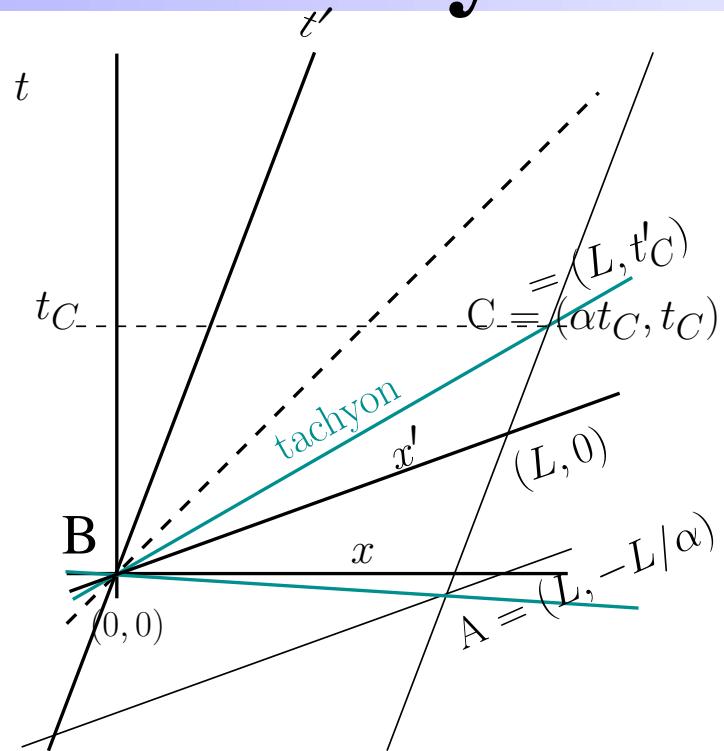
$$t'_C = L \frac{1 - \alpha\beta}{\alpha - \beta}$$

$$t'_C - t'_A = L \left( \frac{1 - \alpha\beta}{\alpha - \beta} + \frac{1}{\alpha} \right)$$





# SR: tachyonic antitelephone



B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame

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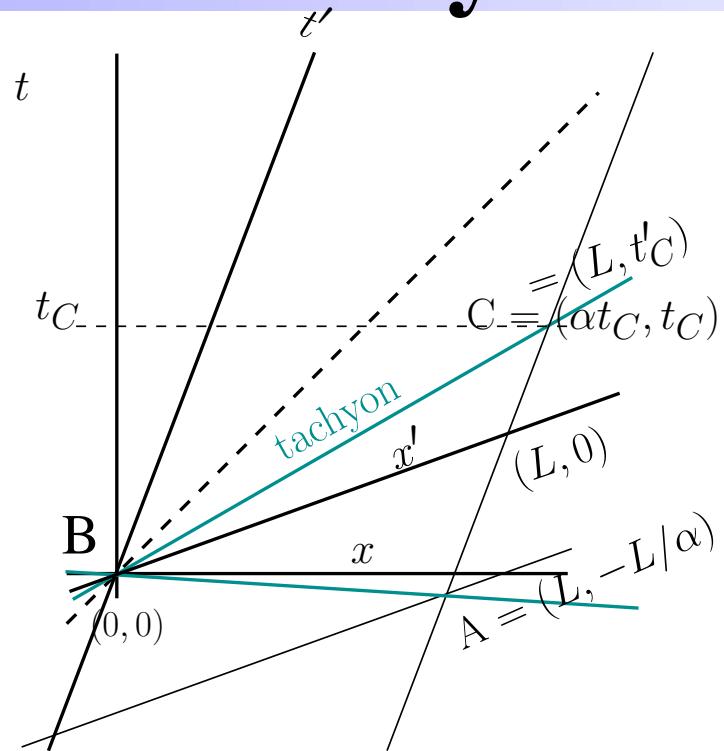
$$t'_C = L \frac{1 - \alpha\beta}{\alpha - \beta}$$

$$t'_C - t'_A = L \frac{\alpha - \alpha^2\beta + \alpha - \beta}{\alpha(\alpha - \beta)}$$





# SR: tachyonic antitelephone



B stationary:  $(x, t)$   
frame

A moving at speed  $\beta$ :  
 $(x', t')$  frame

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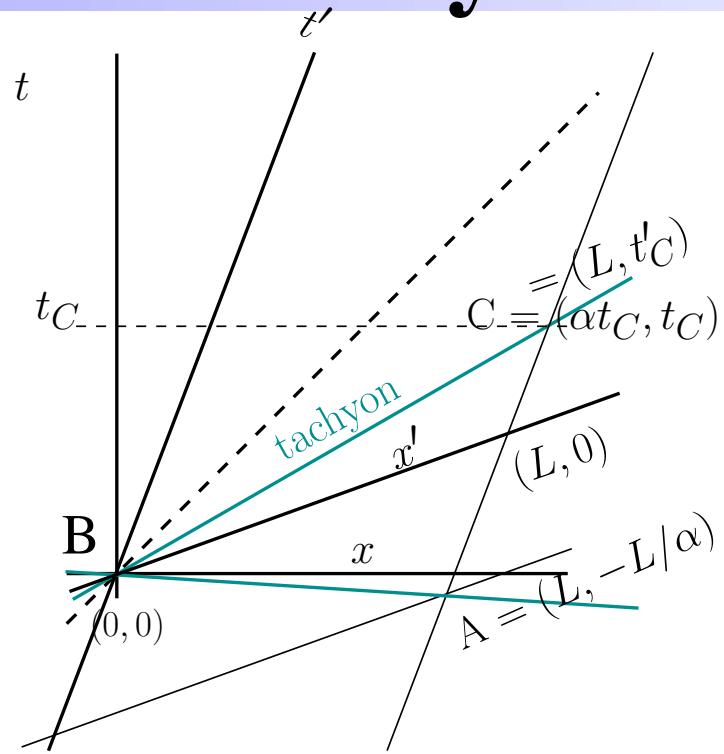
$$t'_C = L \frac{1 - \alpha\beta}{\alpha - \beta}$$

$$t'_C - t'_A = L \frac{2\alpha - (\alpha^2 + 1)\beta}{\alpha(\alpha - \beta)}$$





# SR: tachyonic antitelephone



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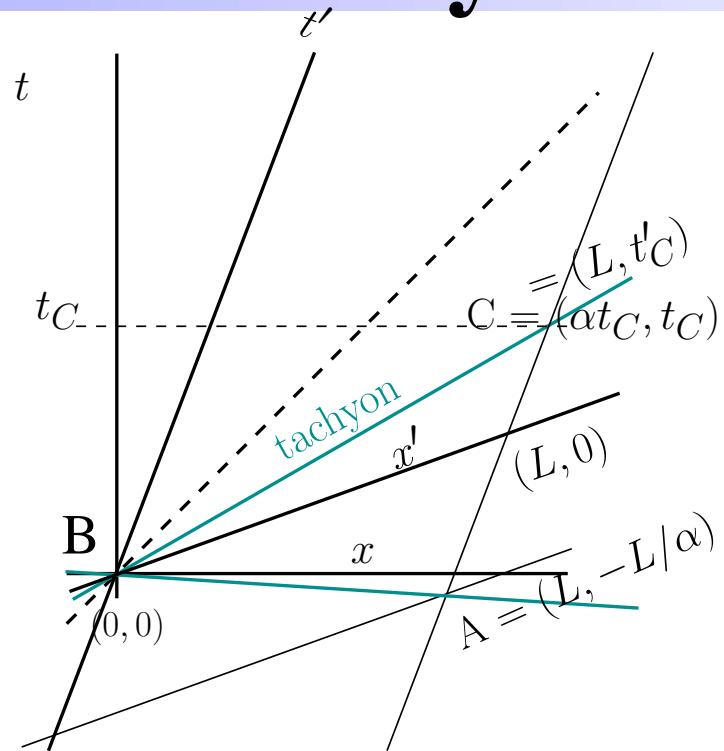
$$t'_C - t'_A = L \frac{2\alpha - (\alpha^2 + 1)\beta}{\alpha(\alpha - \beta)}$$

$$< 0 \text{ if } \beta > \frac{2\alpha}{\alpha^2 + 1}$$





# SR: tachyonic antitelephone



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frame  
A moving at speed  $\beta$ :  
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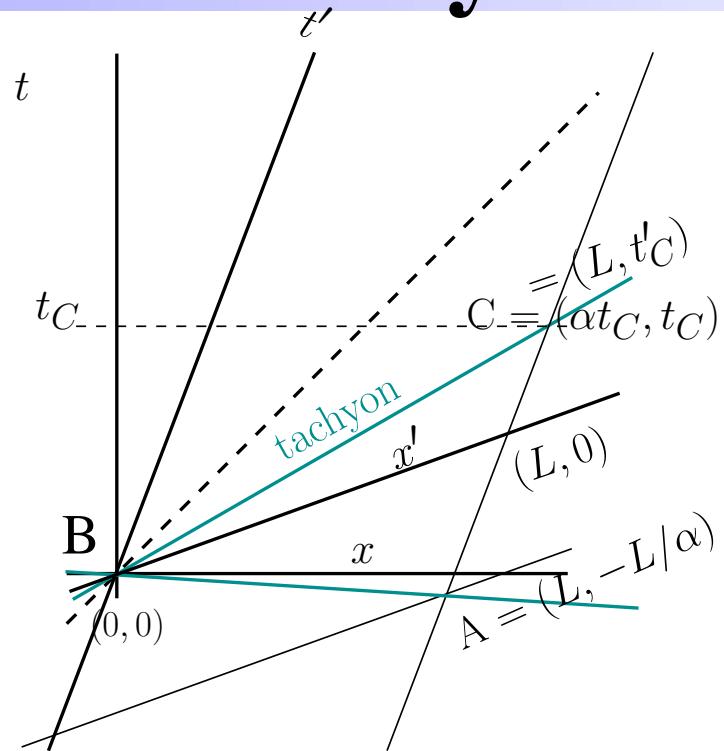
$$< 0 \text{ if } \beta > \frac{2\alpha}{\alpha^2 + 1}$$

A receives tachyonic response at C before sending it





# SR: tachyonic antitelephone



B stationary:  $(x, t)$   
frame  
A moving at speed  $\beta$ :  
 $(x', t')$  frame

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$$t'_C - t'_A = L \frac{2\alpha - (\alpha^2 + 1)\beta}{\alpha(\alpha - \beta)}$$

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A receives tachyonic response at C before sending it

w:tachyonic antitelephone



# SR: pole-barn/ladder paradox





# SR: pole-barn/ladder paradox

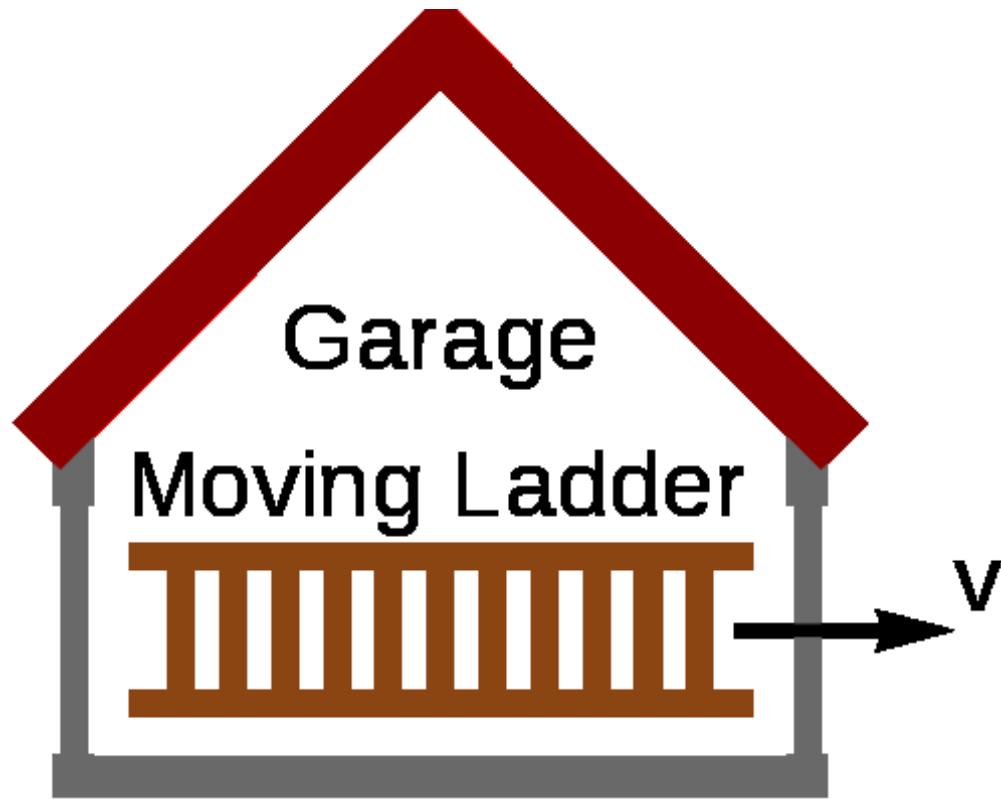


- ladder of length  $29.9\gamma$  ns, garage length 30 ns





# SR: pole-barn/ladder paradox

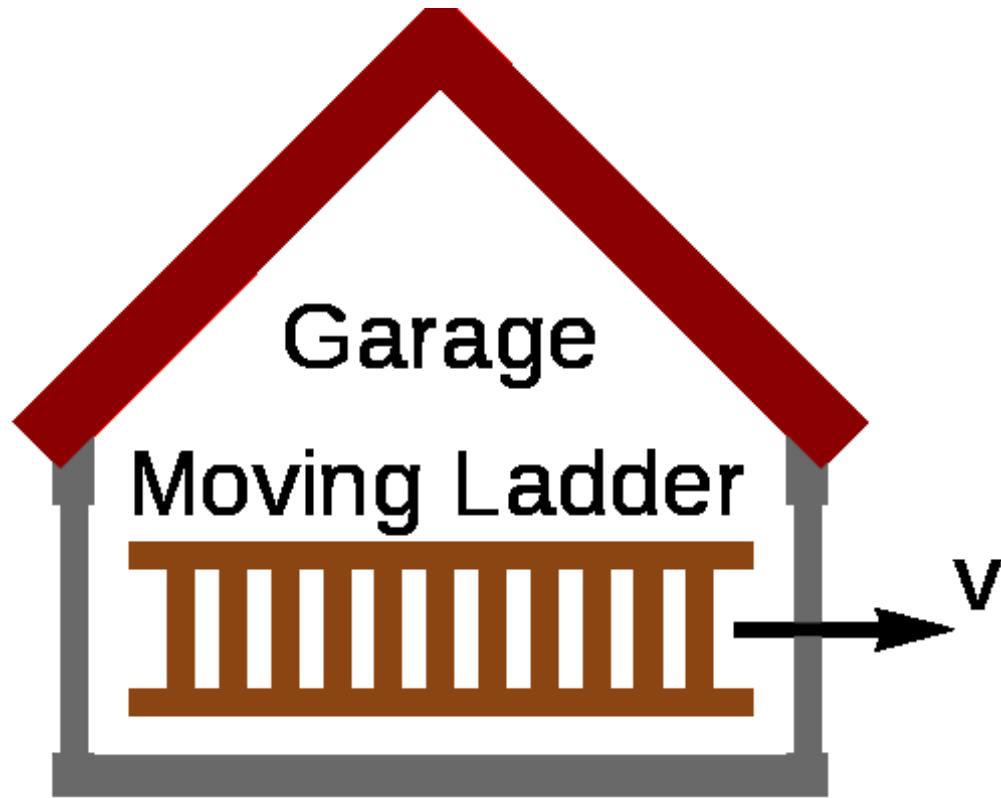


- ladder of length  $29.9\gamma$  ns, garage length 30 ns
- instantaneously close front + back doors





# SR: pole-barn/ladder paradox



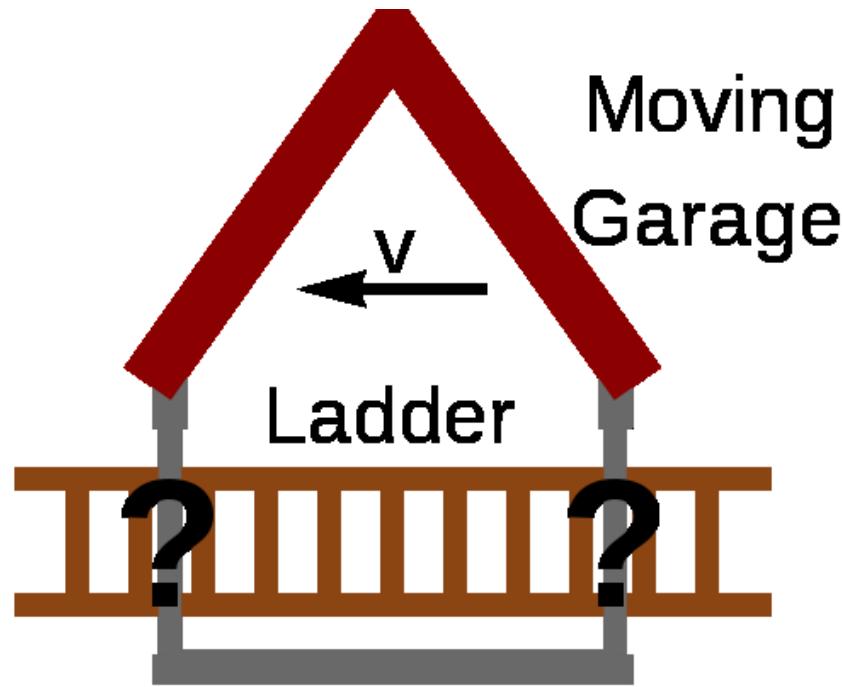
Wikimedia  
COMMONS

- ladder of length  $29.9\gamma$  ns, garage length 30 ns
- instantaneously close front + back doors
- $29.9\gamma$  ns /  $\gamma < 30$  ns  $\Rightarrow$  OK





# SR: pole-barn/ladder paradox



Wikimedia Commons

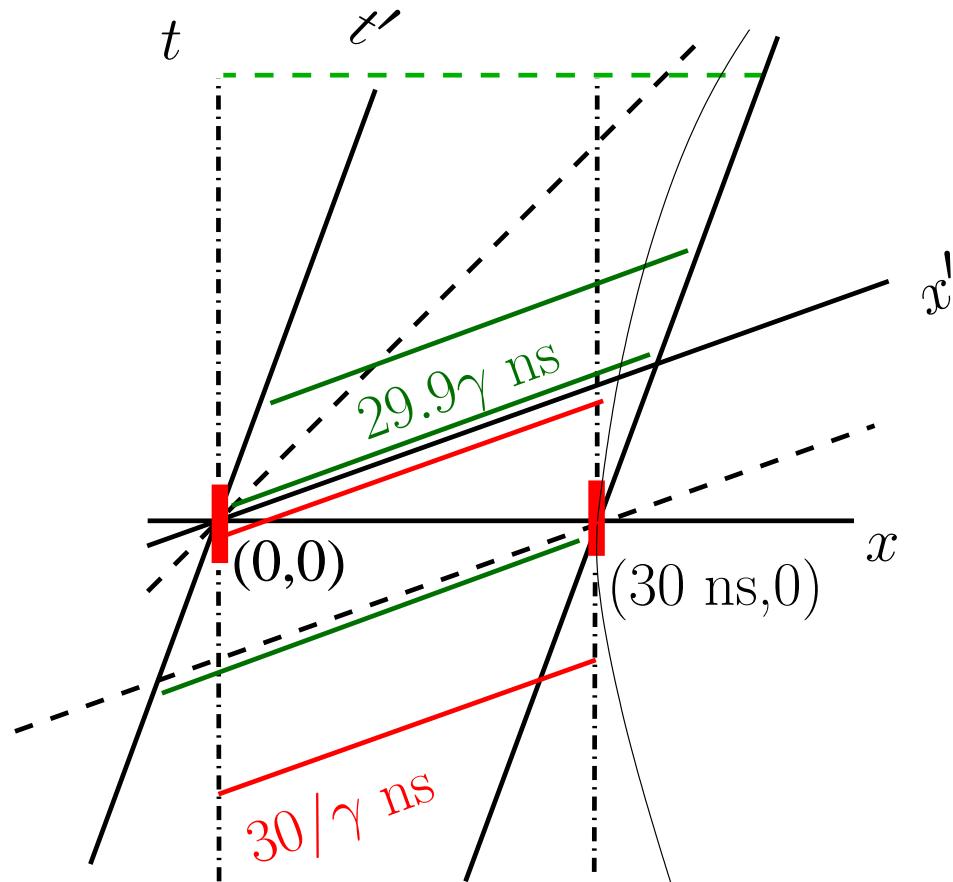
- ladder of length  $29.9\gamma$  ns, garage length 30 ns
- instantaneously close front + back doors
- ladder frame: garage  $30/\gamma$  ns long  $\ll 29.9\gamma$  ns!!

Is this possible or not? Make a spacetime diagram.



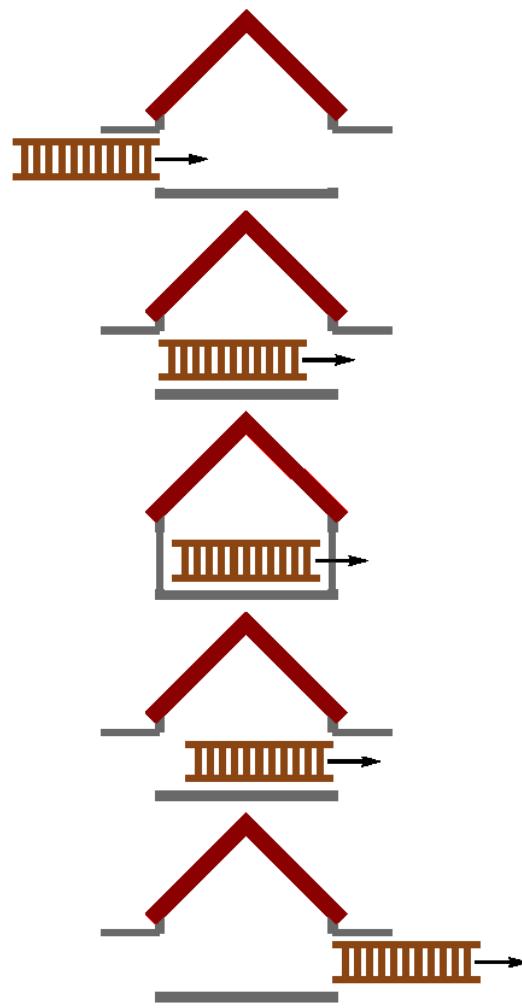


# SR: pole-barn/ladder paradox



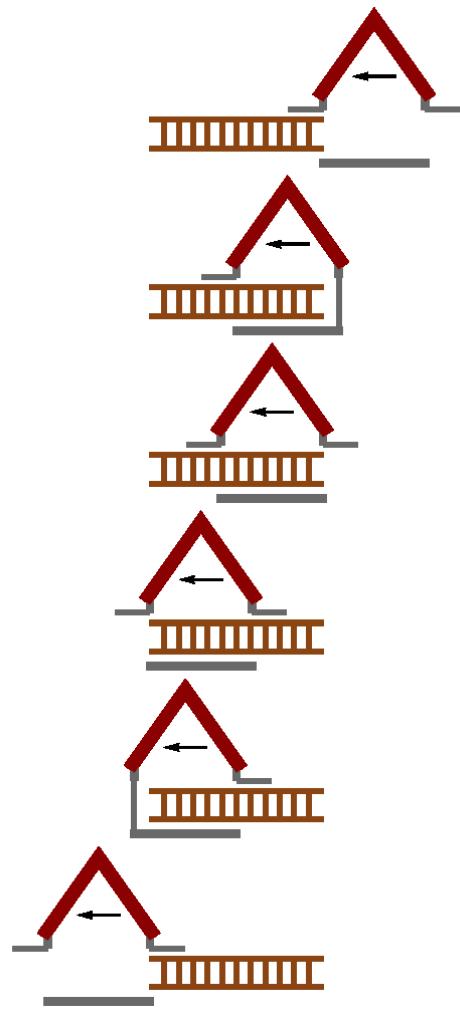


# SR: pole-barn/ladder paradox



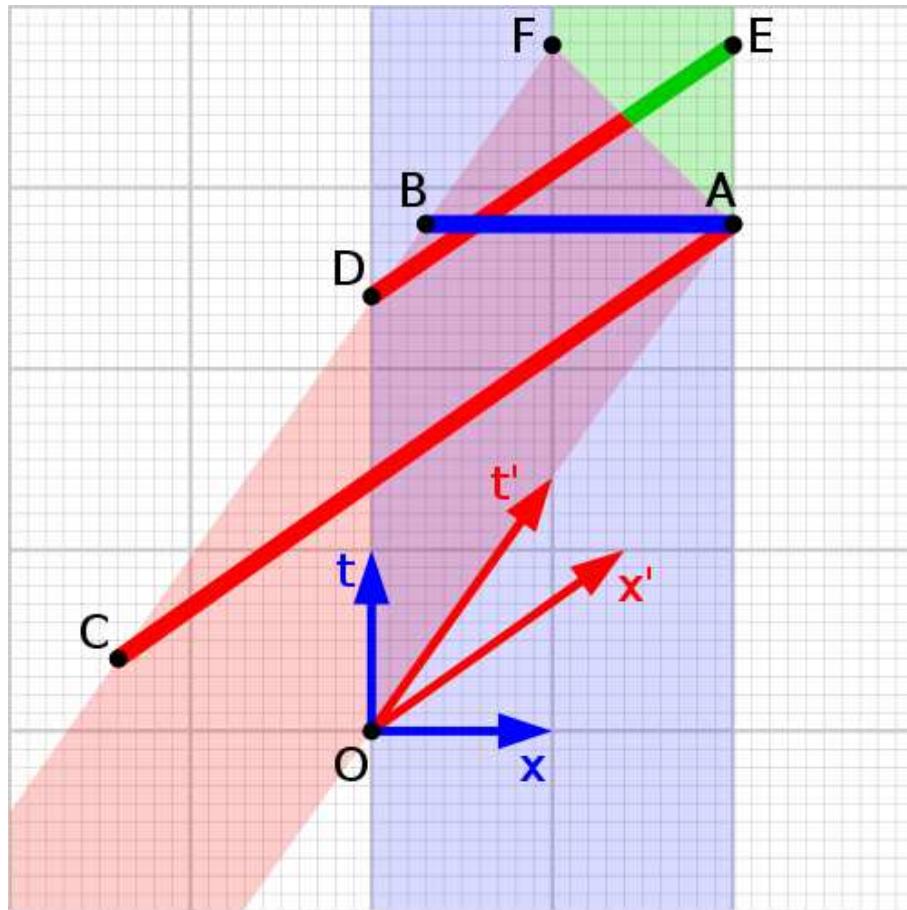


# SR: pole-barn/ladder paradox





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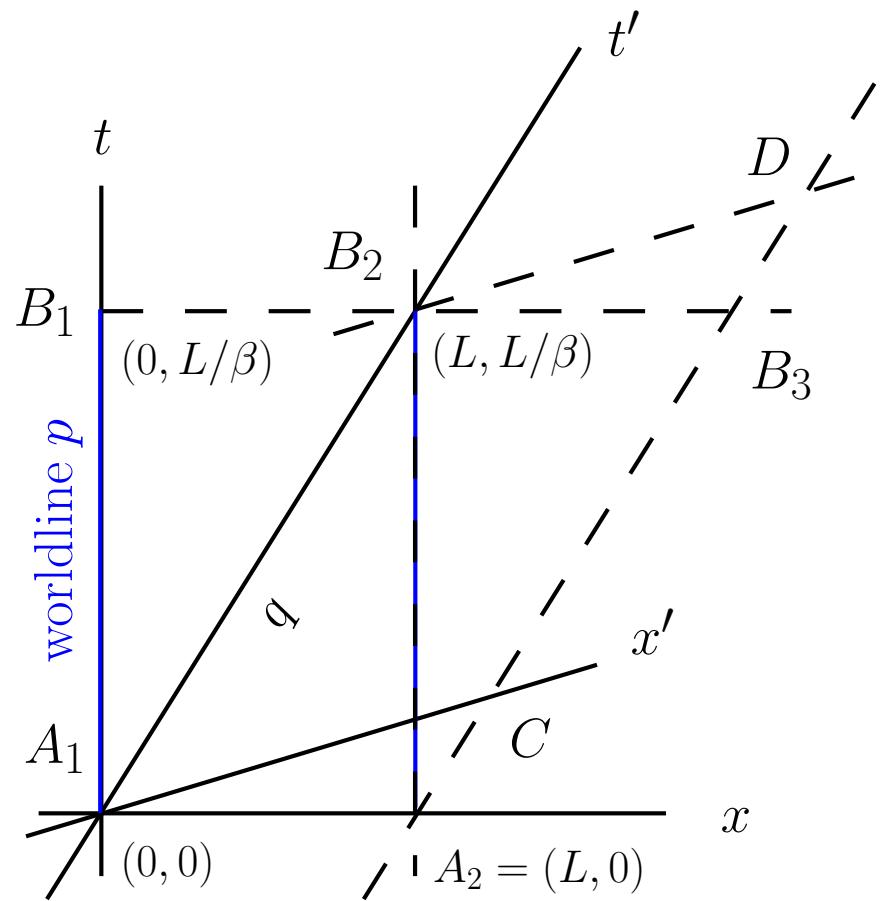


[w:Ladder paradox](#)





# SR: twins paradox

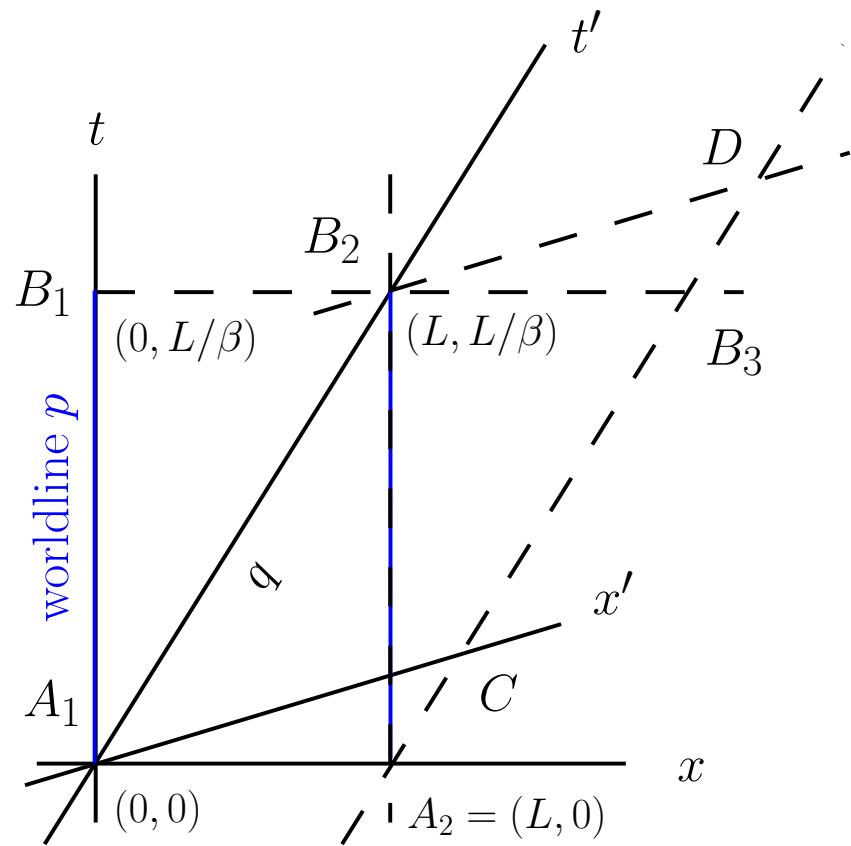


simply connected Minkowski





# SR: twins paradox

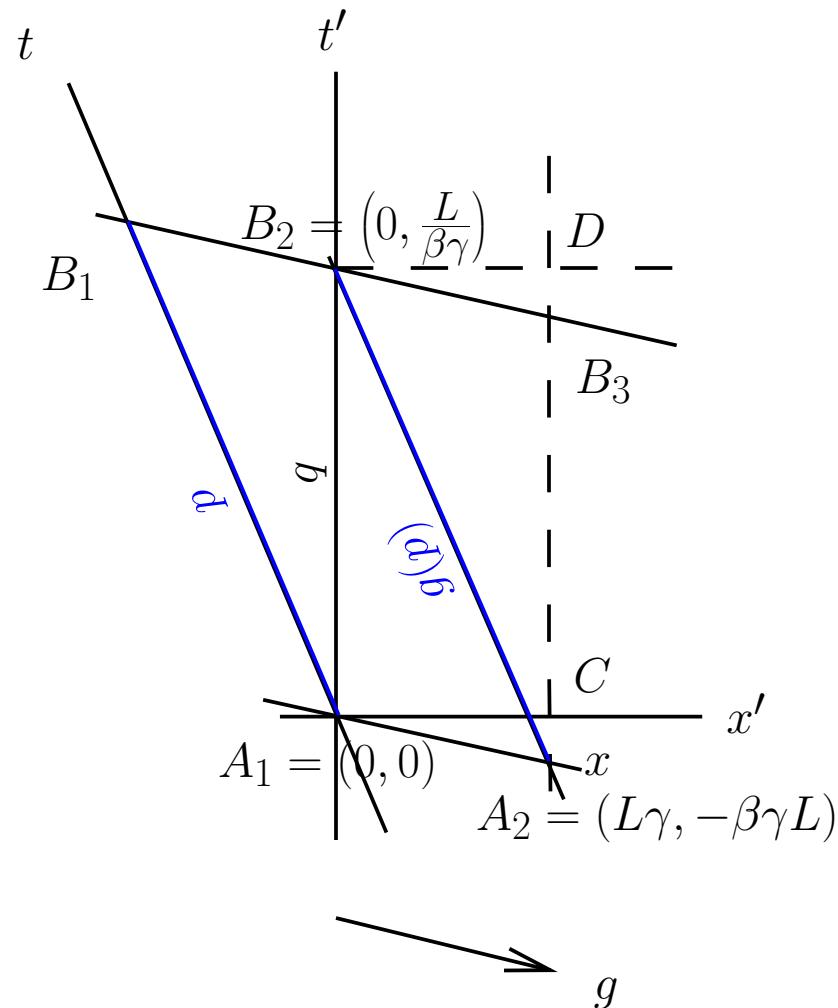


holonomy  $g$   
 $\xrightarrow{\hspace{1cm}}$   
identify spacetime events

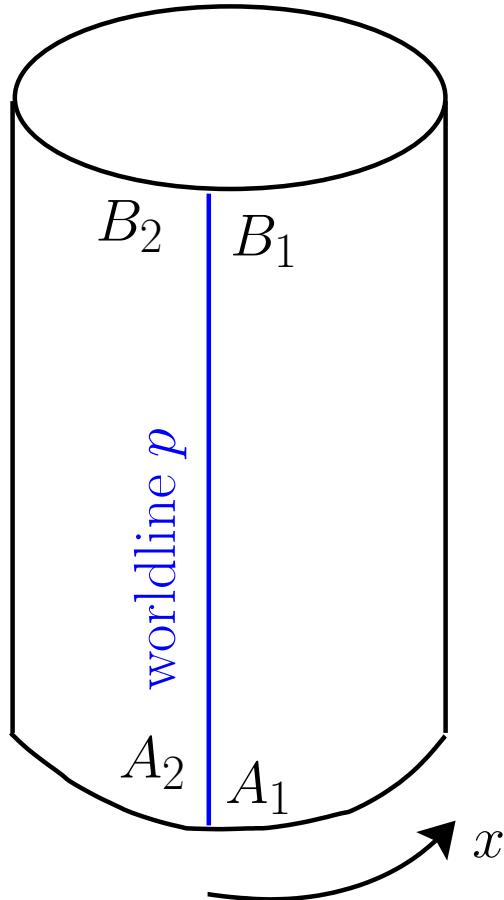




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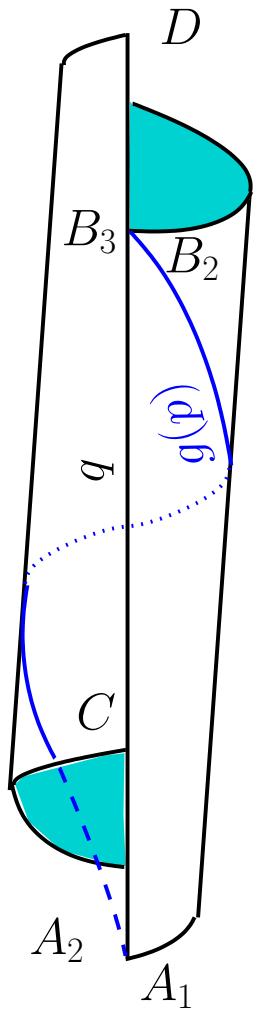


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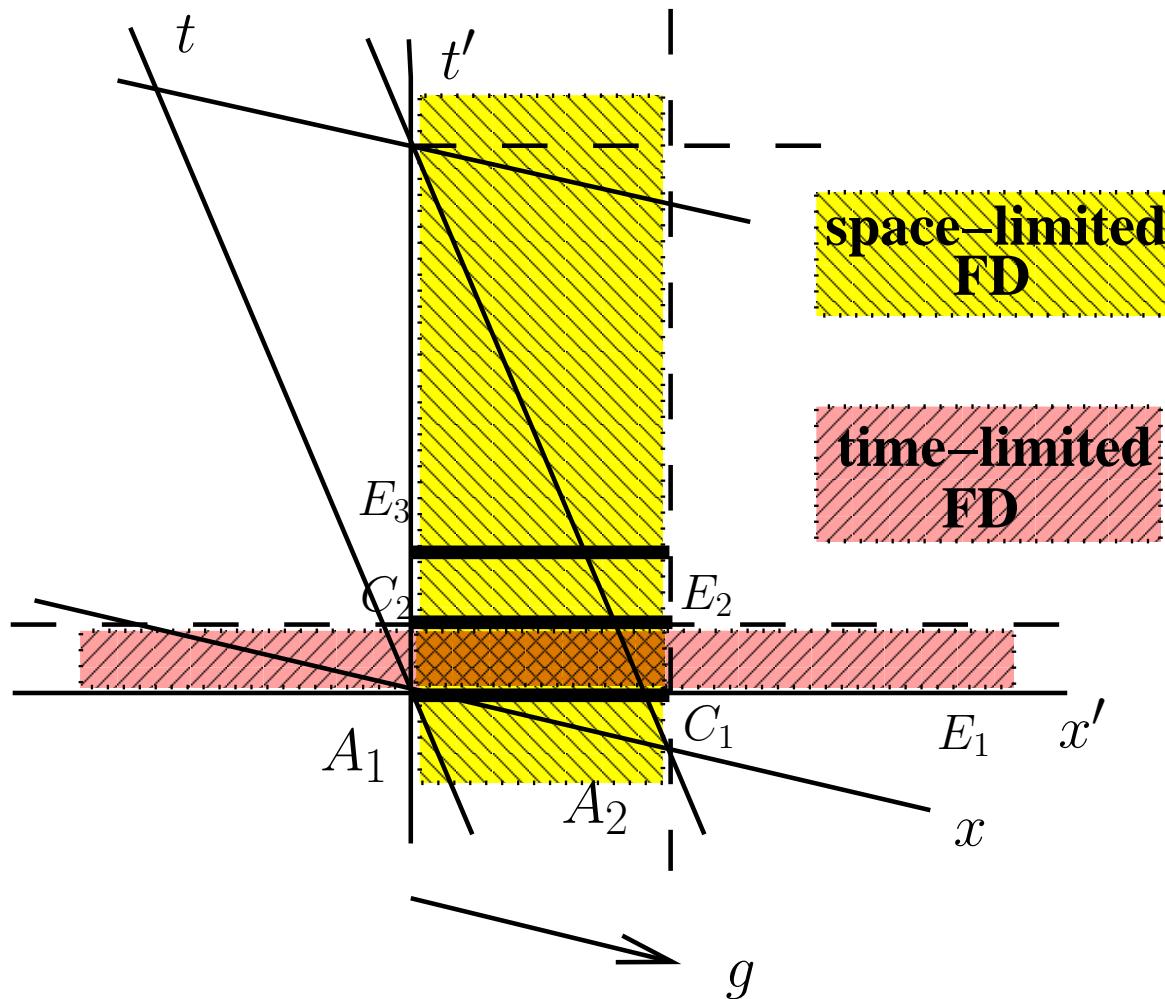




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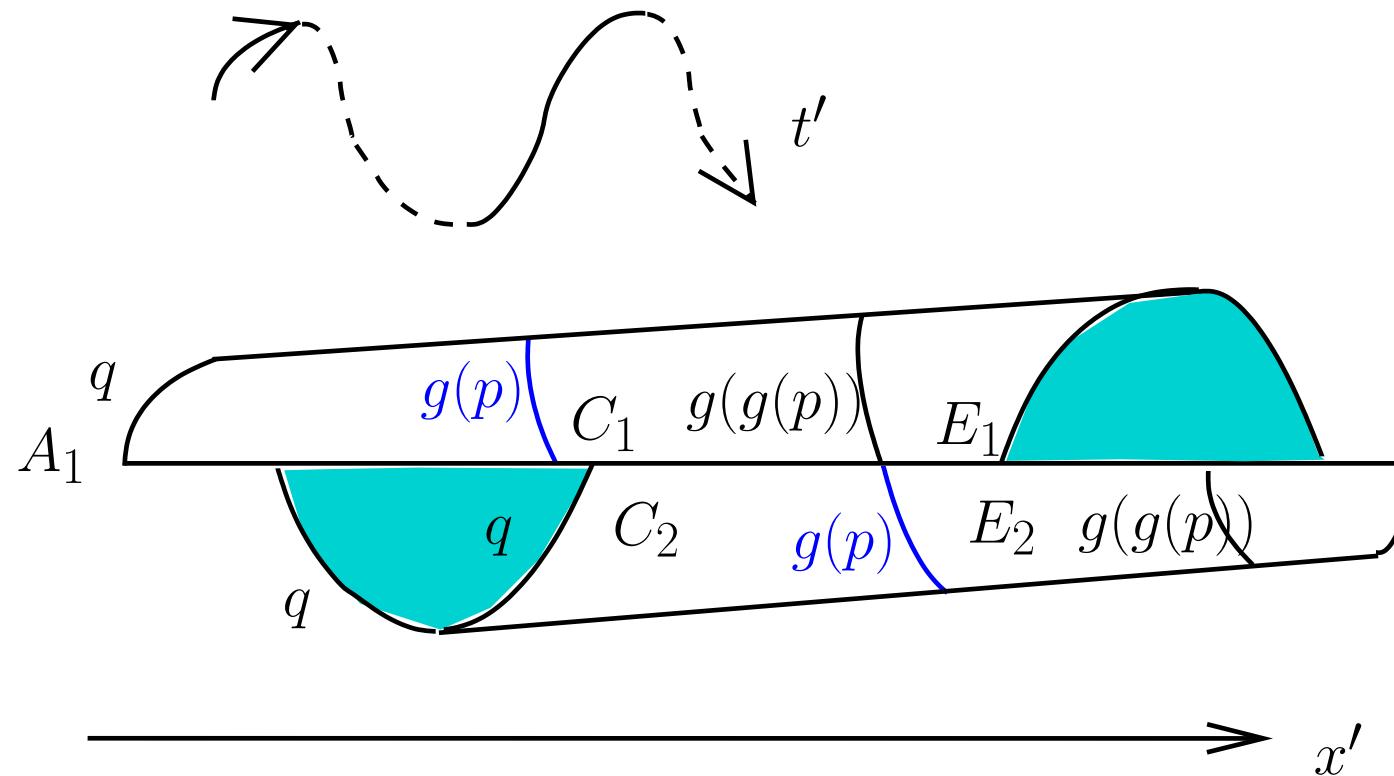


# SR: twins paradox





# SR: twins paradox



Roukema & Bajtlik 2008, MNRAS, 390, 655  
arXiv:astro-ph/0612155



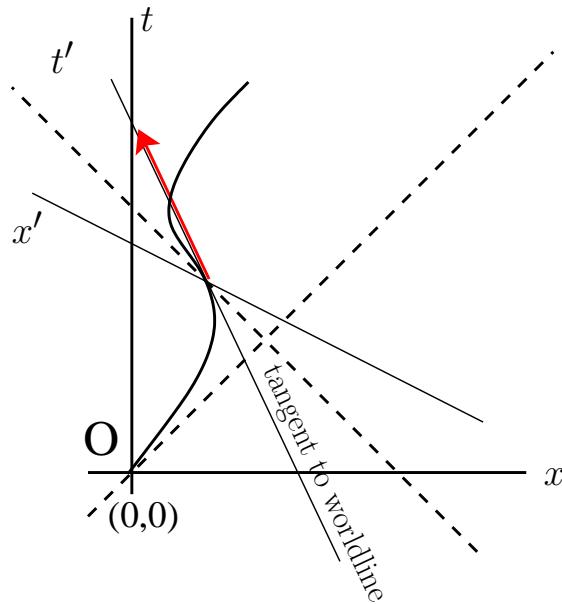
# SR: four-velocity, four-momentum

choose  $x$  axis so that 3-velocity  $u_{\text{Galilean}} = (\beta, 0, 0)^T$  for observer with  $(t, x, y, z)^T$  coord system



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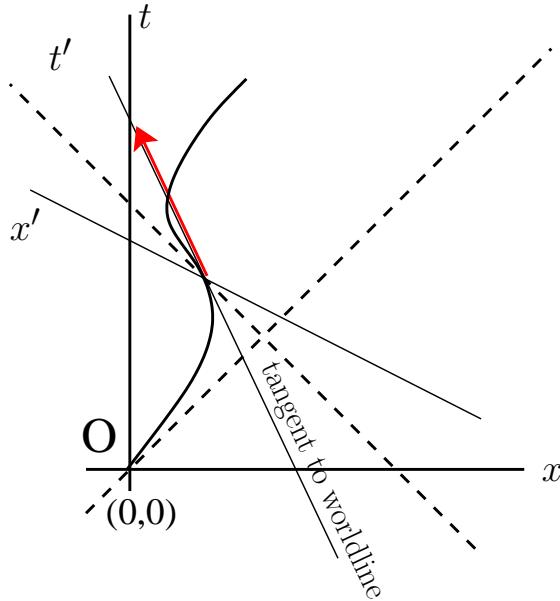
- in  $(t, x)$  spacetime 2-plane, extend from scalar speed  $\beta$  to spacetime vector = tangent to worldline

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$$(u^t, u^x) := \left( \frac{d}{d\tau} t(\tau), \frac{d}{d\tau} x(\tau) \right)$$

w:four-velocity



- in  $(t, x)$  spacetime 2-plane, extend from scalar speed  $\beta$  to spacetime vector =
- L** tangent to worldline

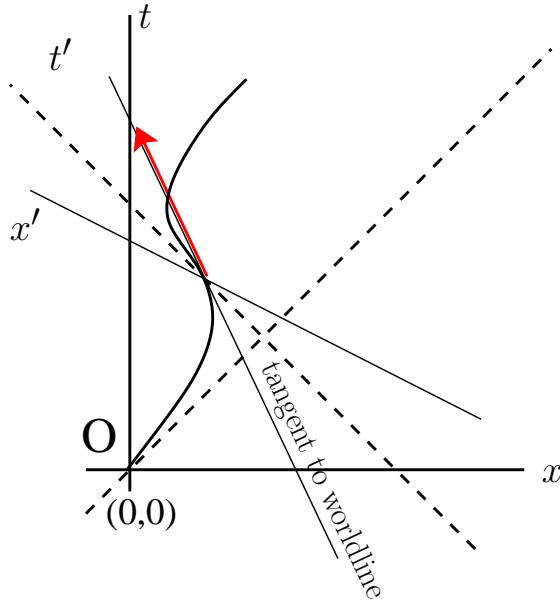
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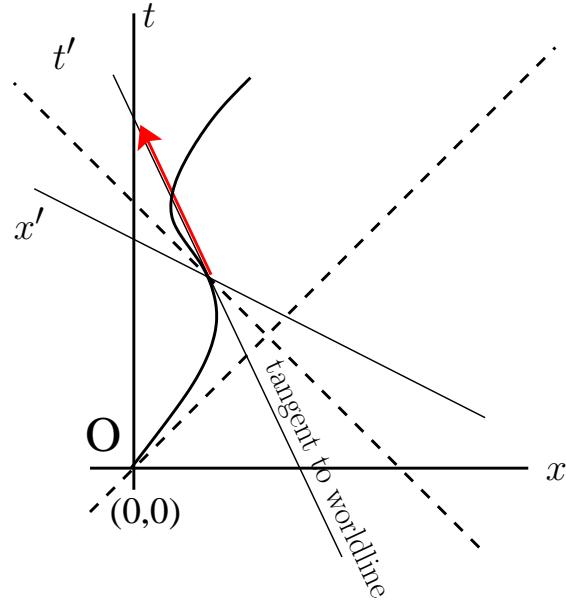
similarly  $(u^{t'}, u^{x'}) = \left( \frac{d}{d\tau} t'(\tau), \frac{d}{d\tau} x'(\tau) \right)$



in  $(t, x)$  spacetime  
2-plane, extend from

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similarly  $(u^{t'}, u^{x'}) = \left( \frac{d}{d\tau} t'(\tau), \frac{d}{d\tau} x'(\tau) \right) = (1, 0)$

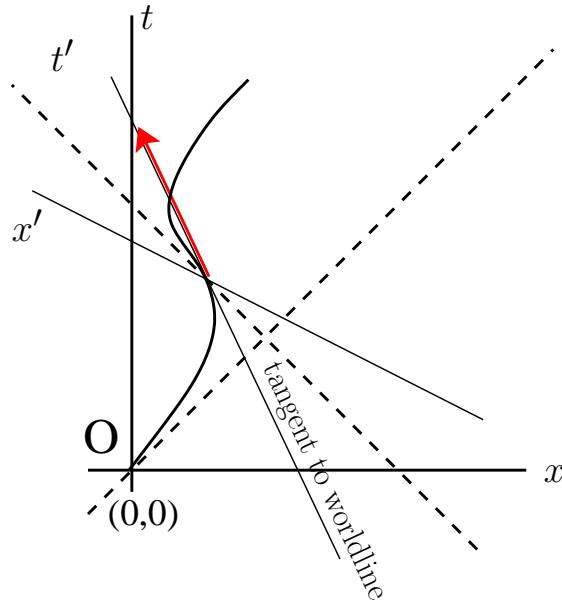
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$$\text{similarly } (u^{t'}, u^{x'}) = \left( \frac{d}{d\tau} t'(\tau), \frac{d}{d\tau} x'(\tau) \right) = (1, 0)$$

want  $\vec{u}$  Lorentz invariant  $\Rightarrow$   
 $(u^t, u^x)^T = \Lambda^{-1}(1, 0)^T$

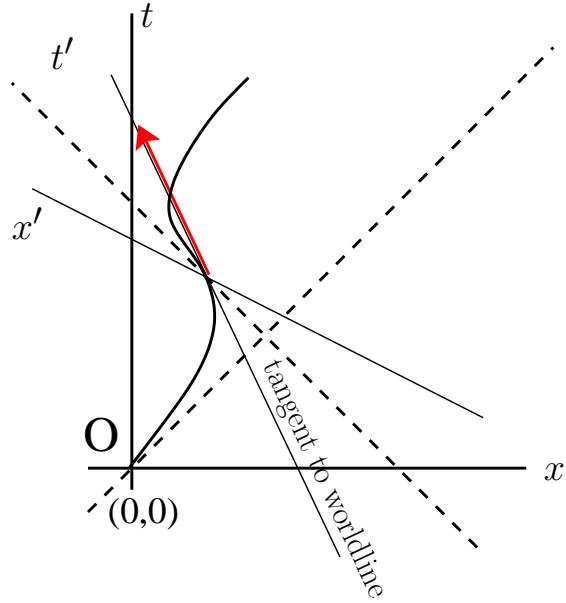
- in  $(t, x)$  spacetime

L<sup>2</sup>-plane, extend from

scalar speed  $\beta$  to  
spacetime vector

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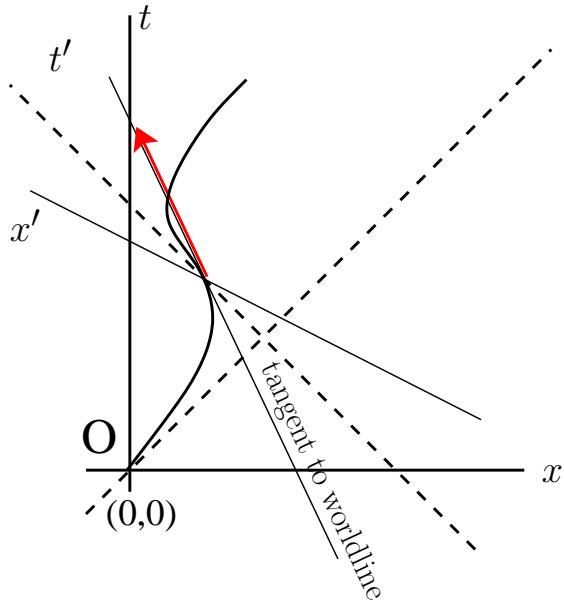
similarly  $(u^{t'}, u^{x'}) = \left( \frac{d}{d\tau} t'(\tau), \frac{d}{d\tau} x'(\tau) \right) = (1, 0)$

want  $\vec{u}$  Lorentz invariant  
 $\Rightarrow (u^t, u^x)^T = \Lambda^{-1}(1, 0)^T = \gamma(1, \beta)^T$

- in  $(t, x)$  spacetime 2-plane, extend from scalar speed  $\beta$  to spacetime vector = tangent to worldline

# SR: four-velocity, four-momentum

choose  $x$  axis so that 3-velocity  $u_{\text{Galilean}} = (\beta, 0, 0)^T$  for observer with  $(t, x, y, z)$  coord system



$$(u^t, u^x) := \left( \frac{d}{d\tau} t(\tau), \frac{d}{d\tau} x(\tau) \right)$$

w:four-velocity

$$\text{similarly } (u^{t'}, u^{x'}) = \left( \frac{d}{d\tau} t'(\tau), \frac{d}{d\tau} x'(\tau) \right) = (1, 0)$$

want  $\vec{u}$  Lorentz invariant  
 $\Rightarrow (u^t, u^x)^T = \Lambda^{-1}(1, 0)^T = \gamma(1, \beta)^T$

4D:  $\vec{u} = \gamma(1, \beta^x, \beta^y, \beta^z)^T$

notation in this pdf:

$\vec{u}$  = 4-vector,  ${}^{(3)}\vec{u}$  = spatial component

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Is the 3-component (spatial component) of  $\vec{u}$  the same as the non-relativistic velocity?



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$$= \gamma \frac{d}{dt}(x, y, z)^T$$

$$\neq \frac{d}{dt}(x, y, z)^T \text{ except if } \beta = 0 \Leftrightarrow \gamma = 1$$



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momentum:  $\vec{p} := m\vec{u} = m\gamma(1, \beta^x, \beta^y, \beta^z)^T$ , where  $m = \text{constant}$  w:invariant mass

<sup>x</sup> ... = tensor-style component notation, not powers



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What does the time component of momentum =  $p^0 = m\gamma$  mean physically?

- first look at spatial component in a given ref. frame



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$$\begin{aligned}(3) \vec{p} &= m \frac{d}{d\tau} (x, y, z)^T \\&= m\gamma \frac{d}{dt} (x, y, z)^T \\&\neq m \frac{d}{dt} (x, y, z)^T \text{ except if } \beta = 0 \Leftrightarrow \gamma = 1\end{aligned}$$



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let us define 4-acceleration, 4-force



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# SR: invariance of ${}^{(4)}u$ , ${}^{(4)}a$ , ${}^{(4)}f$

Euclidean norm:  $\|\vec{x}\|^2 = \sum_\mu (x^\mu)^2$





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w:Einstein summation sum is implicit





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$\delta_{ij} = 1$  if  $i = j$ , otherwise = 0;  $i, j \in 1, 2, 3$

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similarly:  $\|\vec{a}\|^2$ ,  $\|\vec{f}\|^2$  invariant



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3-force :=  $\frac{d}{dt}^{(3)}\vec{p} \neq \frac{d}{d\tau}^{(3)}\vec{p}$





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$$\frac{d}{dt} {}^{(3)}\vec{p} = \frac{{}^{(3)}\vec{f}}{\gamma}$$



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$$\begin{aligned} &= \int_0^{\beta_2} \frac{(3)\vec{f}}{\gamma} \cdot d\vec{x} \\ &= \int_0^{\beta_2} \frac{d}{dt} (m\beta\gamma) dx \end{aligned}$$

(assume  $(3)\vec{f}/\gamma \parallel \vec{x}$ )



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$$\Rightarrow K + m = m\gamma \text{ drop "2"}$$





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so  $p^0$  = kinetic energy + rest mass



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momentum time component:

$$\begin{aligned} p^0 &= m\gamma = m(1 - \beta^2)^{-1/2} \\ &= m[1 - (1/2)(-\beta^2) + \mathcal{O}(\beta^4)] \text{ if } \beta \ll 1 \end{aligned}$$



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Yes.





# SR: $\vec{p} \dots$ : invariant or not?

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**WARNING:** assume that 4-momentum vectors at different space-time positions can be parallel-transported; not the case in curved spacetime





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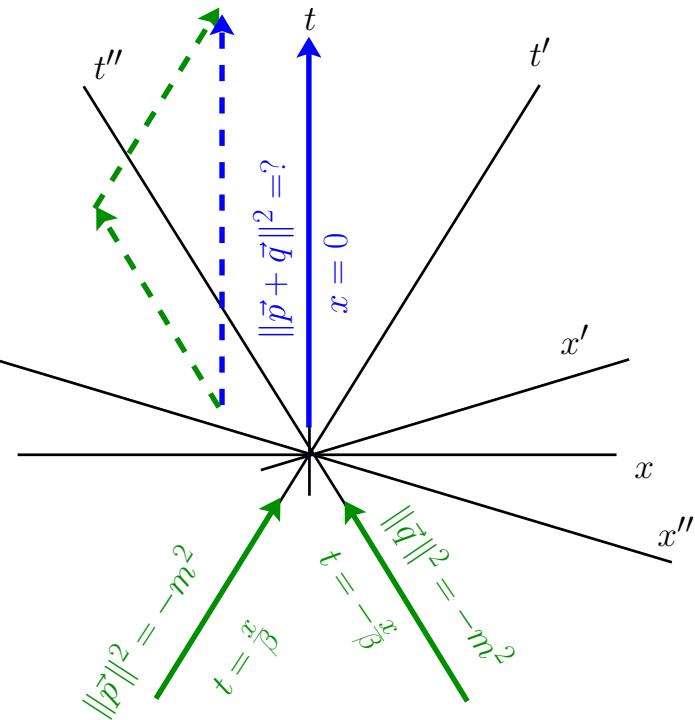
vector space  $\Rightarrow p^0 + q^0 = r^0$

= conservation of (relativistic) “total energy” =  $m + K$   
(Newtonian:  $m$  conserved,  $K$  not conserved,  $K+$  potential energy conserved)

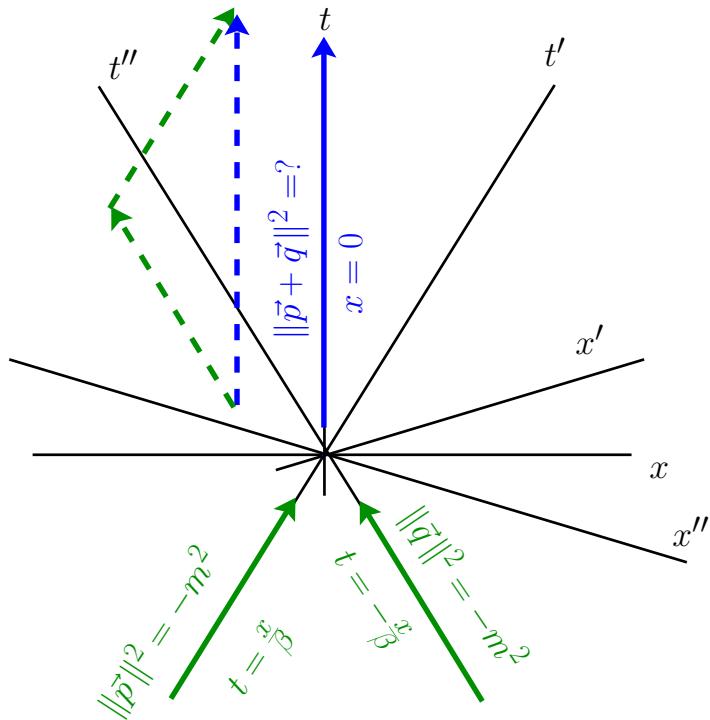
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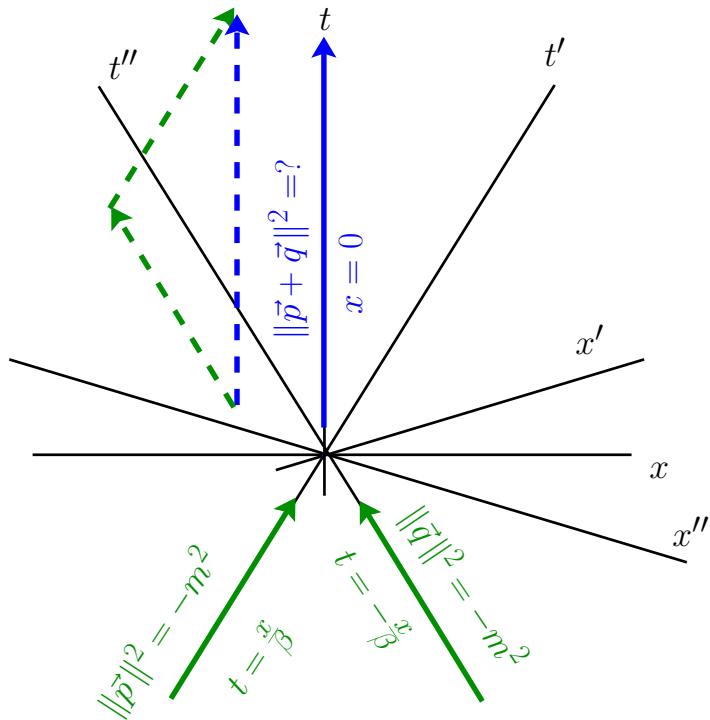
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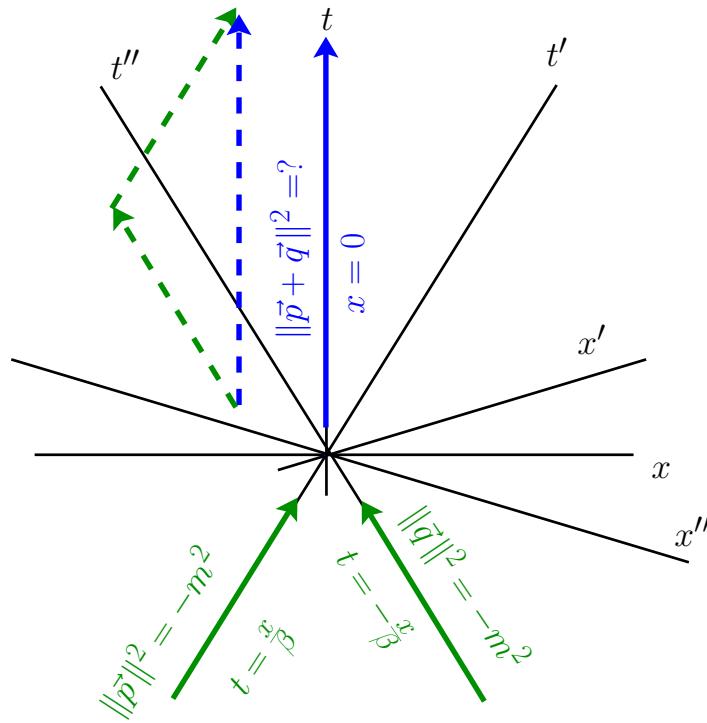
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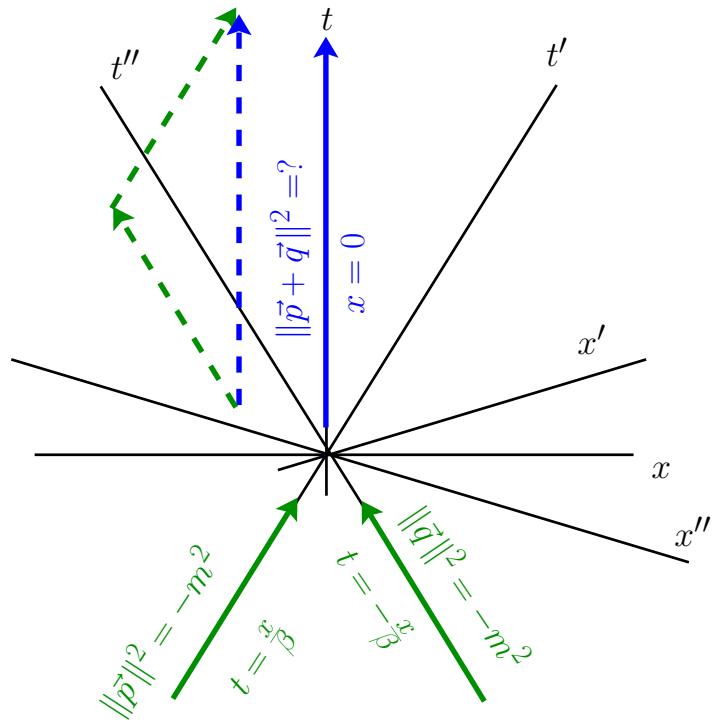
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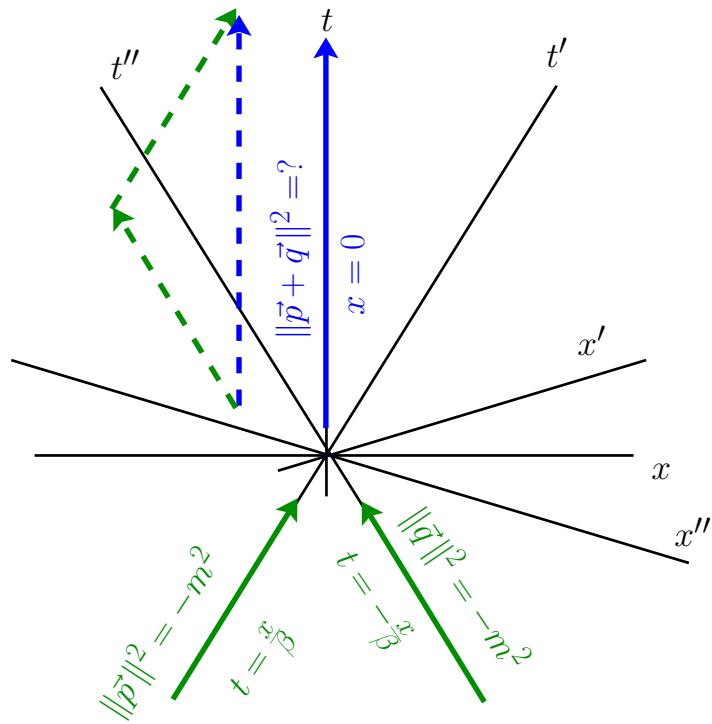
$$\vec{q} = m(\gamma, -\beta\gamma, 0, 0)^T$$

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$$L = m[2\gamma, (-\beta + \beta)\gamma, 0, 0]^T$$



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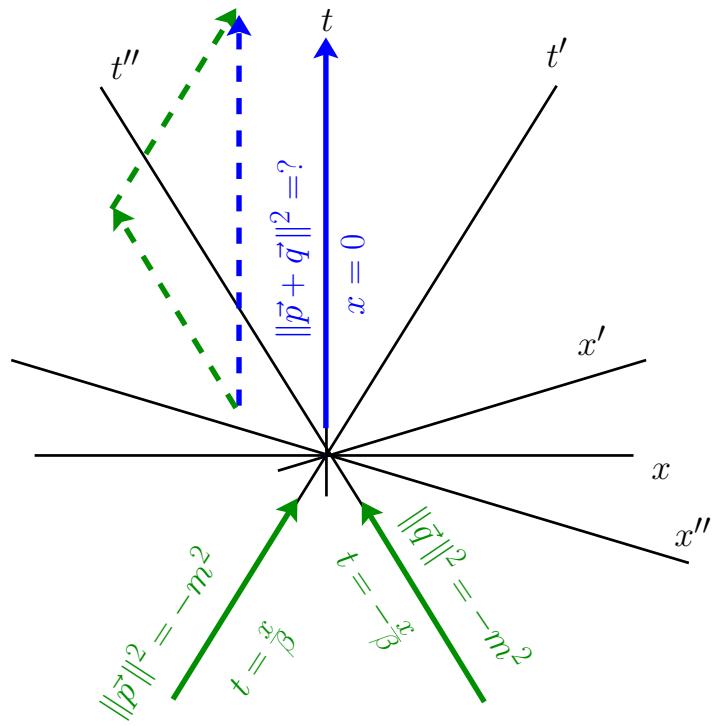
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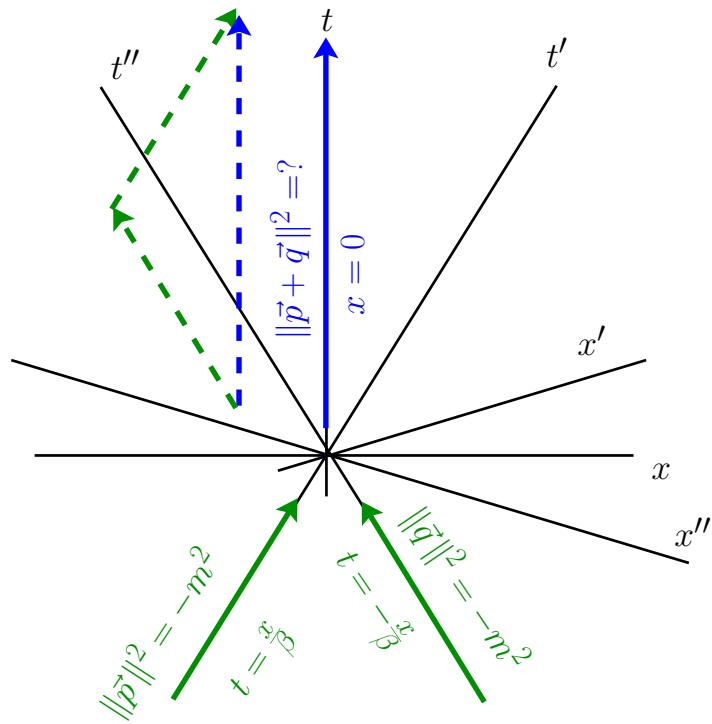
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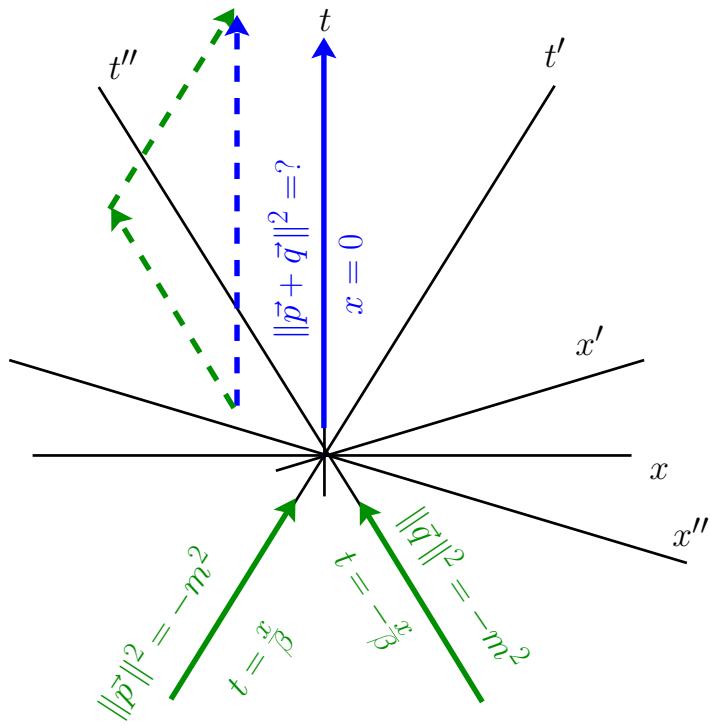
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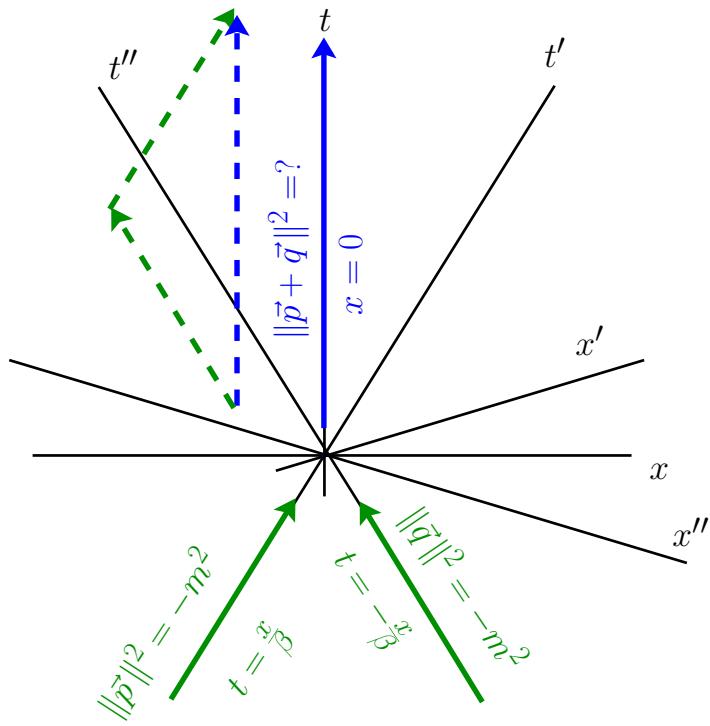
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refuse the assumption of absolute simultaneity (time)

